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NODAL DOMAINS FOR MAASS FORMS

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\mathbb{H} UPPER HALF PLANE

$$\Gamma = \mathrm{SL}_2(\mathbb{Z})$$

$$X = \Gamma \backslash \mathbb{H} \quad \text{MODULAR SURFACE}$$

THE EVERYWHERE UNRAMIFIED
AUTOMORPHIC CUSP FORMS FOR GL_2/\mathbb{Q}
ARE THE EVEN HECKE/MAASS FORMS
ON X :

$$\Delta \phi + \lambda \phi = 0, \quad \phi \in L^2_{\mathrm{cusp}}(X)$$

$$\phi(\sigma z) = \phi(z), \quad \sigma: X \rightarrow X$$

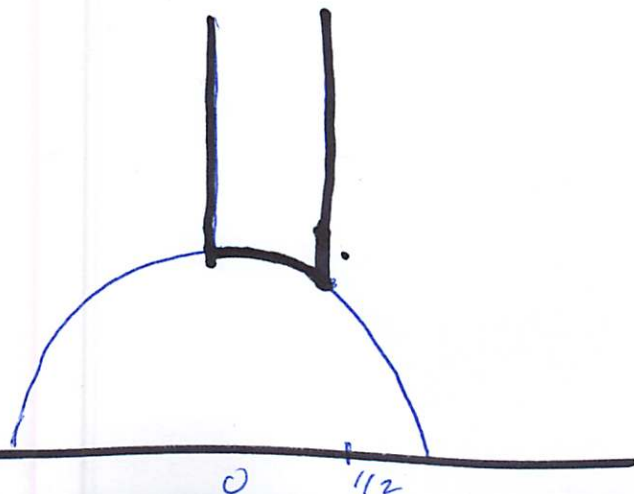
THE INVOLUTION INDUCED BY

$$z \rightarrow -\bar{z}$$

$$\mathrm{Fix}(\sigma) = \mathcal{S} =$$

$$\partial_n \phi|_{\mathcal{S}} = 0$$

/NEUMANN COND.



[3]

OUR INTEREST IS IN THE NODAL DOMAINS OF ϕ .

$$Z(\phi) = \{z \in X : \phi(z) = 0\}, \text{ NODAL LINE}$$

$$X \setminus Z(\phi) = \bigsqcup_{j=1}^N \Omega_j, \quad \Omega_j \text{ CONNECTED}$$

COMPONENTS: NODAL DOMAINS.

HOW MANY ARE THERE?

THE NUMBER MEETING A LARGE BUT FIXED COMPACT $K \subset X$ SATISFIES

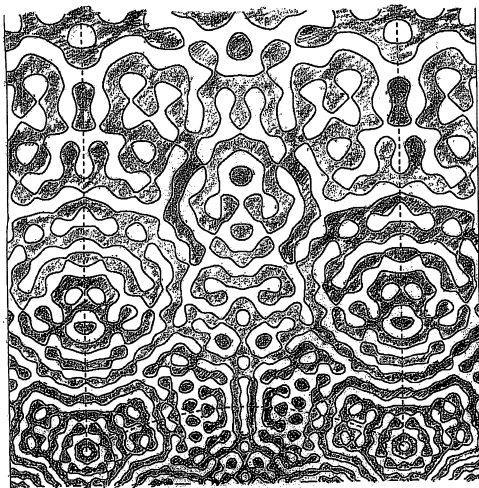
$$N_K(\phi_\lambda) \ll \lambda \quad (\text{COURANT}).$$

[FOR COMPACT MANIFOLDS IN ANY DIMENSION

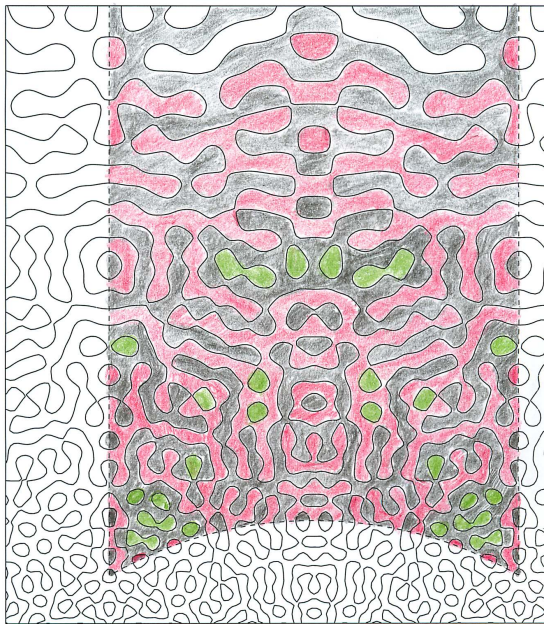
$$N(\phi_n) \leq n$$

WHERE n DENOTES THE n -TH EIGENFUNCTION]

Below a picture of the zero set $\phi = 0$ of such a “Maass form” for $SL_2(\mathbb{Z})$, $\lambda = \frac{1}{4} + t^2$, $t = 125.34 \dots$ (Hejhal–Rackner).
Is the zero set behaving randomly? How many components does it have?

