

# Jean Bourgain and quasiperiodic (ergodic) Schrödinger operators

May 23, 2016



$$(h_\theta \Psi)_n = \Delta \Psi + \lambda v(T^n \theta) \Psi_n \quad (0.1)$$

acting on  $\ell^2(\mathbb{Z}^d)$ ,

$T$  is an ergodic action of  $\mathbb{Z}^d$  on a probability space  $(\Theta, \mu)$ .

$v$  is a sampling function

$$(h_\theta \Psi)_n = \Delta \Psi + \lambda v(T^n \theta) \Psi_n \quad (0.1)$$

acting on  $\ell^2(\mathbb{Z}^d)$ ,

$T$  is an ergodic action of  $\mathbb{Z}^d$  on a probability space  $(\Theta, \mu)$ .

$v$  is a sampling function

Spectrum, and spectral components are the same for a.e.  $\theta$ .

$$(h_\theta \Psi)_n = \Delta \Psi + \lambda v(T^n \theta) \Psi_n \quad (0.1)$$

acting on  $\ell^2(\mathbb{Z}^d)$ ,

$T$  is an ergodic action of  $\mathbb{Z}^d$  on a probability space  $(\Theta, \mu)$ .

$v$  is a sampling function

Spectrum, and spectral components are the same for a.e.  $\theta$ .

## Quasiperiodic operators:

$$\Theta = \mathbb{T}^b$$

$$T_i \theta = \theta + \omega_i, \quad \omega \text{ incommensurate.}$$

Periodic if  $\omega$  rational.

$$(h_\theta \Psi)_n = \Delta \Psi + \lambda v(T^n \theta) \Psi_n \quad (0.1)$$

acting on  $\ell^2(\mathbb{Z}^d)$ ,

$T$  is an ergodic action of  $\mathbb{Z}^d$  on a probability space  $(\Theta, \mu)$ .

$v$  is a sampling function

Spectrum, and spectral components are the same for a.e.  $\theta$ .

## Quasiperiodic operators:

$$\Theta = \mathbb{T}^b$$

$$T_i \theta = \theta + \omega_i, \quad \omega \text{ incommensurate.}$$

Periodic if  $\omega$  rational.

## Anderson model:

$v(T^n \theta)$  i.i.d.r.v.

# Exotic Spectral properties

- Metal-insulator transitions
- Dense point spectrum
- Unusual absolutely continuous spectrum
- Eigenfunctions decaying at the non-Lyapunov rate
- Singular continuous spectrum
- Cantor spectrum
- Unusual eigenvalue statistics

- Anderson localization (Nobel, 1977)
- Integer Quantum Hall Effect (Nobel, 1998)
- Graphene (Nobel, 2010)

# Prizes

- Anderson localization (Nobel, 1977)
- Integer Quantum Hall Effect (Nobel, 1998)
- Graphene (Nobel, 2010)
- Fields medal 2014





# Quasiperiodic operators: KAM methods

$$(H_{\lambda,\alpha,\theta}\Psi)_n = \Psi_{n+1} + \Psi_{n-1} + \lambda v(T^n\theta)\Psi_n$$

$$v(T^n\theta) = v(\theta + n\alpha)$$

$\theta, \alpha \in \mathbb{T}^d$ , Diophantine  $\alpha$

**Small coupling**

Dinaburg-Sinai (76), Eliasson (91)

- KAM in the momentum space
- Reducibility of transfer-matrix cocycles

**Large coupling,  $d = 1$**

Sinai, Fröhlich-Spencer-Wittwer (late 80s) cos-type

Eliasson (97), analytic (Gevrey)

# Quasiperiodic operators: KAM methods

$$(H_{\lambda,\alpha,\theta}\Psi)_n = \Psi_{n+1} + \Psi_{n-1} + \lambda v(T^n\theta)\Psi_n$$

$$v(T^n\theta) = v(\theta + n\alpha)$$

$\theta, \alpha \in \mathbb{T}^d$ , Diophantine  $\alpha$

**Small coupling**

Dinaburg-Sinai (76), Eliasson (91)

