

Liouville Equations and Functional Determinants

Andrea Malchiodi (SNS, Pisa)

IAS, March 5th, 2019

The determinant of the Laplacian

The determinant of the Laplacian

Consider a compact, closed manifold M with metric g , and Laplace-Beltrami operator Δ_g .

The determinant of the Laplacian

Consider a compact, closed manifold M with metric g , and Laplace-Beltrami operator Δ_g . The eigenvalues $\{\lambda_j\}$, with eigenfunctions $\{\varphi_j\}_j$

$$-\Delta_g \varphi_j = \lambda_j \varphi_j \quad \text{on } M,$$

The determinant of the Laplacian

Consider a compact, closed manifold M with metric g , and Laplace-Beltrami operator Δ_g . The eigenvalues $\{\lambda_j\}$, with eigenfunctions $\{\varphi_j\}_j$

$$-\Delta_g \varphi_j = \lambda_j \varphi_j \quad \text{on } M,$$

satisfy

$$\lambda_j \rightarrow +\infty \quad \text{as } j \rightarrow +\infty.$$

The determinant of the Laplacian

Consider a compact, closed manifold M with metric g , and Laplace-Beltrami operator Δ_g . The eigenvalues $\{\lambda_j\}$, with eigenfunctions $\{\varphi_j\}_j$

$$-\Delta_g \varphi_j = \lambda_j \varphi_j \quad \text{on } M,$$

satisfy

$$\lambda_j \rightarrow +\infty \quad \text{as } j \rightarrow +\infty.$$

Formally, the *determinant* of $-\Delta_g$ is defined as

$$\det(-\Delta_g) = \prod_j \lambda_j.$$

The determinant of the Laplacian

Consider a compact, closed manifold M with metric g , and Laplace-Beltrami operator Δ_g . The eigenvalues $\{\lambda_j\}$, with eigenfunctions $\{\varphi_j\}_j$

$$-\Delta_g \varphi_j = \lambda_j \varphi_j \quad \text{on } M,$$

satisfy

$$\lambda_j \rightarrow +\infty \quad \text{as } j \rightarrow +\infty.$$

Formally, the *determinant* of $-\Delta_g$ is defined as

$$\det(-\Delta_g) = \prod_j \lambda_j.$$

- While physicists may like these formulas, mathematicians usually have problems with infinite products of diverging numbers.

Regularized determinant

Regularized determinant

The *spectral zeta function* of (M^n, g) is

$$(1) \quad \zeta(s) = \sum_{j=1}^{\infty} \lambda_j^{-s}.$$

Regularized determinant

The *spectral zeta function* of (M^n, g) is

$$(1) \quad \zeta(s) = \sum_{j=1}^{\infty} \lambda_j^{-s}.$$

By Weyl's asymptotic law,

$$\lambda_j \sim j^{2/n}, \quad j \rightarrow \infty.$$

Regularized determinant

The *spectral zeta function* of (M^n, g) is

$$(1) \quad \zeta(s) = \sum_{j=1}^{\infty} \lambda_j^{-s}.$$

By Weyl's asymptotic law,

$$\lambda_j \sim j^{2/n}, \quad j \rightarrow \infty.$$

Consequently, (1) defines an analytic function provided $\operatorname{Re}(s) > n/2$.

Regularized determinant

The *spectral zeta function* of (M^n, g) is

$$(1) \quad \zeta(s) = \sum_{j=1}^{\infty} \lambda_j^{-s}.$$

By Weyl's asymptotic law,

$$\lambda_j \sim j^{2/n}, \quad j \rightarrow \infty.$$

Consequently, (1) defines an analytic function provided $\operatorname{Re}(s) > n/2$. Differentiating in s one finds

$$\zeta'(s) = \frac{d}{ds} \sum_{j=1}^{\infty} e^{-s \log \lambda_j} = - \sum_{j=1}^{\infty} \log \lambda_j e^{-s \log \lambda_j}.$$

Regularized determinant

The *spectral zeta function* of (M^n, g) is

$$(1) \quad \zeta(s) = \sum_{j=1}^{\infty} \lambda_j^{-s}.$$

By Weyl's asymptotic law,

$$\lambda_j \sim j^{2/n}, \quad j \rightarrow \infty.$$

Consequently, (1) defines an analytic function provided $\operatorname{Re}(s) > n/2$. Differentiating in s one finds

$$\zeta'(s) = \frac{d}{ds} \sum_{j=1}^{\infty} e^{-s \log \lambda_j} = - \sum_{j=1}^{\infty} \log \lambda_j e^{-s \log \lambda_j}.$$

If ζ is regular near $s = 0$ one can define the *regularized determinant* $\det'(-\Delta_g)$ via the following formula

$$\det'(-\Delta_g) = e^{-\zeta'(0)}.$$

Regularity of ζ at $s = 0$ (in 2D)

Regularity of ζ at $s = 0$ (in 2D)

Let (Σ, g) be a surface.

