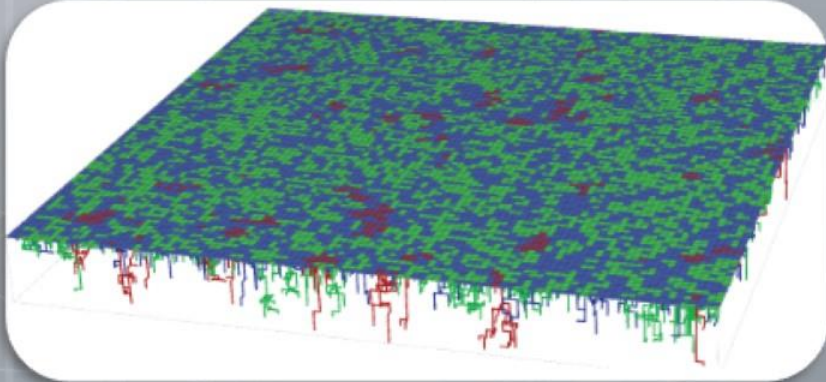


Nov 2014
IAS, Princeton

Information Percolation for the Ising model

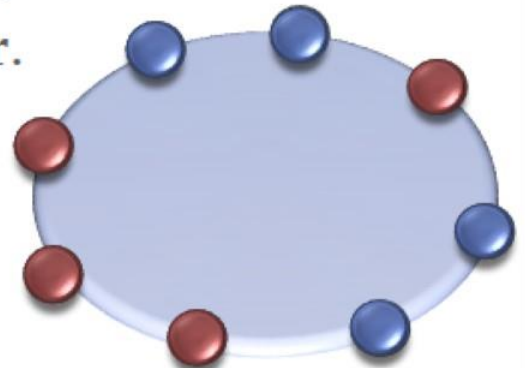


Eyal Lubetzky
Courant Institute (NYU)

Joint with Allan Sly

Noisy Election Day (on a cycle)

- ▶ Setup: (1D noisy voter model with noise $0 < \varepsilon < 1$)
 - n binary voters on a cycle.
 - Every step, a uniformly chosen voter updates its vote:
 - **prob. $1 - \varepsilon$** : copy a random neighbor.
 - **prob. ε** : new vote is a fair coin toss.



- ▶ How long does it take to reach equilibrium?
 - from all-1? from 01010...? from 001100...?
 - from a typical state? from a random IID state?

Definition: the classical Ising model

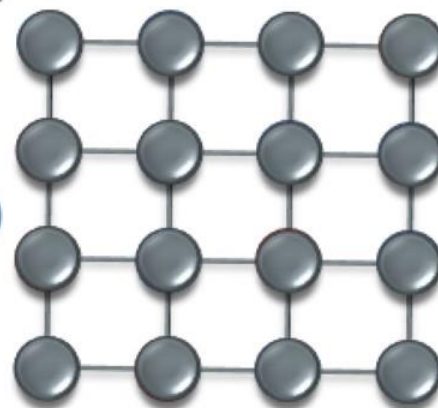
▶ Underlying geometry: $\Lambda =$ finite 2D grid.

▶ Set of possible configurations:

$$\Omega = \{\pm 1\}^\Lambda$$

(each *site* receives a plus/minus *spin*)

▶ Probability of a configuration $\sigma \in \Omega$ given by the *Gibbs distribution*:



$$\mu(\sigma) = \frac{1}{Z(\beta)} \exp\left(\beta \sum_{x \sim y} \sigma(x)\sigma(y)\right)$$

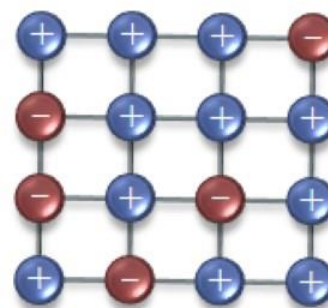
Partition
function

Inverse
temperature
 $\beta \geq 0$

The classical Ising model

▶ $\mu(\sigma) \propto \exp(\beta \sum_{x \sim y} \sigma(x)\sigma(y))$ for $\sigma \in \Omega = \{\pm 1\}^\Lambda$

- ▶ Larger β favors configurations with aligned spins at neighboring sites.
- ▶ Spin interactions: local, justified by rapid decay of magnetic force with distance.



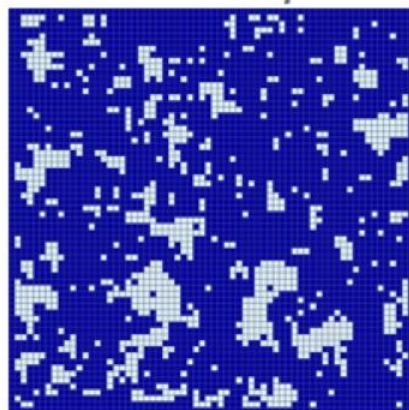
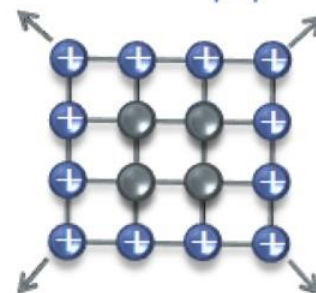
- ▶ The *magnetization* is the (normalized) sum of spins:

$$M(\sigma) = \frac{1}{|\Lambda|} \sum_{x \in \Lambda} \sigma(x)$$

- ▶ Distinguishes between disorder ($M \approx 0$) and order.
- ▶ Symmetry: $\mathbb{E}[M(\sigma)] = 0$. What if we *break the symmetry*?