- KAM in the momentum space
- Reducibility of transfer-matrix cocycles

**Large coupling,  $d = 1$**

Sinai, Fröhlich-Spencer-Wittwer (late 80s) cos-type

Eliasson (97), analytic (Gevrey)

# Bourgain: robust non-KAM methods

Direct estimates of the Green's function in the regime of positive Lyapunov exponents (large coupling). No multiple scales!  
non-perturbative ac spectrum by dual localization

# Bourgain: robust non-KAM methods

Direct estimates of the Green's function in the regime of positive Lyapunov exponents (large coupling). No multiple scales!  
non-perturbative ac spectrum by dual localization

Bourgain (+Goldstein, Schlag, J.) (2000-2005): analytic,  
multi-frequency, multidimensional

# Bourgain: robust non-KAM methods

Direct estimates of the Green's function in the regime of positive Lyapunov exponents (large coupling). No multiple scales!  
non-perturbative ac spectrum by dual localization

Bourgain (+Goldstein, Schlag, J.) (2000-2005): analytic,  
multi-frequency, multidimensional  
applications e.g. to random NLS (Bourgain, 2005, Bourgain-Wang,  
2008)



1. Massiveness of the spectrum (Bourgain, 2005)

**Corollary:** absolutely continuous spectrum for multi-dimensional quasiperiodic operators at low coupling  
(Bourgain, J., Parnovsky, 2016)

2. Continuity of the Lyapunov exponent of analytic  $SL(2, \mathbb{R})$  cocycles **in frequency** at irrational frequencies (Bourgain-J, 2002)

2. Continuity of the Lyapunov exponent of analytic  $SL(2, \mathbb{R})$  cocycles **in frequency** at irrational frequencies (Bourgain-J, 2002)  
optimal:  
can be discontinuous at rational frequencies  
can be discontinuous in  $C^\infty$ , (Wang-You, 2014)



2. Continuity of the Lyapunov exponent of analytic  $SL(2, \mathbb{R})$  cocycles **in frequency** at irrational frequencies (Bourgain-J, 2002)

optimal:

can be discontinuous at rational frequencies

can be discontinuous in  $C^\infty$ , (Wang-You, 2014)

**Prequel:**

Hölder continuity in energy for Diophantine frequencies  
(Goldstein-Schlag, 2001)

2. Continuity of the Lyapunov exponent of analytic  $SL(2, \mathbb{R})$  cocycles **in frequency** at irrational frequencies (Bourgain-J, 2002)

optimal:

can be discontinuous at rational frequencies

can be discontinuous in  $C^\infty$ , (Wang-You, 2014)

**Prequel:**

Hölder continuity in energy for Diophantine frequencies  
(Goldstein-Schlag, 2001)

**Sequels:**

2. Continuity of the Lyapunov exponent of analytic  $SL(2, \mathbb{R})$  cocycles **in frequency** at irrational frequencies (Bourgain-J, 2002)

optimal:

can be discontinuous at rational frequencies

can be discontinuous in  $C^\infty$ , (Wang-You, 2014)

**Prequel:**

Hölder continuity in energy for Diophantine frequencies  
(Goldstein-Schlag, 2001)

**Sequels:**

- Ten Martini problem (Avila-J, 2009)
- Avila's global theory of quasiperiodic cocycles (Avila, 2009-2015)

2. Continuity of the Lyapunov exponent of analytic  $SL(2, \mathbb{R})$  cocycles **in frequency** at irrational frequencies (Bourgain-J, 2002)

optimal:

can be discontinuous at rational frequencies

can be discontinuous in  $C^\infty$ , (Wang-You, 2014)

**Prequel:**

Hölder continuity in energy for Diophantine frequencies  
(Goldstein-Schlag, 2001)