Regularity of ζ at $s = 0$ (in 2D)

Let (Σ, g) be a surface. One can write

$$\zeta(s) = \sum_{j=1}^{\infty} \lambda_j^{-s} = \frac{1}{\Gamma(s)} \int_0^{\infty} \left(\sum_{j=1}^{\infty} e^{-\lambda_j t} \right) t^s \frac{dt}{t}$$

Regularity of ζ at $s = 0$ (in 2D)

Let (Σ, g) be a surface. One can write

$$\begin{aligned}\zeta(s) &= \sum_{j=1}^{\infty} \lambda_j^{-s} = \frac{1}{\Gamma(s)} \int_0^{\infty} \left(\sum_{j=1}^{\infty} e^{-\lambda_j t} \right) t^s \frac{dt}{t} \\ &= \frac{1}{\Gamma(s)} \int_0^{\infty} (\text{Tr}(e^{\Delta t} - 1)) t^s \frac{dt}{t}.\end{aligned}$$

Regularity of ζ at $s = 0$ (in 2D)

Let (Σ, g) be a surface. One can write

$$\begin{aligned}\zeta(s) &= \sum_{j=1}^{\infty} \lambda_j^{-s} = \frac{1}{\Gamma(s)} \int_0^{\infty} \left(\sum_{j=1}^{\infty} e^{-\lambda_j t} \right) t^s \frac{dt}{t} \\ &= \frac{1}{\Gamma(s)} \int_0^{\infty} (\text{Tr}(e^{\Delta t} - 1)) t^s \frac{dt}{t}.\end{aligned}$$

It is known that (Taylor expand the heat kernel)

$$\sum_{j=1}^{\infty} e^{-\lambda_j t} \varphi_j^2(x) = H_t(x, x)$$

Regularity of ζ at $s = 0$ (in 2D)

Let (Σ, g) be a surface. One can write

$$\begin{aligned}\zeta(s) &= \sum_{j=1}^{\infty} \lambda_j^{-s} = \frac{1}{\Gamma(s)} \int_0^{\infty} \left(\sum_{j=1}^{\infty} e^{-\lambda_j t} \right) t^s \frac{dt}{t} \\ &= \frac{1}{\Gamma(s)} \int_0^{\infty} (\text{Tr}(e^{\Delta t} - 1)) t^s \frac{dt}{t}.\end{aligned}$$

It is known that (Taylor expand the heat kernel)

$$\sum_{j=1}^{\infty} e^{-\lambda_j t} \varphi_j^2(x) = H_t(x, x) = \frac{1}{4\pi t} + \frac{K(x)}{12\pi} + O(t),$$

where K is the Gaussian curvature.

Regularity of ζ at $s = 0$ (in 2D)

Let (Σ, g) be a surface. One can write

$$\begin{aligned}\zeta(s) &= \sum_{j=1}^{\infty} \lambda_j^{-s} = \frac{1}{\Gamma(s)} \int_0^{\infty} \left(\sum_{j=1}^{\infty} e^{-\lambda_j t} \right) t^s \frac{dt}{t} \\ &= \frac{1}{\Gamma(s)} \int_0^{\infty} (\text{Tr}(e^{\Delta t} - 1)) t^s \frac{dt}{t}.\end{aligned}$$

It is known that (Taylor expand the heat kernel)

$$\sum_{j=1}^{\infty} e^{-\lambda_j t} \varphi_j^2(x) = H_t(x, x) = \frac{1}{4\pi t} + \frac{K(x)}{12\pi} + O(t),$$

where K is the Gaussian curvature. Therefore one finds

$$\zeta(s) = \frac{1}{\Gamma(s)} \left\{ \frac{A(\Sigma)}{4\pi(s-1)} + \left(\frac{\chi(\Sigma)}{6} - 1 \right) + \text{holom. in } s \right\},$$

which is regular near zero.