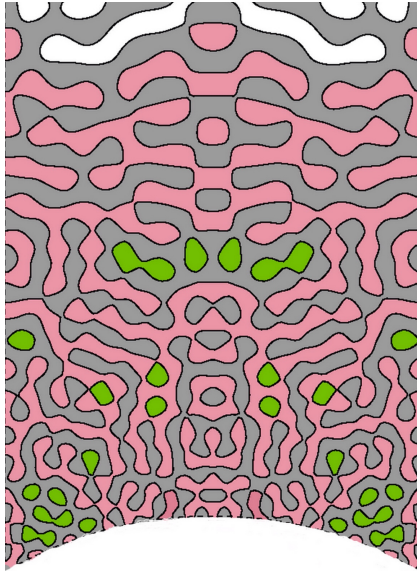


58 nodal domains in A

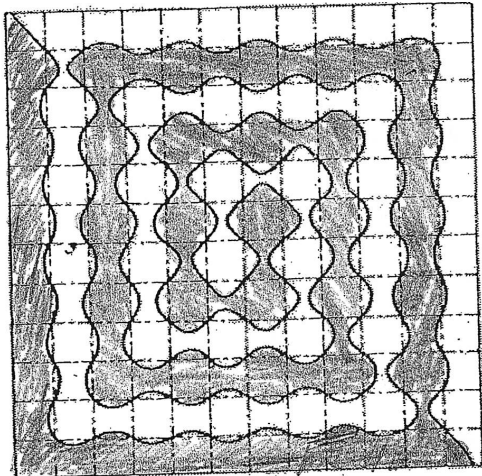


Hejhal–Rackner nodal lines for $\lambda = 1/4 + R^2$, $R = 125.313840$

Nodal portrait



The trouble is giving a lower bound for $N(\phi)$. In general it need not grow!



Nodal domain of an eigenfunction on the square, $N(\phi) = 2$.
From Courant–Hilbert; Vol I. Thesis A. Stern, Gottingen, 1925.

WHAT SHOULD WE EXPECT?

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M. BERRY: THE EIGENSTATES OF THE QUANTIZATION OF A CLASSICALLY CHAOTIC HAMILTONIAN IN THE SEMI-CLASSICAL LIMIT BEHAVE LIKE RANDOM MONOCHROMATIC WAVES.

• THE CLASSICAL GEODESIC MOTION ON $T_1^*(X) \cong \mathbb{P}^1 \backslash \text{SL}_2(\mathbb{R})$.

IS WELL KNOWN TO BE CHAOTIC.

RANDOM MONOCHROMATIC WAVES:

S^2 (SIMILAR DEFN IN ANY DIMENSION)

H_k : THE $2k+1$ DIMENSIONAL SPACE OF SPHERICAL HARMONICS OF DEGREE k . THESE YIELD THE EIGENFUNCTIONS OF Δ_{S^2} WITH EIGENVALUE $\lambda_k = \frac{k(k+1)}{2}$.

ϕ_j , $j = -k, \dots, k$ O.N.B OF H_k 15

$$\psi_\lambda(z) = \sum_{j=-k}^k a_j \phi_j(z)$$

WITH a_j I.I.D. $N(0,1)$ VARIABLES.

ψ_λ IS A RANDOM WAVE WITH FREQ. λ .
WHAT DO ITS NODAL DOMAINS LOOK LIKE?