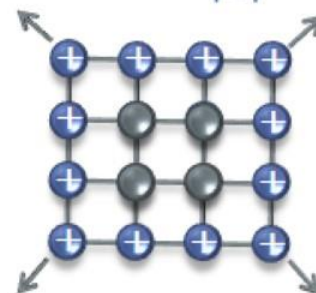
The Ising phase-transition

- ▶ Ferromagnetism in this setting: [recall $M(\sigma) = \frac{1}{|\Lambda|} \sum \sigma(x)$]
 - Condition on the boundary sites all having *plus* spins.
 - Let the system size $|\Lambda|$ tend $\rightarrow \infty$ (\approx a magnetic field with effect $\rightarrow 0$).
- ▶ What is the typical $M(\sigma)$ for large $|\Lambda|$?
Does the effect of *plus* boundary vanish in the limit?



The Ising phase-transition (ctd.)

- ▶ Ferromagnetism in this setting: [recall $M(\sigma) = \frac{1}{|\Lambda|} \sum \sigma(x)$]
 - Condition on the boundary sites all having *plus* spins.
 - Let the system size $|\Lambda|$ tend $\rightarrow \infty$



- ▶ Phase-transition at some critical β_c :

$$\lim_{|\Lambda| \rightarrow \infty} \mathbb{E}^+ [M(\sigma)] = \begin{cases} 0 & \text{if } \beta < \beta_c \\ m_\beta > 0 & \text{if } \beta > \beta_c \end{cases}$$

all-plus boundary

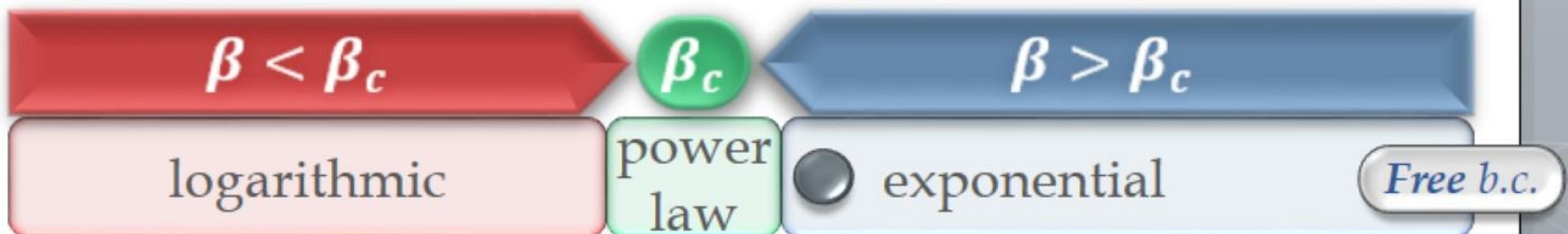
spontaneous magnetization

Static vs. stochastic Ising

- ▶ Expected behavior for the Ising distribution:



- ▶ Expected behavior for the mixing time of dynamics:



Glauber dynamics for Ising

(*a.k.a.* the *Stochastic Ising model*)

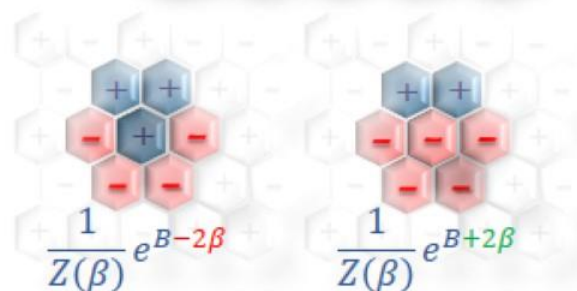
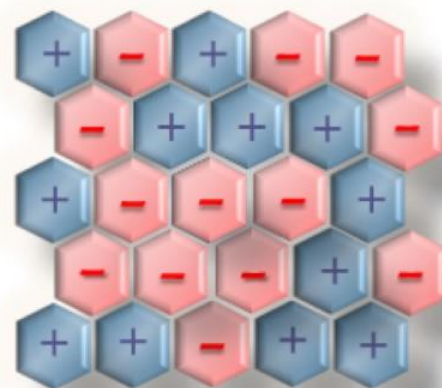
- ▶ Introduced in 1963 by Roy Glauber. (*heat-bath* version; famous other flavor: *Metropolis*)

Time-dependent statistics of the Ising model

RJ Glauber – *Journal of mathematical physics*, 1963

Cited by 2749

- ▶ One of the most commonly used samplers for the Ising distribution μ :
 - Update sites via IID Poisson(1) clocks
 - Each update replaces a spin at $x \in V$ by a new spin $\sim \mu$ given spins at $V \setminus \{x\}$.
- ▶ How long does it take it to converge to μ ?