Regularity of ζ at $s = 0$ (in 2D)

Let (Σ, g) be a surface. One can write

$$\begin{aligned}\zeta(s) &= \sum_{j=1}^{\infty} \lambda_j^{-s} = \frac{1}{\Gamma(s)} \int_0^{\infty} \left(\sum_{j=1}^{\infty} e^{-\lambda_j t} \right) t^s \frac{dt}{t} \\ &= \frac{1}{\Gamma(s)} \int_0^{\infty} (\text{Tr}(e^{\Delta t} - 1)) t^s \frac{dt}{t}.\end{aligned}$$

It is known that (Taylor expand the heat kernel)

$$\sum_{j=1}^{\infty} e^{-\lambda_j t} \varphi_j^2(x) = H_t(x, x) = \frac{1}{4\pi t} + \frac{K(x)}{12\pi} + O(t),$$

where K is the Gaussian curvature. Therefore one finds

$$\zeta(s) = \frac{1}{\Gamma(s)} \left\{ \frac{A(\Sigma)}{4\pi(s-1)} + \left(\frac{\chi(\Sigma)}{6} - 1 \right) + \text{holom. in } s \right\},$$

which is regular near zero. $\Rightarrow \det'(-\Delta_g)$ is well defined.

Polyakov's formula for conformal metrics

Polyakov's formula for conformal metrics

In 2D the Laplacian is conformally covariant.

Polyakov's formula for conformal metrics

In 2D the Laplacian is conformally covariant. If $\tilde{g}(x) := e^{2w(x)}g(x)$ is a metric conformal to the original one g , then

$$\Delta_{\tilde{g}} = e^{-2w(x)} \Delta_g; \quad -\Delta_g w + K_g = K_{\tilde{g}} e^{2w}.$$

Polyakov's formula for conformal metrics

In 2D the Laplacian is conformally covariant. If $\tilde{g}(x) := e^{2w(x)}g(x)$ is a metric conformal to the original one g , then

$$\Delta_{\tilde{g}} = e^{-2w(x)}\Delta_g; \quad -\Delta_g w + K_g = K_{\tilde{g}}e^{2w}.$$

These properties allowed Polyakov in '81 to compute the variation of the determinant for conformal metrics with the same volume

Polyakov's formula for conformal metrics

In 2D the Laplacian is conformally covariant. If $\tilde{g}(x) := e^{2w(x)}g(x)$ is a metric conformal to the original one g , then

$$\Delta_{\tilde{g}} = e^{-2w(x)} \Delta_g; \quad -\Delta_g w + K_g = K_{\tilde{g}} e^{2w}.$$

These properties allowed Polyakov in '81 to compute the variation of the determinant for conformal metrics with the same volume

$$\log \det'(-\Delta_{\tilde{g}}) - \log \det'(-\Delta_g) = -\frac{1}{12\pi} \int_{\Sigma} (|\nabla w|^2 + 2Kw) dv.$$

Polyakov's formula for conformal metrics

In 2D the Laplacian is conformally covariant. If $\tilde{g}(x) := e^{2w(x)}g(x)$ is a metric conformal to the original one g , then

$$\Delta_{\tilde{g}} = e^{-2w(x)} \Delta_g; \quad -\Delta_g w + K_g = K_{\tilde{g}} e^{2w}.$$

These properties allowed Polyakov in '81 to compute the variation of the determinant for conformal metrics with the same volume

$$\log \det'(-\Delta_{\tilde{g}}) - \log \det'(-\Delta_g) = -\frac{1}{12\pi} \int_{\Sigma} (|\nabla w|^2 + 2Kw) dv.$$

This formula appears in a partition function in string theory, and is related to the *Moser-Trudinger-Onofri inequality*.

Polyakov's formula for conformal metrics

In 2D the Laplacian is conformally covariant. If $\tilde{g}(x) := e^{2w(x)}g(x)$ is a metric conformal to the original one g , then

$$\Delta_{\tilde{g}} = e^{-2w(x)} \Delta_g; \quad -\Delta_g w + K_g = K_{\tilde{g}} e^{2w}.$$

These properties allowed Polyakov in '81 to compute the variation of the determinant for conformal metrics with the same volume

$$\log \det'(-\Delta_{\tilde{g}}) - \log \det'(-\Delta_g) = -\frac{1}{12\pi} \int_{\Sigma} (|\nabla w|^2 + 2Kw) dv.$$

This formula appears in a partition function in string theory, and is related to the *Moser-Trudinger-Onofri inequality*. On the sphere it is known to be maximised only on conformal factors of Möbius maps.