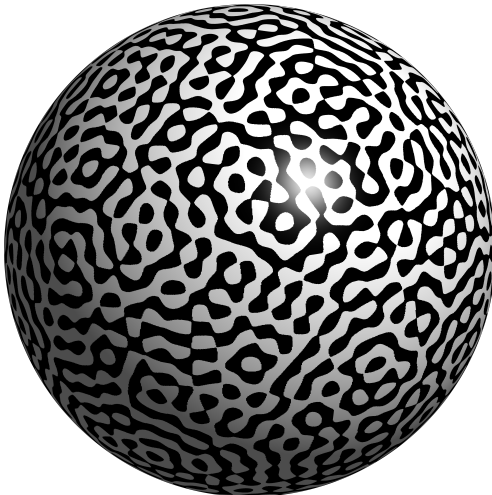
THEOREM 1 (NAZAROV-SODIN 2008)

AS $\lambda \rightarrow \infty$ THE RANDOM ψ_λ HAS

$$N(\psi_\lambda) \sim c\lambda \quad \text{FOR SOME UNIVERSAL } c > 0.$$

FOR $m \geq 1$ LET $N_m(\psi_\lambda)$ BE THE
NUMBER OF NODAL DOMAINS OF ψ
WHICH HAVE CONNECTIVITY m .

Nodal portrait: Random spherical harmonic ($\alpha = 1$)



random spherical harmonic of degree = 80. (A. Barnett)

THEOREM 2: (CANZANI-WIGMAN-S 2017) 16

THERE IS A PROBABILITY MEASURE μ ON \mathbb{N} WITH $\mu(m) > 0$ FOR ALL m S.T. FOR RANDOM ψ

$$N_m(\psi_\lambda) \sim \mu(m) N(\psi_\lambda) \text{ AS } \lambda \rightarrow \infty.$$

NUMERICAL MONTE CARLO EXPERIMENTS
(NASTASESCU, KONRAD, BARNETT, JIN, ...)
YIELD VALUES ;

$$C \approx 1/20 \quad \left(\begin{array}{l} \text{RIGOROUS BOUND} \\ C \geq 10^{-70} \\ \text{DE COURCY-IRELAND} \end{array} \right)$$

$$\begin{aligned} \mu(1) &= 0.906, \quad \mu(2) = 0.055, \\ \mu(3) &= 0.010, \quad \mu(4) = 0.006, \dots \end{aligned}$$

CHECKING THESE NUMBERS AGAINST
THE NODAL DOMAINS OF MAASS
FORMS ON $SL_2(\mathbb{Z})$ GIVES GOOD
AGREEMENT.

• SO WE EXPECT AN INDIVIDUAL WEYL
TYPE LAW FOR THE NODAL COUNT.

REMARK: THE THEOREMS FOR THE
 NODAL SETS AND DOMAINS OF RANDOM
 MONOCHROMATIC WAVES CAN BE CARRIED
 OUT IN ANY DIMENSION YIELDING A
 UNIVERSAL DISTRIBUTION FOR THE
 TOPOLOGIES OF $Z(\psi)$ (C-W-S).
 HOWEVER THE CORRESPONDING CONSTANTS
 IN DIMENSION 3 AND HIGHER APPEAR
 TO BE ABSURDLY SMALL. WE APPEAR
 TO BE IN A SUPERCRITICAL REGIME
 (NUMERICAL EXPERIMENTS - BARNETT)
 SO THAT FOR PRACTICAL PURPOSES
 THERE ESSENTIALLY TWO NODAL
 DOMAINS!

WHAT CAN ONE PROVE FOR
 MAASS FORMS AND HOW
 TO PRODUCE NODAL DOMAINS?

THEOREM 3 (GHOSH-REZNIKOV-S 2014):

$\Gamma = \mathrm{SL}_2(\mathbb{Z})$, ϕ AS ABOVE, K
A FIXED LARGE COMPACT IN X ,
ASSUME THE LINDELOFF HYPOTHESIS
FOR $L(s, \phi_\lambda)$ THEN

$$N(\phi_\lambda) \gg \lambda^{1/24}.$$

MORE GENERALLY LET X BE A
CONGRUENCE ARITHMETIC SURFACE
(I.E. A CONGRUENCE COVER OF THE
MODULAR SURFACE OR A SHIMURA "CURVE")
ASSUME THAT X CARRIES A REFLECTION
SYMMETRY $\sigma: X \rightarrow X$.
 $S = \mathrm{FIX}(\sigma) \subset X$ IS
A FINITE UNION OF CLOSED GEODESICS.
WE ASSUME THAT ALL ~~MASS~~ MASS
EIGENFUNCTION ϕ ARE ALSO
EIGENFUNCTIONS OF σ .