**Sequels:**

- Ten Martini problem (Avila-J, 2009)
- Avila's global theory of quasiperiodic cocycles (Avila, 2009-2015)

Bourgain, 2005: the same continuity for multifrequency

# Bernoulli-Anderson model

Large coupling (or bottom of the spectrum for continuous) implies  
Anderson localization (pure point spectrum with exponentially  
decaying eigenfunctions)

# Bernoulli-Anderson model

Large coupling (or bottom of the spectrum for continuous) implies  
Anderson localization (pure point spectrum with exponentially  
decaying eigenfunctions)

Multi-scale analysis: Fröhlich-Spencer (1983)  
Klein + coauthors (1989-2004)

# Bernoulli-Anderson model

Large coupling (or bottom of the spectrum for continuous) implies Anderson localization (pure point spectrum with exponentially decaying eigenfunctions)

Multi-scale analysis: Fröhlich-Spencer (1983)  
Klein + coauthors (1989-2004)

Fractional moments: Aizenman-Molchanov,  
Aizenman+collaborators (1994+)

# Bernoulli-Anderson model

Large coupling (or bottom of the spectrum for continuous) implies Anderson localization (pure point spectrum with exponentially decaying eigenfunctions)

Multi-scale analysis: Fröhlich-Spencer (1983)  
Klein + coauthors (1989-2004)

Fractional moments: Aizenman-Molchanov,  
Aizenman+collaborators (1994+)

Both rely on continuity of the distribution of randomness  
(essentially requiring continuous rank one (compact) perturbations)



# Bernoulli-Anderson model

Large coupling (or bottom of the spectrum for continuous) implies Anderson localization (pure point spectrum with exponentially decaying eigenfunctions)

Multi-scale analysis: Fröhlich-Spencer (1983)  
Klein + coauthors (1989-2004)

Fractional moments: Aizenman-Molchanov,  
Aizenman+collaborators (1994+)

Both rely on continuity of the distribution of randomness (essentially requiring continuous rank one (compact) perturbations)  
Bourgain-Kenig (2005): a.e. localization near the bottom of the spectrum for continuous operators with Bernoulli potential.

# Bernoulli-Anderson model

Large coupling (or bottom of the spectrum for continuous) implies Anderson localization (pure point spectrum with exponentially decaying eigenfunctions)

Multi-scale analysis: Fröhlich-Spencer (1983)  
Klein + coauthors (1989-2004)

Fractional moments: Aizenman-Molchanov,  
Aizenman+collaborators (1994+)

Both rely on continuity of the distribution of randomness (essentially requiring continuous rank one (compact) perturbations)

Bourgain-Kenig (2005): a.e. localization near the bottom of the spectrum for continuous operators with Bernoulli potential.  
new powerful combinatorics techniques

# Bernoulli-Anderson model

Large coupling (or bottom of the spectrum for continuous) implies Anderson localization (pure point spectrum with exponentially decaying eigenfunctions)

Multi-scale analysis: Fröhlich-Spencer (1983)  
Klein + coauthors (1989-2004)

Fractional moments: Aizenman-Molchanov,  
Aizenman+collaborators (1994+)

Both rely on continuity of the distribution of randomness (essentially requiring continuous rank one (compact) perturbations)

Bourgain-Kenig (2005): a.e. localization near the bottom of the spectrum for continuous operators with Bernoulli potential.

new powerful combinatorics techniques

Still open in the discrete case!

# Integrated density of states

One of Simon's 15 problems in mathematical physics (1985): prove continuity of the Integrated density of states for continuum operators

# Integrated density of states

One of Simon's 15 problems in mathematical physics (1985): prove continuity of the Integrated density of states for continuum operators

Bourgain-Klein (2012): proof for  $d = 2, 3$ .

# Integrated density of states

One of Simon's 15 problems in mathematical physics (1985): prove continuity of the Integrated density of states for continuum operators

Bourgain-Klein (2012): proof for  $d = 2, 3$ .