Polyakov's formula for conformal metrics

In 2D the Laplacian is conformally covariant. If $\tilde{g}(x) := e^{2w(x)}g(x)$ is a metric conformal to the original one g , then

$$\Delta_{\tilde{g}} = e^{-2w(x)} \Delta_g; \quad -\Delta_g w + K_g = K_{\tilde{g}} e^{2w}.$$

These properties allowed Polyakov in '81 to compute the variation of the determinant for conformal metrics with the same volume

$$\log \det'(-\Delta_{\tilde{g}}) - \log \det'(-\Delta_g) = -\frac{1}{12\pi} \int_{\Sigma} (|\nabla w|^2 + 2Kw) dv.$$

This formula appears in a partition function in string theory, and is related to the *Moser-Trudinger-Onofri inequality*. On the sphere it is known to be maximised only on conformal factors of Möbius maps.

Existence of extremals is easy for positive genus.

Polyakov's formula for conformal metrics

In 2D the Laplacian is conformally covariant. If $\tilde{g}(x) := e^{2w(x)}g(x)$ is a metric conformal to the original one g , then

$$\Delta_{\tilde{g}} = e^{-2w(x)} \Delta_g; \quad -\Delta_g w + K_g = K_{\tilde{g}} e^{2w}.$$

These properties allowed Polyakov in '81 to compute the variation of the determinant for conformal metrics with the same volume

$$\log \det'(-\Delta_{\tilde{g}}) - \log \det'(-\Delta_g) = -\frac{1}{12\pi} \int_{\Sigma} (|\nabla w|^2 + 2Kw) \, dv.$$

This formula appears in a partition function in string theory, and is related to the *Moser-Trudinger-Onofri inequality*. On the sphere it is known to be maximised only on conformal factors of Möbius maps.

Existence of extremals is easy for positive genus. On spheres it can be achieved via a *balancing condition* and Möbius invariance, ([Aubin, '76], [Osgood-Phillips-Sarnak, '88], [Gui-Moradifam, '18]).

Isospectral metrics ([Osgood-Phillips-Sarnak, '88])

Isospectral metrics ([Osgood-Phillips-Sarnak, '88])

Isospectral metrics on a closed surface are compact in any C^k sense.

Isospectral metrics ([Osgood-Phillips-Sarnak, '88])

Isospectral metrics on a closed surface are compact in any C^k sense.

Case of the sphere. On S^2 all metrics are conformally equivalent (up to diffeomorphisms).

Isospectral metrics ([Osgood-Phillips-Sarnak, '88])

Isospectral metrics on a closed surface are compact in any C^k sense.

Case of the sphere. On S^2 all metrics are conformally equivalent (up to diffeomorphisms). Since the determinant is bounded, one gets a uniform bound on the $W^{1,2}$ norm of the conformal factor.

Isospectral metrics ([Osgood-Phillips-Sarnak, '88])

Isospectral metrics on a closed surface are compact in any C^k sense.

Case of the sphere. On S^2 all metrics are conformally equivalent (up to diffeomorphisms). Since the determinant is bounded, one gets a uniform bound on the $W^{1,2}$ norm of the conformal factor.

Expanding the heat kernel (via parametrix) one can prove that

$$\sum_{j=1}^{\infty} e^{-\lambda_j t} =: \text{Tr}(e^{\Delta t}) = \frac{1}{t} \sum_{j=0}^l t^j \int_{\Sigma} \Omega_j(x) dV + o(t^l),$$

where Ω_j is a universal polynomial in K_g and Δ_g of degree $2j$.

Isospectral metrics ([Osgood-Phillips-Sarnak, '88])

Isospectral metrics on a closed surface are compact in any C^k sense.

Case of the sphere. On S^2 all metrics are conformally equivalent (up to diffeomorphisms). Since the determinant is bounded, one gets a uniform bound on the $W^{1,2}$ norm of the conformal factor.