THEOREM 4 (G-R-S 2017):

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ASSUMING THE LINDELOFF HYPOTHESIS
FOR GL_2 AUTOMORPHIC L-FUNCTIONS
(OVER ANY F) WE HAVE THAT

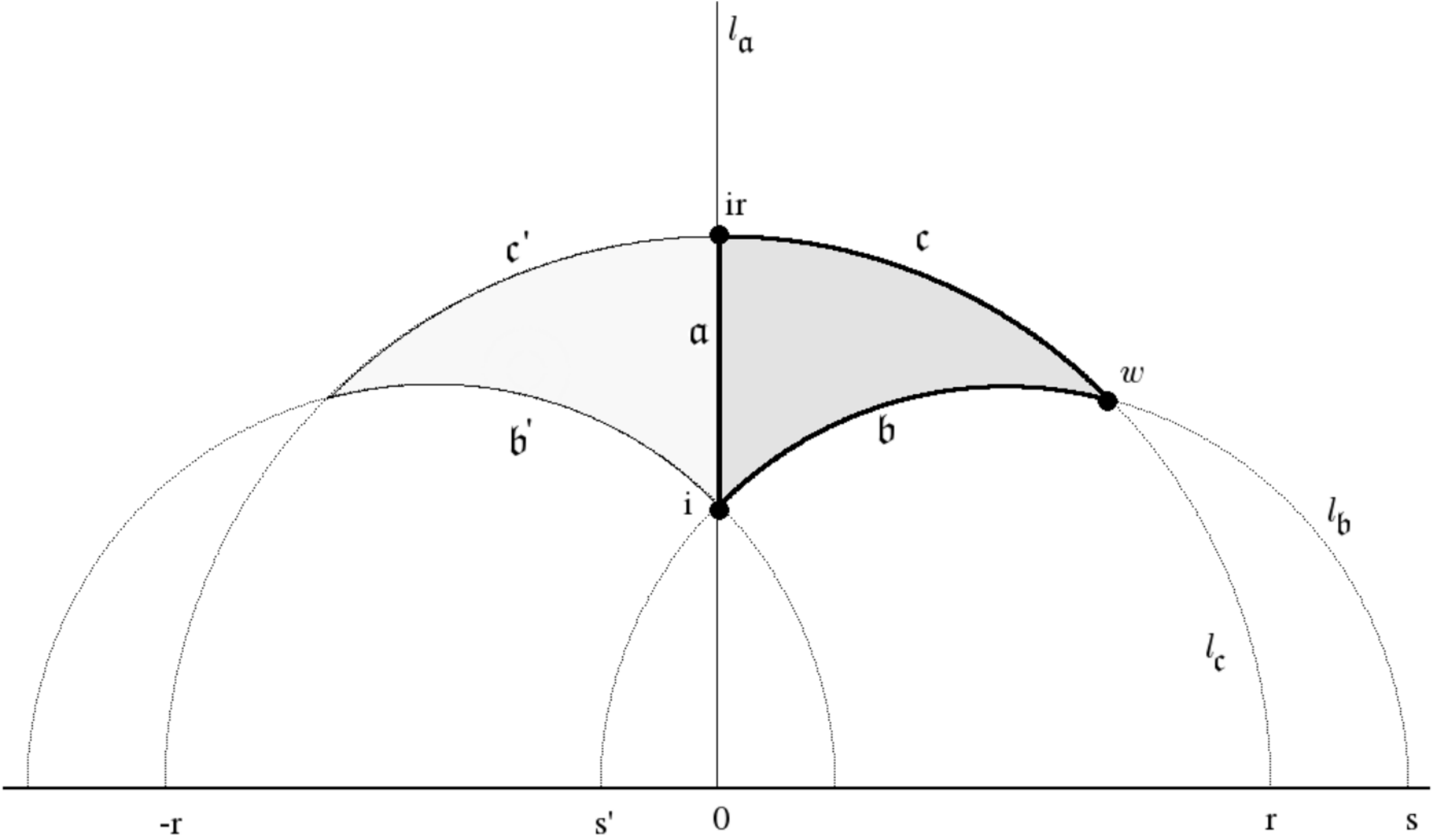
$$N(\phi_\lambda) \gg \lambda^{1/27}.$$

RECENTLY J. JANG AND J. JUNG
PROVED AN UNCONDITIONAL RESULT
IN THE ABOVE COMPACT X SETTING.

THEOREM 5 (JANG-JUNG 2018):

LET X BE A COMPACT ARITHMETIC
SURFACE AS ABOVE AND ϕ_λ ALSO
AS ABOVE THEN $N(\phi_\lambda) \rightarrow \infty$ AS $\lambda \rightarrow \infty$.