Measuring convergence to equilibrium

- ▶ Mixing time : (according to a given metric).
Standard choice: L^1 (total-variation) mixing time to within distance ε is defined as

$$t_{\text{mix}}(\varepsilon) = \inf \left\{ t : \max_{x_0} \|p^t(x_0, \cdot) - \mu\|_{\text{tv}} \leq \varepsilon \right\}$$

(where $\|\mu - \nu\|_{\text{tv}} = \sup_{A \subset \Omega} [\mu(A) - \nu(A)]$)



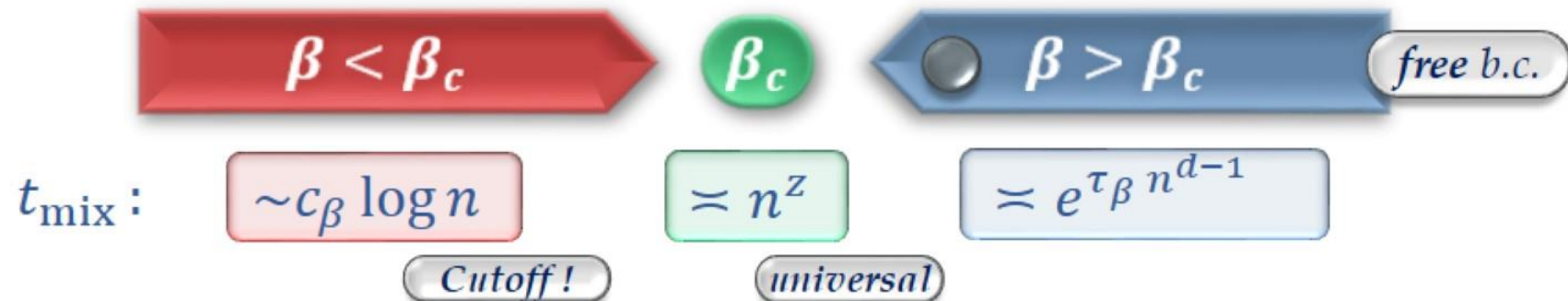
- ▶ Dependence on ε : (*cutoff phenomenon* [DS81], [A83],[AD86])

We say there is *cutoff* $\Leftrightarrow t_{\text{mix}}(\varepsilon) \sim t_{\text{mix}}(\varepsilon') \quad \forall$ fixed $\varepsilon, \varepsilon'$



Believed picture for Ising on \mathbb{Z}_n^d

- ▶ For some critical inverse-temperature β_c :



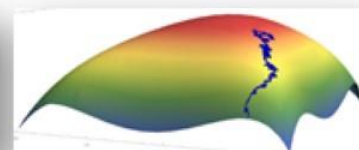
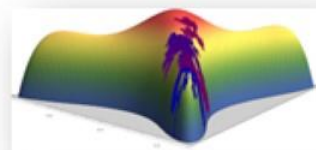
- ▶ Analogous picture verified for:

- ▶ Complete graph [Ding, L., Peres '09a, '09b], [Levin, Luczak, Peres '10] :

$\frac{1}{2(1-\beta)} \log n + O(1)$	$\asymp \sqrt{n}$	$\asymp \frac{1}{\beta-1} \exp\left[\frac{3}{4}(\beta-1)^2 n\right]$
--------------------------------------	-------------------	--

- ▶ Regular tree [Berger, Kenyon, Mossel, Peres '05] (high T /low T)
[Ding, L., Peres '10] (critical T)

- ▶ Potts model on complete graph
[Cuff, Ding, L., Loidor, Peres, Sly '12]



Glauber dynamics for 2D Ising

▶ Fast mixing at **high** temperatures:

- [Aizenman, Holley '84]
- [Dobrushin, Shlosman '87]
- [Holley, Stroock '87, '89]
- [Holley '91]
- [Stroock, Zegarlinski '92a, '92b, '92c]
- [Lu, Yau '93]
- [Martinelli, Olivieri '94a, '94b]
- [Martinelli, Olivieri, Schonmann '94]

$$\beta < \beta_c$$

$$t_{\text{mix}} \asymp \log n$$

▶ Slow mixing at **low** temperatures:

- [Schonmann '87]
- [Chayes, Chayes, Schonmann '87]
- [Martinelli '94]
- [Cesi, Guadagni, Martinelli, Schonmann '96]

$$\beta > \beta_c$$

free b.c.

$$t_{\text{mix}} = e^{(\tau\beta + o(1))n^{d-1}}$$

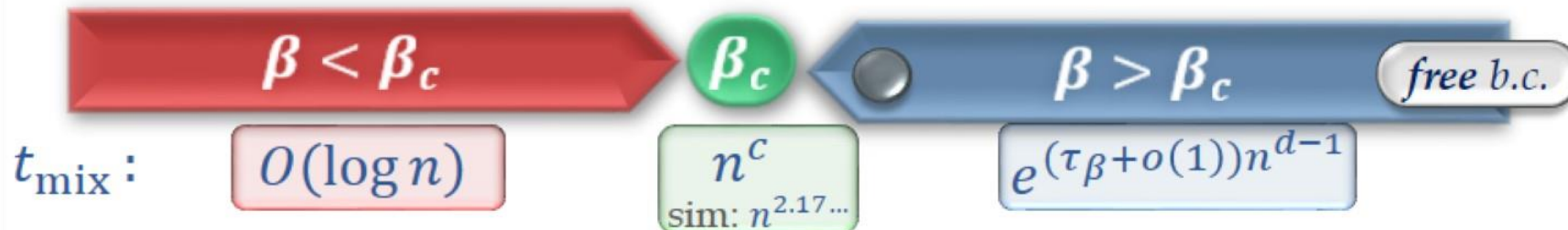
$$\beta_c$$

▶ **Critical** power-law:

- *simulations*: [Ito '93],[Wang, Hatano, Suzuki '95],[Grassberger '95],...: $n^{2.17\dots}$
- *lower bound*: [Aizenman, Holley '84], [Holley '91]
- *upper bound (polynomial mixing)*: [L., Sly '12]

$$n^{c_1} \leq t_{\text{mix}} \leq n^{c_2}$$

Glauber dynamics for 2D Ising

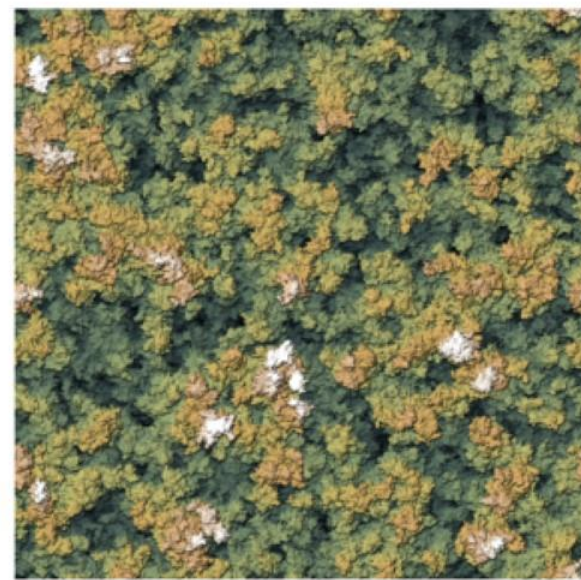


High temperature in 2D:

- [L., Sly '13]: *cutoff*
for any $\beta < \beta_c = \frac{1}{2} \log(1 + \sqrt{2})$:

$$t_{\text{mix}}(\varepsilon) = \frac{1}{2} \lambda_\infty^{-1} \log n + O(\log \log n)$$

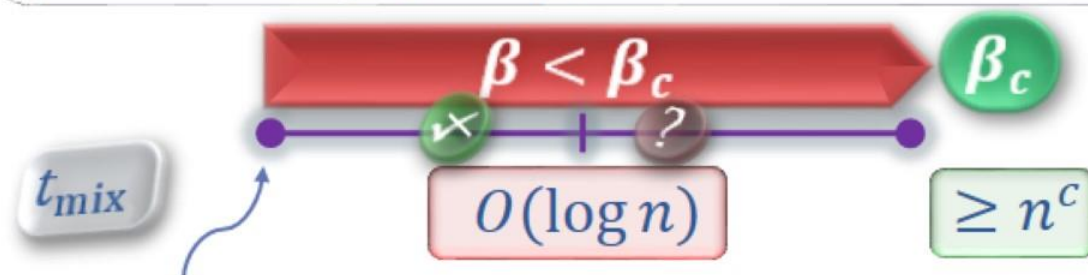
- Method caveat: needs *strong spatial mixing*; e.g., breaks on 3D Ising for β close to β_c .



High temperature unknowns (I)

▶ High temperature \leftrightarrow Infinite temperature:

Qualitatively, $\beta < \beta_c$ believed to behave \approx as $\beta = 0$.



[Martinelli, Olivieri '94],
[Aizenman, Holley '84]

▶ $\beta = 0$: (*independent spins*) one of the first examples of **cutoff**:

$$t_{\text{mix}}(\varepsilon) = c \log n + O(1)$$

[Aldous '83], [Diaconis Shahshahani '87]
[Diaconis, Graham, Morisson '90]

▶ \Rightarrow expect **cutoff** $\forall \beta < \beta_c$ (conj. [Peres '04]) & with $O(1)$ -window

▶ Concretely: for 3D Ising (*e.g.* on a torus) at $\beta = 0.99 \beta_c$:

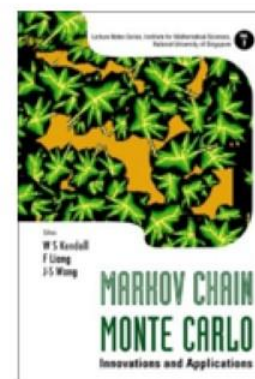
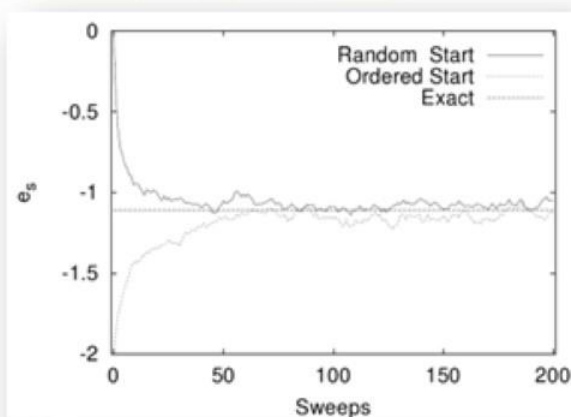
? ▶ does the dynamics exhibit cutoff?
if so, where & what is the window?

High temperature unknowns (II)

- ▶ Warm (random) start vs. cold (ordered) start:
random start is better than ordered



▶ e.g.



- ▶ Concretely: for 3D Ising at $\beta = 0.01$:

? ▶ what is $t_{\text{mix}}^{(U)}(\varepsilon) = \inf \left\{ t : \left\| \frac{1}{|\Omega|} \sum_{x_0} p^t(x_0, \cdot) - \mu \right\|_{\text{tv}} \leq \varepsilon \right\}$?
how does it compare with $t_{\text{mix}}(\varepsilon)$?

CFTP

High temperature unknowns (III)

▶ Universality of cutoff:

on any locally finite geometry there should be cutoff if the temperature is high enough (function of max-degree)

▶ $\exists c_0 > 0$: The Ising model on any graph G on n vertices with maximal degree d at $\beta < c_0/d$ has $t_{\text{mix}} = O(\log n)$

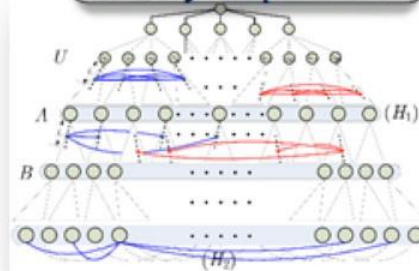
[Dobrushin '71],[Holley '72],[Dobrushin-Shlosman '85], [Aizenman-Holley '87]

▶ \Rightarrow expect **cutoff** $\forall \beta < \kappa/d$, and with $O(1)$ -window.