Expanding the heat kernel (via parametrix) one can prove that

$$\sum_{j=1}^{\infty} e^{-\lambda_j t} =: \text{Tr}(e^{\Delta t}) = \frac{1}{t} \sum_{j=0}^l t^j \int_{\Sigma} \Omega_j(x) dV + o(t^l),$$

where Ω_j is a universal polynomial in K_g and Δ_g of degree $2j$.

It was proved in [McKean-Singer, '67], [Gilkey, '79] that

$$\Omega_j \simeq \int_{\Sigma} K_g \Delta^{j-2} K_g dv \simeq \|u\|_{W^{j,2}(\Sigma)}$$

Isospectral metrics ([Osgood-Phillips-Sarnak, '88])

Isospectral metrics on a closed surface are compact in any C^k sense.

Case of the sphere. On S^2 all metrics are conformally equivalent (up to diffeomorphisms). Since the determinant is bounded, one gets a uniform bound on the $W^{1,2}$ norm of the conformal factor.

Expanding the heat kernel (via parametrix) one can prove that

$$\sum_{j=1}^{\infty} e^{-\lambda_j t} =: \text{Tr}(e^{\Delta t}) = \frac{1}{t} \sum_{j=0}^l t^j \int_{\Sigma} \Omega_j(x) dV + o(t^l),$$

where Ω_j is a universal polynomial in K_g and Δ_g of degree $2j$.

It was proved in [McKean-Singer, '67], [Gilkey, '79] that

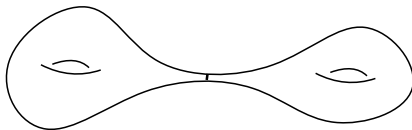
$$\Omega_j \simeq \int_{\Sigma} K_g \Delta^{j-2} K_g dv \simeq \|u\|_{W^{j,2}(\Sigma)},$$

therefore one gets bounds even in higher Sobolev norms.

Isospectral metrics: positive genus

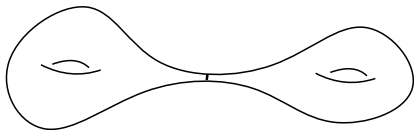
Isospectral metrics: positive genus

Even if, after a conformal change $g \mapsto \hat{g}$, the Gaussian curvature is identically -1 , one could *diverge* in *Teichmüller's space* forming *necks*



Isospectral metrics: positive genus

Even if, after a conformal change $g \mapsto \hat{g}$, the Gaussian curvature is identically -1 , one could *diverge* in *Teichmüller's space* forming *necks*



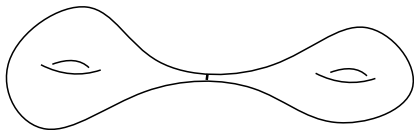
It was however shown in [Wolpert, '87] that

$$\det'(\hat{g}) \leq \frac{1}{l} e^{-\frac{c_1}{l}}; \quad c_1 = c_1(\chi(\Sigma)),$$

where l is the length of the shortest geodesic, so $l \not\rightarrow 0$.

Isospectral metrics: positive genus

Even if, after a conformal change $g \mapsto \hat{g}$, the Gaussian curvature is identically -1 , one could *diverge* in *Teichmüller's space* forming *necks*



It was however shown in [Wolpert, '87] that

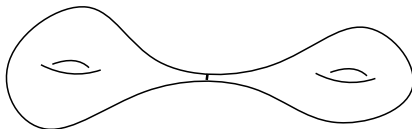
$$\det'(\hat{g}) \leq \frac{1}{l} e^{-\frac{c_1}{l}}; \quad c_1 = c_1(\chi(\Sigma)),$$

where l is the length of the shortest geodesic, so $l \not\rightarrow 0$.

Finally, a theorem in [Mumford, '71] shows that if l is bounded below and if $K_{\hat{g}} = \text{const.}$, then there is smooth convergence of the metrics.

Isospectral metrics: positive genus

Even if, after a conformal change $g \mapsto \hat{g}$, the Gaussian curvature is identically -1 , one could *diverge* in *Teichmüller's space* forming *necks*



It was however shown in [Wolpert, '87] that

$$\det'(\hat{g}) \leq \frac{1}{l} e^{-\frac{c_1}{l}}; \quad c_1 = c_1(\chi(\Sigma)),$$

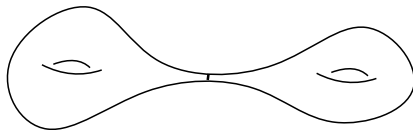
where l is the length of the shortest geodesic, so $l \not\rightarrow 0$.