Still open for  $d > 3$ ...

# Almost Mathieu operators

$$(H_{\lambda,\alpha,\theta}\Psi)_n = \Psi_{n+1} + \Psi_{n-1} + \lambda v(T^n\theta)\Psi_n$$

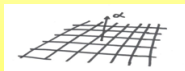
$$v(T^n\theta) = 2 \cos 2\pi(\theta + n\alpha), \alpha \text{ irrational},$$

# Almost Mathieu operators

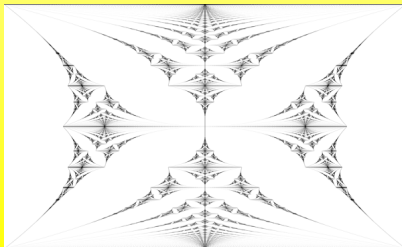
$$(H_{\lambda,\alpha,\theta}\Psi)_n = \Psi_{n+1} + \Psi_{n-1} + \lambda v(T^n\theta)\Psi_n$$

$$v(T^n\theta) = 2 \cos 2\pi(\theta + n\alpha), \alpha \text{ irrational},$$

Tight-binding model of 2D Bloch electrons in magnetic fields



## Hofstadter butterfly





$$(H_{\lambda,\alpha,\theta}\Psi)_n = \Psi_{n+1} + \Psi_{n-1} + 2\lambda \cos 2\pi(\theta + n\alpha)\Psi_n$$

- $|\lambda| < 1$  :  $H_{\lambda,\alpha,\theta}$  has purely absolutely continuous spectrum for all  $\alpha, \theta$  (Avila, 2008).

# Arithmetic spectral transitions

$$(H_{\lambda,\alpha,\theta}\Psi)_n = \Psi_{n+1} + \Psi_{n-1} + 2\lambda \cos 2\pi(\theta + n\alpha)\Psi_n$$

- $|\lambda| < 1$  :  $H_{\lambda,\alpha,\theta}$  has purely absolutely continuous spectrum for all  $\alpha, \theta$  (Avila, 2008).
- $|\lambda| > 1$  : arithmetic transitions.

$$\beta(\alpha) \equiv \limsup_{n \rightarrow \infty} - \frac{\ln \|n\alpha\|_{\mathbb{R}/\mathbb{Z}}}{|n|}$$

# Arithmetic spectral transitions

$$(H_{\lambda,\alpha,\theta}\Psi)_n = \Psi_{n+1} + \Psi_{n-1} + 2\lambda \cos 2\pi(\theta + n\alpha)\Psi_n$$

- $|\lambda| < 1$  :  $H_{\lambda,\alpha,\theta}$  has purely absolutely continuous spectrum for all  $\alpha, \theta$  (Avila, 2008).
- $|\lambda| > 1$  : arithmetic transitions.

$$\beta(\alpha) \equiv \limsup_{n \rightarrow \infty} - \frac{\ln \|n\alpha\|_{\mathbb{R}/\mathbb{Z}}}{|n|}$$

and

$$\delta(\alpha, \theta) \equiv \limsup_{n \rightarrow \infty} - \frac{\ln \|2\theta + n\alpha\|_{\mathbb{R}/\mathbb{Z}}}{|n|}$$

# Arithmetic spectral transitions

$$(H_{\lambda,\alpha,\theta}\Psi)_n = \Psi_{n+1} + \Psi_{n-1} + 2\lambda \cos 2\pi(\theta + n\alpha)\Psi_n$$

- $|\lambda| < 1$  :  $H_{\lambda,\alpha,\theta}$  has purely absolutely continuous spectrum for all  $\alpha, \theta$  (Avila, 2008).
- $|\lambda| > 1$  : arithmetic transitions.