THEOREMS 4 AND 5 APPLY TO THE
COMPACT ARITHMETIC TRIANGLES
WITH DIRICHLET OR NEUMANN BDRY
CONDITIONS. EG $(\pi/2, \pi/6, \pi/6)$



THE PROOFS MAKE USE OF VARIOUS RESULTS IN THE SEMI-CLASSICS OF MAASS FORMS "QUANTUM CHAOS" AS WELL AS PERIOD FORMULAE.

(1) QUE: QUANTUM UNIQUE ERGODICITY (LINDENSTRAUSS, SOUNDARARAJAN)

$$\phi_\lambda \text{ ON } X_\Gamma, a \in C(\Gamma \backslash SL_2(\mathbb{R})) \simeq T_1^*(X_\Gamma)$$

$$\|\phi_\lambda\|_2 = 1, \text{ THEN}$$

$$\langle \mathcal{O}_\Gamma(a) \phi_\lambda, \phi_\lambda \rangle \rightarrow \int_{\Gamma \backslash SL_2(\mathbb{R})} a(g) dg \text{ AS } \lambda \rightarrow \infty.$$

(2) SUBCONVEX BOUND FOR L^∞ :

$$\|\phi_\lambda\|_\infty \ll \lambda^{5/24} \cdot \|\phi_\lambda\|_2 \quad \begin{matrix} \text{(IWANIEC)} \\ \Gamma/\mathbb{Q} \end{matrix}$$

$$\|\phi_\lambda\|_\infty \ll_{X_\Gamma} \lambda^{1/4 - \frac{1}{27}}, \quad \begin{matrix} (\Gamma/F) \\ \text{(BLUMER-MICHEL)} \end{matrix}$$

III

(3) WALDSPURGER'S FORMULA

$$\left| \int_{\mathcal{S}} \phi(s) w(s) ds \right|^2 = * L\left(\frac{1}{2}, \pi \times w\right)$$

↑ explicit
and can be
estimated

w A CHARACTER
OF THE TORUS DEFINED BY \mathcal{S} ,
 π BASE CHANGE OF ϕ TO E THE
CORRESPONDING QUADRATIC EXTN OF F .

• HOW TO PRODUCE NODAL DOMAINS?
CAN WE LOCALIZE?

NOTE THAT IN STERN'S EXAMPLE
 $\phi|_{\partial\Omega}$ HAS ONLY TWO SIGN CHANGES.

TOPOLOGICAL REDUCTION TO \mathcal{S} .

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TOPOLOGICAL LEMMA (GENERAL SURFACE)

X, δ, σ AS ABOVE

$$\partial_n \phi|_{\delta} = 0 \quad (\text{IE } \phi \text{ IS } \sigma \text{ EVEN})$$

LET $\eta(\phi)$ BE THE NUMBER OF SIGN CHANGES OF ϕ AS ONE TRAVERSES δ , THEN

$$N(\phi) \geq \frac{\eta(\phi)}{2} + 1 - \text{genus}(X)$$



IF $z(\phi)$ HAS NO SELF CROSSINGS THE

BOUND IS EASY TO PROVE BY INDUCTION ; THE POINT WHICH WORKS IN OUR FAVOR HERE IS THAT SELF CROSSINGS ONLY INCREASE THE NUMBER OF NODAL DOMAINS.

[13]

- THIS LOCALIZES OUR PROBLEM TO PRODUCING SIGN CHANGES OF ϕ ALONG δ .
- THE NODAL DOMAINS THAT WE PRODUCE IN THIS WAY ARE ALL " σ -INERT", THAT IS THEY ARE INVARIANT UNDER σ OR WHAT IS THE SAME THING THEY CROSS δ . THEIR NUMBER IN THIS REAL ANALYTIC SETTING CAN BE BOUNDED FROM ABOVE (TOTH-ZELDITCH 2008)

$$N_{\text{inert}}(\phi_\lambda) \ll \lambda^{1/2}.$$

SO WITH THIS METHOD WE WILL NEVER PRODUCE THE EXPECTED $c\lambda$ DOMAINS MOST OF WHICH ARE SPLIT.