Finally, a theorem in [Mumford, '71] shows that if l is bounded below and if $K_{\hat{g}} = \text{const.}$, then there is smooth convergence of the metrics.

In higher dimensions very little is known.

Isospectral metrics: positive genus

Even if, after a conformal change $g \mapsto \hat{g}$, the Gaussian curvature is identically -1 , one could *diverge* in *Teichmüller's space* forming *necks*



It was however shown in [Wolpert, '87] that

$$\det'(\hat{g}) \leq \frac{1}{l} e^{-\frac{c_1}{l}}; \quad c_1 = c_1(\chi(\Sigma)),$$

where l is the length of the shortest geodesic, so $l \not\rightarrow 0$.

Finally, a theorem in [Mumford, '71] shows that if l is bounded below and if $K_{\hat{g}} = \text{const.}$, then there is smooth convergence of the metrics.

In higher dimensions very little is known. There are results in special cases like within a conformal class in 3D [Chang-Yang, '90] or under bounded curvature assumptions [G.Zhou, '97].

Conformally covariant operators

Conformally covariant operators

Definition. A linear operator $A = A_g$ is *conformally covariant of bi-degree* (a, b) if $\tilde{g} = e^{2w}g$ implies

$$A_{\tilde{g}}\psi = e^{-bw}A_g(e^{aw}\psi) \quad \text{for each smooth } \psi.$$

Conformally covariant operators

Definition. A linear operator $A = A_g$ is *conformally covariant of bi-degree* (a, b) if $\tilde{g} = e^{2w}g$ implies

$$A_{\tilde{g}}\psi = e^{-bw}A_g(e^{aw}\psi) \quad \text{for each smooth } \psi.$$

Examples

Conformally covariant operators

Definition. A linear operator $A = A_g$ is *conformally covariant of bi-degree* (a, b) if $\tilde{g} = e^{2w}g$ implies

$$A_{\tilde{g}}\psi = e^{-bw}A_g(e^{aw}\psi) \quad \text{for each smooth } \psi.$$

Examples 0. The Laplacian Δ_g for $n = 2$: $(a, b) = (0, 2)$.

Conformally covariant operators

Definition. A linear operator $A = A_g$ is *conformally covariant of bi-degree* (a, b) if $\tilde{g} = e^{2w}g$ implies

$$A_{\tilde{g}}\psi = e^{-bw}A_g(e^{aw}\psi) \quad \text{for each smooth } \psi.$$

Examples **0.** The Laplacian Δ_g for $n = 2$: $(a, b) = (0, 2)$.

1. The *conformal Laplacian* for $n \geq 3$

$$L_g = -\frac{4(n-1)}{(n-2)}\Delta_g + R_g \quad (a, b) = \left(\frac{n-2}{2}, \frac{n+2}{2}\right).$$

Conformally covariant operators

Definition. A linear operator $A = A_g$ is *conformally covariant of bi-degree* (a, b) if $\tilde{g} = e^{2w}g$ implies

$$A_{\tilde{g}}\psi = e^{-bw}A_g(e^{aw}\psi) \quad \text{for each smooth } \psi.$$

Examples **0.** The Laplacian Δ_g for $n = 2$: $(a, b) = (0, 2)$.

1. The *conformal Laplacian* for $n \geq 3$

$$L_g = -\frac{4(n-1)}{(n-2)}\Delta_g + R_g \quad (a, b) = \left(\frac{n-2}{2}, \frac{n+2}{2}\right).$$

2. The *Paneitz operator* P_g for $n = 4$

$$P_g\varphi = (-\Delta_g)^2\varphi + \operatorname{div} \left[\left(\frac{2}{3}Rg - 2\operatorname{Ric} \right) \circ \nabla\varphi \right], \quad (a, b) = (0, 4).$$

Conformally covariant operators

Definition. A linear operator $A = A_g$ is *conformally covariant of bi-degree* (a, b) if $\tilde{g} = e^{2w}g$ implies

$$A_{\tilde{g}}\psi = e^{-bw}A_g(e^{aw}\psi) \quad \text{for each smooth } \psi.$$

Examples **0.** The Laplacian Δ_g for $n = 2$: $(a, b) = (0, 2)$.