$$\beta(\alpha) \equiv \limsup_{n \rightarrow \infty} - \frac{\ln \|n\alpha\|_{\mathbb{R}/\mathbb{Z}}}{|n|}$$

and

$$\delta(\alpha, \theta) \equiv \limsup_{n \rightarrow \infty} - \frac{\ln \|2\theta + n\alpha\|_{\mathbb{R}/\mathbb{Z}}}{|n|}$$

Diophantine  $\alpha$  :  $\beta(\alpha) = 0$ ,

$\alpha$ -Diophantine  $\theta$  :  $\delta(\alpha, \theta) = 0$

$$(H_{\lambda,\alpha,\theta}\Psi)_n = \Psi_{n+1} + \Psi_{n-1} + \lambda \cos 2\pi(\theta + n\alpha)\Psi_n$$

J.- Liu (2015-2016):

- For Diophantine  $\alpha$ , there is a sharp transition at  $\lambda_0 = e^{\delta(\alpha,\theta)}$   
 $|\lambda| < e^{\delta(\alpha,\theta)} \implies H_{\lambda,\alpha,\theta}$  has purely singular continuous spectrum  
 $|\lambda| > e^{\delta(\alpha,\theta)} \implies H_{\lambda,\alpha,\theta}$  has Anderson localization.
- For  $\alpha$ -Diophantine  $\theta$  (any  $\alpha$ ) there is a sharp transition at  $\lambda_0 = e^{\beta(\alpha)}$   
 $|\lambda| < e^{\beta(\alpha)} \implies H_{\lambda,\alpha,\theta}$  has purely singular continuous spectrum ( Avila-You-Zhou, 2015),  
 $|\lambda| > e^{\beta(\alpha)} \implies H_{\lambda,\alpha,\theta}$  has Anderson localization.

$$(H_{\lambda,\alpha,\theta}\Psi)_n = \Psi_{n+1} + \Psi_{n-1} + \lambda \cos 2\pi(\theta + n\alpha)\Psi_n$$

J.- Liu (2015-2016):

- For Diophantine  $\alpha$ , there is a sharp transition at  $\lambda_0 = e^{\delta(\alpha,\theta)}$   
 $|\lambda| < e^{\delta(\alpha,\theta)} \implies H_{\lambda,\alpha,\theta}$  has purely singular continuous spectrum  
 $|\lambda| > e^{\delta(\alpha,\theta)} \implies H_{\lambda,\alpha,\theta}$  has Anderson localization.
- For  $\alpha$ -Diophantine  $\theta$  (any  $\alpha$ ) there is a sharp transition at  $\lambda_0 = e^{\beta(\alpha)}$   
 $|\lambda| < e^{\beta(\alpha)} \implies H_{\lambda,\alpha,\theta}$  has purely singular continuous spectrum ( Avila-You-Zhou, 2015),  
 $|\lambda| > e^{\beta(\alpha)} \implies H_{\lambda,\alpha,\theta}$  has Anderson localization.

Ten Martini paper: localization for  $|\lambda| > e^{16/9\beta(\alpha)}$ .

## Theorem 1

*In both cases, throughout the point spectrum regime, suppose  $E$  is an eigenvalue of  $H_{\lambda,\alpha,\theta}$  and  $\phi$  is the eigenfunction.*

## Theorem 1

*In both cases, throughout the point spectrum regime, suppose  $E$  is an eigenvalue of  $H_{\lambda,\alpha,\theta}$  and  $\phi$  is the eigenfunction. Let*

$$U(k) = \begin{pmatrix} \phi(k) \\ \phi(k-1) \end{pmatrix}.$$



## Theorem 1

*In both cases, throughout the point spectrum regime, suppose  $E$  is an eigenvalue of  $H_{\lambda,\alpha,\theta}$  and  $\phi$  is the eigenfunction. Let*

*$U(k) = \begin{pmatrix} \phi(k) \\ \phi(k-1) \end{pmatrix}$ . Then for any  $\varepsilon > 0$ , there exists  $K$  such that for any  $|k| \geq K$ ,  $U(k)$  satisfies*