L^2 -RESTRICTIONS TO δ (G-R-S)

IF X IS COMPACT (AND IN THE NON-COMPACT CASE SOME ARITHMETIC INPUT IS USED (AND δ BELOW IS REPLACED BY A LARGE COMPACT SUBSEGMENT) THEN QUE IMPLIES THAT OUR ϕ_λ 'S SATISFY

$$\int_{\delta} |\phi_\lambda(s)|^2 ds \gg 1 \quad (\|\phi_\lambda\|_2 = 1)$$

(IF ϕ_λ IS NOT EVEN ABOUT δ THEN THIS IS REPLACED BY A COMBINATION OF THE L^2 RESTRICTION OF ϕ_λ TO δ PLUS THE L^2 -RESTRICTION OF $\partial_n \phi$ (NORMALIZED PROPERLY).

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THE IDEA TO PRODUCE SIGN CHANGES
IS TO COMPARE

$$\left| \int_C \phi ds \right| \quad \text{WITH} \quad \int_C |\phi| ds$$

FOR SUITABLE SEGMENTS C OF δ AND
TO SHOW THESE AREN'T EQUAL.

WE HAVE:

$$1 \ll \int_{\delta} |\phi(s)|^2 ds \leq \|\phi\|_{\infty} \int_{\delta} |\phi(s)| ds$$

$$= \|\phi\|_{\infty} \sum_{j=1}^n \left| \int_{C_j} \phi_{\lambda}(s) ds \right|$$

→ segments
of fixed
sign.

Waldpinger's formula
+ Lindeloff

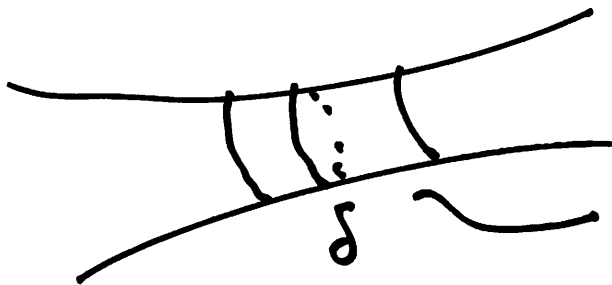
$$\Rightarrow \int_{\epsilon} \phi_{\lambda}(s) ds \ll_{\epsilon} \lambda^{-1/4 + \epsilon}.$$

[16]

SO THE SUBCONVEX BOUND FOR L^∞ IS PRECISELY WHAT IS NEEDED TO OBTAIN A (POWER) LOWER BOUND FOR η .

JANG-JUNG:

THEY START BY EXPLICATING IN THIS CASE AN INTEGRATION BY PARTS ARGUMENT OF CHRISTIANSON - TOTH - ZELDITCH



CONSTRUCT A CUTOFF
Y.d.o.p. IN A NBH
OF δ

AND USE THE
QUE HYPOTHESIS

LEADS TO:

[17]

IF $f(s) \geq 0$ ON δ IS FIXED, $m \geq 0$
THEN

$$\lim_{\lambda \rightarrow \infty} \left[\int_{\delta} |f(s) \phi_{\lambda}(s)|^2 ds - \lambda^{-m} \int_{\delta} |\partial_s^m (f(s) \phi_{\lambda}(s))|^2 ds \right] \\ = 2 \left(1 - \frac{1}{\pi} \int_{-1}^1 \frac{\bar{z}^{2m} dz}{\sqrt{1-z^2}} \right) \int_{\delta} f^2(s) ds$$

(*)
→

$$\Rightarrow \lim_{\lambda \rightarrow \infty} \int_{\delta} f(s)^2 \phi_{\lambda}^2(s) ds \geq 2 \int_{\delta} f(s)^2 ds$$

WHICH IS A STRONGER AND
LOCALIZED IMPROVEMENT OF THE
 L^2 -RESTRICTION BOUND ABOVE.

(18)

THERE IS NOTHING TO STOP

$$\|\phi_\lambda|_\delta\|_2 \rightarrow \infty \text{ As } \lambda \rightarrow \infty.$$

THEY RENORMALIZE

$$\frac{f(s)\phi_\lambda(s)}{\|\phi_\lambda|_\delta\|_2}$$

AND STUDY THE POSSIBLE WEAK
LIMITS OF THIS SEQUENCE ON \mathcal{S} .
THE FOURIER TRANSFORM $\hat{\phi}_\lambda$ (ON \mathcal{S})

$$\frac{\lambda \hat{\phi}_\lambda(\lambda \xi)}{\int_{\mathcal{S}} |\phi_\lambda|^2}$$

ARE POSITIVE DEFINITE IF $\phi_\lambda(s)f(s)$
IS A FIXED SIGN, THEY SHOW THAT
THEY HAVE ENOUGH INFORMATION FROM CA
TO SHOW THIS CANT HAPPEN!