1. The *conformal Laplacian* for $n \geq 3$

$$L_g = -\frac{4(n-1)}{(n-2)}\Delta_g + R_g \quad (a, b) = \left(\frac{n-2}{2}, \frac{n+2}{2}\right).$$

2. The *Paneitz operator* P_g for $n = 4$

$$P_g\varphi = (-\Delta_g)^2\varphi + \operatorname{div} \left[\left(\frac{2}{3}Rg - 2\operatorname{Ric} \right) \circ \nabla\varphi \right], \quad (a, b) = (0, 4).$$

3. The *Dirac operator* $\mathcal{D}$ for $n \geq 2$: $(a, b) = \left(\frac{n-1}{2}, \frac{n+1}{2}\right)$.

Determinants of conf. covariant operators in 4D

Determinants of conf. covariant operators in 4D

Theorem ([Branson-Ørsted, '91])

Determinants of conf. covariant operators in 4D

Theorem ([Branson-Ørsted, '91])

Let A be conformally covariant on (M^4, g) .

Determinants of conf. covariant operators in 4D

Theorem ([Branson-Ørsted, '91])

Let A be conformally covariant on (M^4, g) . Then $\exists \gamma_1(A), \gamma_2(A), \gamma_3(A)$ such that for $\tilde{g} = e^{2w}g$

Determinants of conf. covariant operators in 4D

Theorem ([Branson-Ørsted, '91])

Let A be conformally covariant on (M^4, g) . Then $\exists \gamma_1(A), \gamma_2(A), \gamma_3(A)$ such that for $\tilde{g} = e^{2w}g$

$$F_A[w] := \log \frac{\det A_{\tilde{g}}}{\det A_g} = \gamma_1(A)I[w] + \gamma_2(A)II[w] + \gamma_3(A)III[w]$$

Determinants of conf. covariant operators in 4D

Theorem ([Branson-Ørsted, '91])

Let A be conformally covariant on (M^4, g) . Then $\exists \gamma_1(A), \gamma_2(A), \gamma_3(A)$ such that for $\tilde{g} = e^{2w}g$

$$F_A[w] := \log \frac{\det A_{\tilde{g}}}{\det A_g} = \gamma_1(A)I[w] + \gamma_2(A)II[w] + \gamma_3(A)III[w],$$

where

$$I[w] = 4 \int_M w |W_g|^2 dv - \left(\int_M |W_g|^2 dv \right) \log \int_M e^{4w} dv,$$

$$II[w] = \int_M w P_g w dv - \left(\int_M Q_g dv \right) \log \int_M e^{4(w-\bar{w})} dv,$$

$$III[w] = 12 \int_M (\Delta_g w + |\nabla w|^2)^2 dv - 4 \int_M (w \Delta_g R_g + R_g |\nabla w|^2) dv.$$

Determinants of conf. covariant operators in 4D

Theorem ([Branson-Ørsted, '91])

Let A be conformally covariant on (M^4, g) . Then $\exists \gamma_1(A), \gamma_2(A), \gamma_3(A)$ such that for $\tilde{g} = e^{2w}g$

$$F_A[w] := \log \frac{\det A_{\tilde{g}}}{\det A_g} = \gamma_1(A)I[w] + \gamma_2(A)II[w] + \gamma_3(A)III[w],$$

where

$$I[w] = 4 \int_M w |W_g|^2 dv - \left(\int_M |W_g|^2 dv \right) \log \int_M e^{4w} dv,$$

$$II[w] = \int_M w P_g w dv - \left(\int_M Q_g dv \right) \log \int_M e^{4(w-\bar{w})} dv,$$

$$III[w] = 12 \int_M (\Delta_g w + |\nabla w|^2)^2 dv - 4 \int_M (w \Delta_g R_g + R_g |\nabla w|^2) dv.$$

Here W_g is Weyl's curvature, while Q_g is the Q -curvature

Determinants of conf. covariant operators in 4D

Theorem ([Branson-Ørsted, '91])