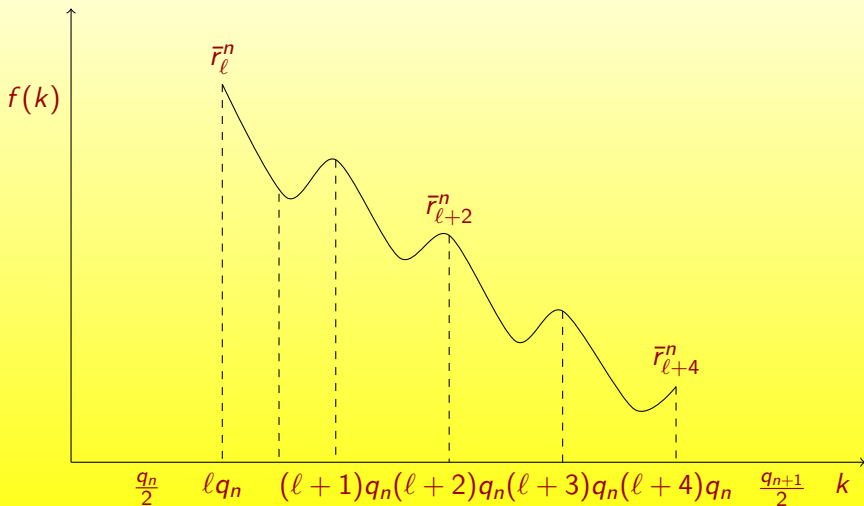
## Theorem 1

*In both cases, throughout the point spectrum regime, suppose  $E$  is an eigenvalue of  $H_{\lambda,\alpha,\theta}$  and  $\phi$  is the eigenfunction. Let*

*$U(k) = \begin{pmatrix} \phi(k) \\ \phi(k-1) \end{pmatrix}$ . Then for any  $\varepsilon > 0$ , there exists  $K$  such that for any  $|k| \geq K$ ,  $U(k)$  satisfies*

$$f(|k|)e^{-\varepsilon|k|} \leq \|U(k)\| \leq f(|k|)e^{\varepsilon|k|},$$

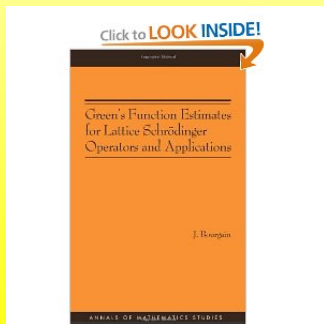
# The behavior of $f(k)$



Wencai Liu:



# Wencai Liu:



Incidentally, I found the reading of Jean's papers as a graduate student to be simultaneously extremely frustrating and extremely rewarding. Decoding an offhand remark or a mysterious step in his paper was often as instructive (and as time-consuming) as reading several pages of arguments by some other authors. (But his papers do become much easier to read once one has internalised enough of his "box of tools"...)

[share](#) [cite](#) [improve this answer](#)

[edited Jul 10 '12 at 20:48](#)

[answered Jul 10 '12 at 17:37](#)



**Terry Tao**

43.4k ● 14 ● 198 ● 282

Incidentally, I found the reading of Jean's papers as a graduate student to be simultaneously extremely frustrating and extremely rewarding. Decoding an offhand remark or a mysterious step in his paper was often as instructive (and as time-consuming) as reading several pages of arguments by some other authors. (But his papers do become much easier to read once one has internalised enough of his "box of tools"...)

[share](#) [cite](#) [improve this answer](#)

[edited Jul 10 '12 at 20:48](#)

[answered Jul 10 '12 at 17:37](#)



**Terry Tao**

43.4k • 14 • 198 • 282

- 
- 12 Well, Terry, I found the reading of Jean's papers as a full professor to be simultaneously extremely frustrating and extremely rewarding. – [Bill Johnson](#) [Jul 11 '12 at 0:57](#)