▶ Concretely: for Ising on a binary tree at $\beta = 0.01$:

? ▶ does the dynamics exhibit cutoff?
if so, where & what is the window?

or any expander

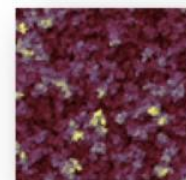


Recipe for stochastic Ising analysis

- ▶ Traditional approach to sharp mixing results
 1. Establish spatial properties of static Ising measure
 2. Use to drive a multi-scale analysis of dynamics.

▶ Example: best-known results on 2D Ising (torus \mathbb{Z}_n^2):

- ▶ [L., Sly '13]: *cutoff* at $\beta < \beta_c = \frac{1}{2} \log(1 + \sqrt{2})$
 - used *log-Sobolev ineq.* & *strong spatial mixing*.



- ▶ [L., Sly '12]: *power-law* at β_c
 - used *SLE behavior* of critical interfaces.



- ▶ [L., Martinelli, Sly, Toninelli '13]: at $\beta > \beta_c$
quasi-polynomial mixing under all-plus b.c.
 - uses interface convergence to *Brownian bridges*

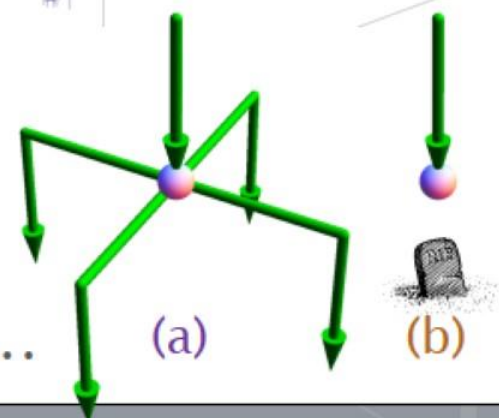


New framework for the analysis

- ▶ Traditional approach to sharp mixing results
 1. Establish spatial properties of static Ising measure
 2. Use to drive a multi-scale analysis of dynamics.
- ▶ New approach: study these *simultaneously* examining *information percolation* clusters in the space-time slab:



- ▶ track update lineage back in time.
- ▶ update either (a) branches out, or (b) terminates (“*oblivious*”)
- ▶ analyze RED/GREEN/BLUE clusters...



Results: cutoff up to β_c in 3D Ising

- ▶ Confirm Peres's conj. on \mathbb{Z}_n^d for any d , with $O(1)$ -window.
- ▶ **THEOREM:** ([L.-Sly '14+])

$\forall d \geq 1$ and $\beta < \beta_c$ there is **cutoff** with an $O(1)$ -window at

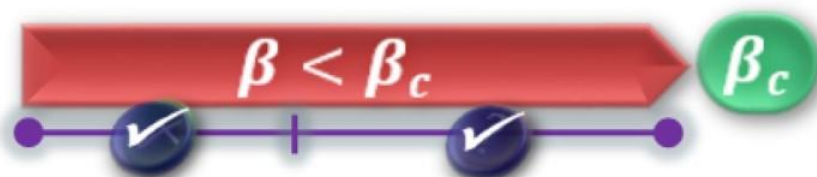
$$t_m = \inf \left\{ t : \mathbb{E}_+ [M(\sigma_t)] \leq \sqrt{n^d} \right\}$$

*cutoff window:
 $O(\log(1/\varepsilon))$*

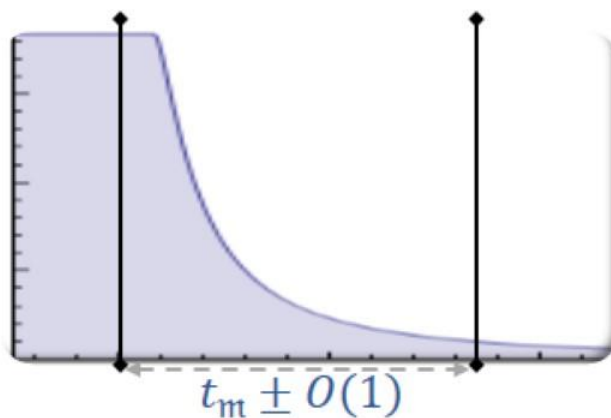
- ▶ Examples:

- ▶ $d = 1$: $t_m = \frac{1}{2(1-\tanh(2\beta))} \log n$.

- ▶ $\beta = 0$: $t_m = \frac{1}{2} \log n$ (matching [Aldous '83])



[recall $M(\sigma) = \frac{1}{|\Lambda|} \sum \sigma(x)$]



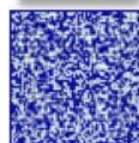
Results: initial states

▶ *Warm start is twice faster:*

▶ **All-plus starting state is worst** (up to an additive $O(1)$)
[but twice faster than naïve monotone coupling bound].

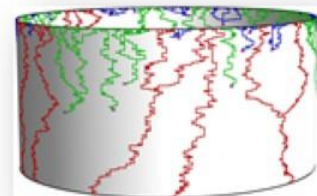


▶ Uniform initial state \approx **twice faster than all-plus.**



▶ Almost \forall deterministic initial state \approx **as bad as all-plus.**

▶ Example: the 1D Ising model (\mathbb{Z}_n):



THEOREM: ([L.-Sly '14+])

Fix $\beta > 0$ and $0 < \varepsilon < 1$; set $t_m = \frac{1}{2(1-\tanh(2\beta))} \log n$.

1. (*Annealed*) $t_{\text{mix}}^{(U)}(\varepsilon) \sim \frac{1}{2} t_m$

2. (*Quenched*) $t_{\text{mix}}^{(x_0)}(\varepsilon) \sim t_{\text{mix}}^{(+)}(\varepsilon) \sim t_m$ for almost $\forall x_0$

Results: universality of cutoff

- ▶ **Paradigm:** cutoff for *any locally finite geometry* at high enough temperature (including expanders, trees, ...)

$$\beta < \kappa/d \quad \text{Cutoff}$$

- ▶ THEOREM: ([L.-Sly '14+])

$\exists \kappa > 0$ so that, if G is any n -vertex graph with degrees $\leq d$ and $\beta < \kappa/d$, then \exists cutoff with an $O(1)$ -window at

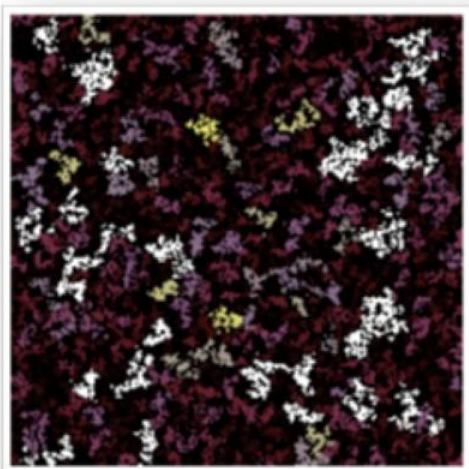
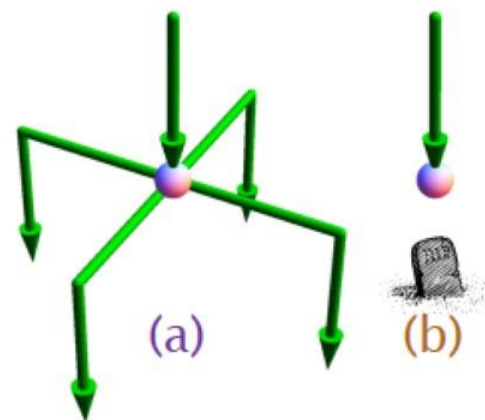
$$t_m = \inf \left\{ t : \sum_x \mathbb{E}_+ \left[M(\sigma_t(x))^2 \right] \leq 1 \right\}.$$

- ▶ Moreover:

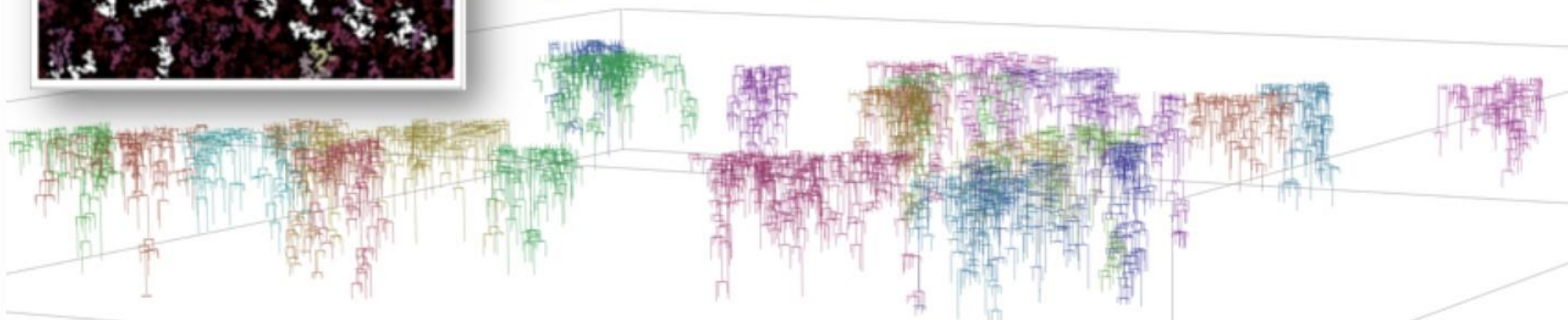
$$t_{\text{mix}}^{(U)} \leq \left(\frac{1}{2} + \varepsilon_\beta \right) t_m \quad \text{yet} \quad t_{\text{mix}}^{(x_0)} \geq (1 - \varepsilon_\beta) t_m \quad \text{a.e. } x_0.$$

The new framework (revisited)

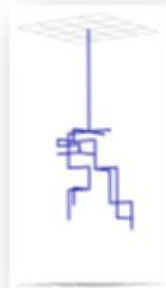
- ▶ *Information percolation* clusters in the space-time slab:
 - ▶ track update lineage back in time.
 - ▶ update either (a) branches out, or (b) terminates (“*oblivious*”)



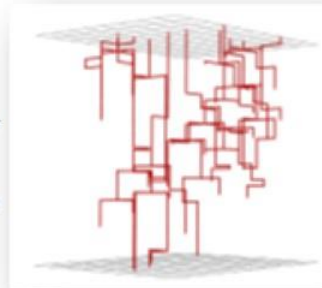
\mathbb{Z}_{200}^2 cluster
(top/side view)



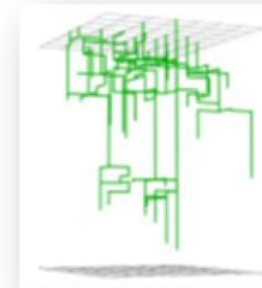
Information percolation clusters



BLUE:
dies out quickly
in space & time.

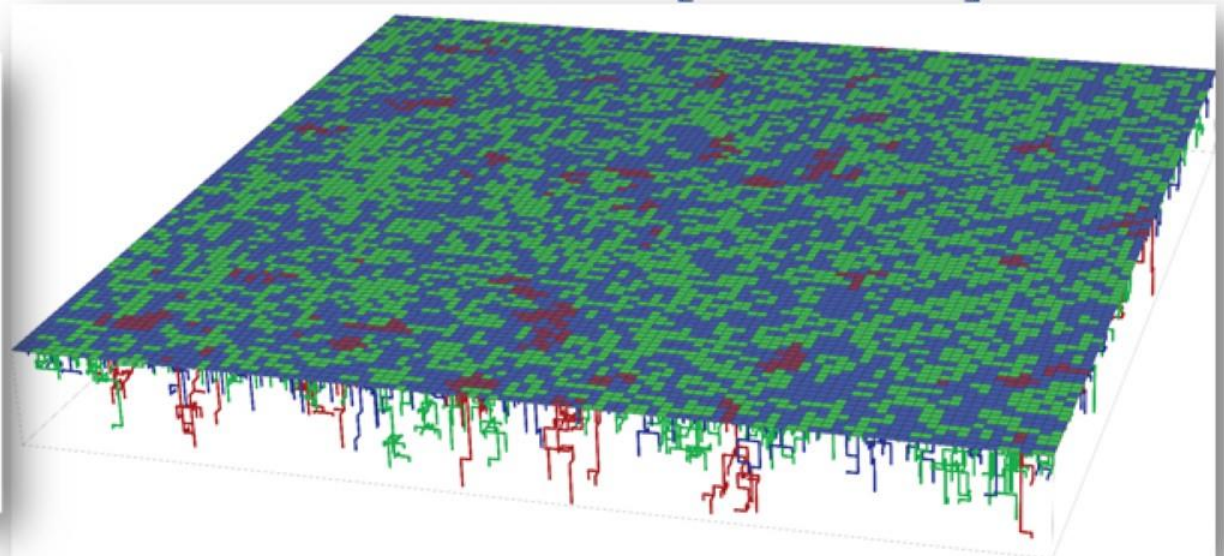
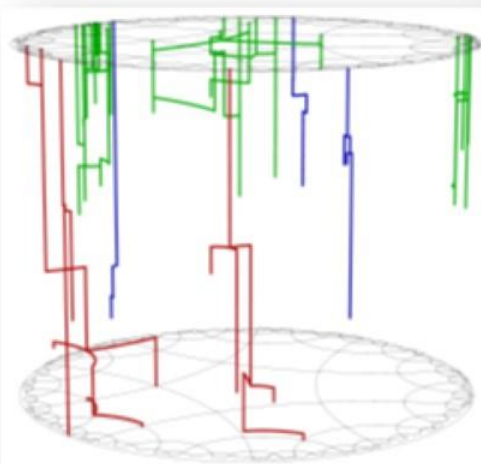


RED:
top spins are
affected by
initial state.



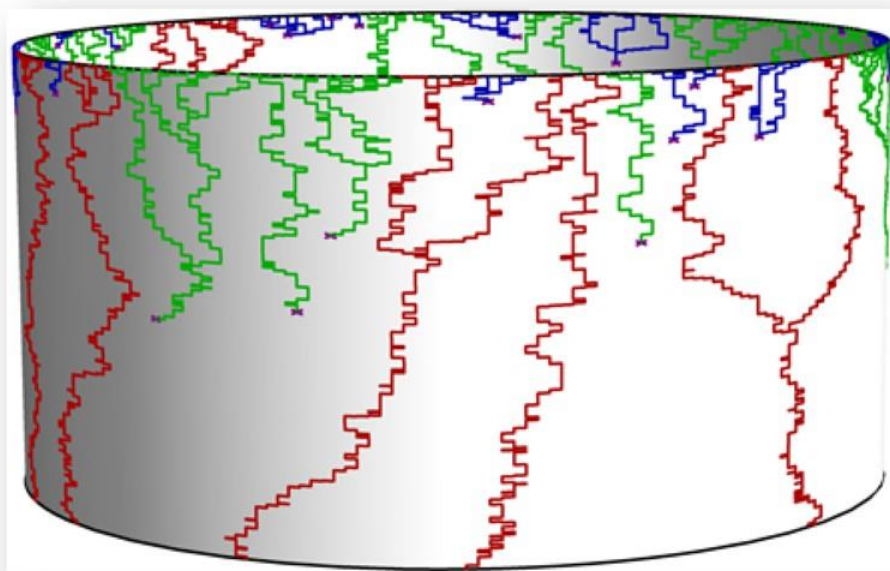
GREEN:
o/w.

- ▶ Rough idea: condition on **GREEN**, let the effect of **RED** clusters vanish among **BLUE** (show $\mathbb{E} \left[2^{|R \cap R'|} \mid G \right] \rightarrow 1$).



Example: the framework in 1D

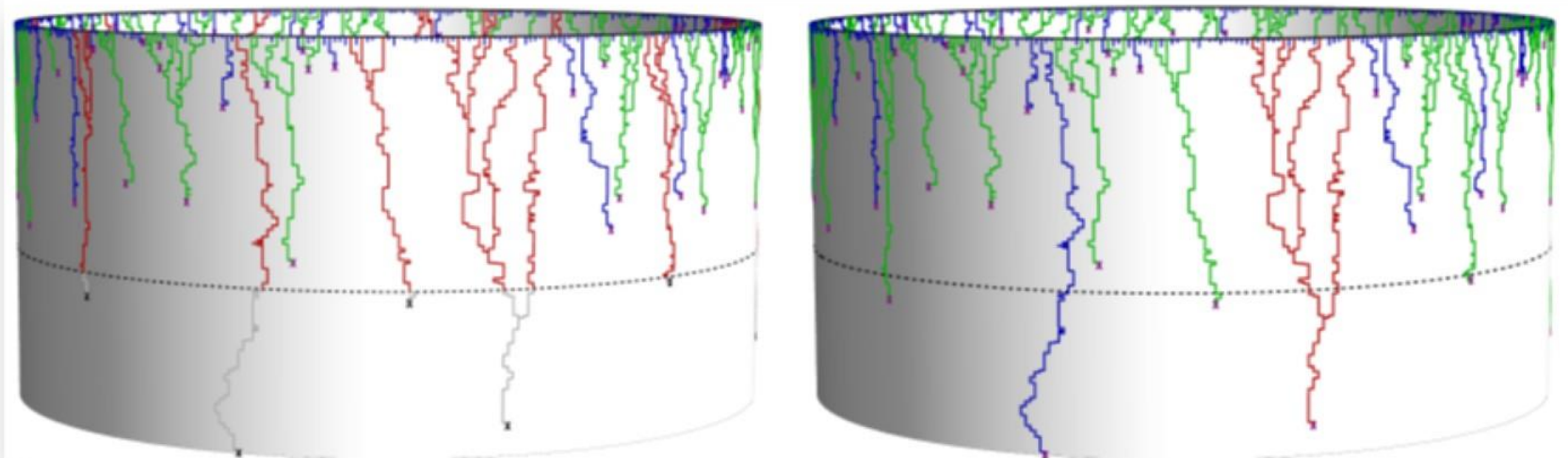
- ▶ In 1D: $\theta = \mathbb{P}(\text{oblivious update}) = 1 - \tanh 2\beta$
- ▶ Update history: continuous-time RW killed at rate θ .
- ▶ $\mathbb{P}(\text{surviving to time } t_m) \text{ is } \approx 1/\sqrt{n}$.
- ▶ Cutoff at $t_m = \frac{1}{2\theta} \log n$
- ▶ Effect of the initial state on the final state is in terms of the bias of the cont.-time RW...



the 3 cluster classes (R/G/B) in \mathbb{Z}_{256}

Example: random initial state

- ▶ Handling a uniform (IID) starting configuration:
 - Compare the dynamics directly with Ising measure: develop history to time $-\infty$ (*coupling from the past*).
 - Redefine **RED** clusters (coalesce before time 0).



Losing red clusters in a blue sea

▶ LEMMA: ([Miller, Peres '12])

Let μ be a measure on $\sigma \in \Omega = \{\pm 1\}^n$ as follows:

1. draw a random variable $R \subseteq [n]$ via a law $\tilde{\mu}$;
 2. let $\sigma_R \sim \text{law } \phi_R$ and $\sigma_{[n] \setminus R} \sim \text{IID Bernoulli } \begin{cases} +1 & 1/2 \\ -1 & 1/2 \end{cases}$
- $$\Rightarrow \|\mu - \nu\|_{L^2(\nu)}^2 \leq \mathbb{E} \left[2^{|R \cap R'|} \right] - 1$$

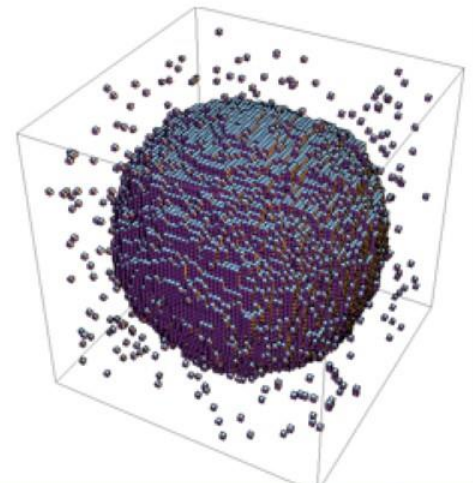
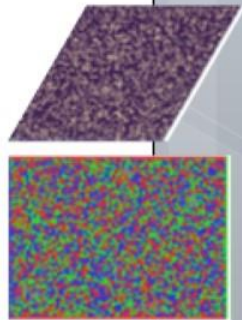
$\nu = \text{uniform measure}$

$R, R' \text{ IID}$

- ▶ the set R embodies the nontrivial part of μ
- ▶ has a negligible effect provided the exponential moment can be controlled...

Open problems

- ▶ High temperature regime for other spin-systems (Potts / Independent sets / Legal colorings / Spin glass,...):
 - ▶ asymptotic mixing on the lattice up to β_c
 - ▶ cutoff on a transitive **expander**
 - ▶ asymptotic mixing from random starting states (e.g., compare ordered/disordered start in Potts)
- ▶ 3D Ising:
 - ▶ no cutoff at criticality
 - ▶ power-law behavior at criticality
 - ▶ sub-exponential upper bound at low temperatures under all-plus b.c.



Thank you