Let A be conformally covariant on (M^4, g) . Then $\exists \gamma_1(A), \gamma_2(A), \gamma_3(A)$ such that for $\tilde{g} = e^{2w}g$

$$F_A[w] := \log \frac{\det A_{\tilde{g}}}{\det A_g} = \gamma_1(A)I[w] + \gamma_2(A)II[w] + \gamma_3(A)III[w],$$

where

$$I[w] = 4 \int_M w |W_g|^2 dv - \left(\int_M |W_g|^2 dv \right) \log \int_M e^{4w} dv,$$

$$II[w] = \int_M w P_g w dv - \left(\int_M Q_g dv \right) \log \int_M e^{4(w-\bar{w})} dv,$$

$$III[w] = 12 \int_M (\Delta_g w + |\nabla w|^2)^2 dv - 4 \int_M (w \Delta_g R_g + R_g |\nabla w|^2) dv.$$

Here W_g is Weyl's curvature, while Q_g is the Q -curvature, a 4D conformal counterpart of the Gaussian curvature.

Comments

The three functionals I, II, III are quite natural since

The three functionals I, II, III are quite natural since

$$\hat{g} = e^{2w}g \text{ is critical for } I \iff |W_{\hat{g}}|^2 = \text{const.},$$

The three functionals I, II, III are quite natural since

$$\hat{g} = e^{2w}g \text{ is critical for } I \iff |W_{\hat{g}}|^2 = \text{const.},$$

$$\hat{g} \text{ is critical for } II \iff Q_{\hat{g}} = \text{const.},$$

The three functionals I, II, III are quite natural since

$$\hat{g} = e^{2w}g \text{ is critical for } I \iff |W_{\hat{g}}|^2 = \text{const.},$$

$$\hat{g} \text{ is critical for } II \iff Q_{\hat{g}} = \text{const.},$$

$$\hat{g} \text{ is critical for } III \iff \Delta_g R_{\hat{g}} = \text{const.} \quad (\text{Yamabe problem}).$$

The three functionals I, II, III are quite natural since

$$\hat{g} = e^{2w}g \text{ is critical for } I \iff |W_{\hat{g}}|^2 = \text{const.},$$

$$\hat{g} \text{ is critical for } II \iff Q_{\hat{g}} = \text{const.},$$

$$\hat{g} \text{ is critical for } III \iff \Delta_g R_{\hat{g}} = \text{const.} \quad (\text{Yamabe problem}).$$

Also, in 4D there is a Gauss-Bonnet formula

$$\int_M \left(Q_g + \frac{1}{8} |W_g|^2 \right) dv = 4\pi^2 \chi(M).$$

The three functionals I, II, III are quite natural since

$$\hat{g} = e^{2w}g \text{ is critical for } I \iff |W_{\hat{g}}|^2 = \text{const.},$$

$$\hat{g} \text{ is critical for } II \iff Q_{\hat{g}} = \text{const.},$$

$$\hat{g} \text{ is critical for } III \iff \Delta_g R_{\hat{g}} = \text{const.} \quad (\text{Yamabe problem}).$$

Also, in 4D there is a Gauss-Bonnet formula

$$\int_M \left(Q_g + \frac{1}{8} |W_g|^2 \right) dv = 4\pi^2 \chi(M).$$

Each term separately is not a topological invariant.

The three functionals I, II, III are quite natural since

$$\hat{g} = e^{2w}g \text{ is critical for } I \iff |W_{\hat{g}}|^2 = \text{const.},$$

$$\hat{g} \text{ is critical for } II \iff Q_{\hat{g}} = \text{const.},$$

$$\hat{g} \text{ is critical for } III \iff \Delta_g R_{\hat{g}} = \text{const.} \quad (\text{Yamabe problem}).$$

Also, in 4D there is a Gauss-Bonnet formula

$$\int_M \left(Q_g + \frac{1}{8} |W_g|^2 \right) dv = 4\pi^2 \chi(M).$$

Each term separately is not a topological invariant. However, both $\int_M Q_g dv$ and $\int_M |W_g|^2 dv$ are conformally invariant.

The three functionals I, II, III are quite natural since

$$\hat{g} = e^{2w}g \text{ is critical for } I \iff |W_{\hat{g}}|^2 = \text{const.},$$

$$\hat{g} \text{ is critical for } II \iff Q_{\hat{g}} = \text{const.},$$

$$\hat{g} \text{ is critical for } III \iff \Delta_g R_{\hat{g}} = \text{const.} \quad (\text{Yamabe problem}).$$

Also, in 4D there is a Gauss-Bonnet formula

$$\int_M \left(Q_g + \frac{1}{8} |W_g|^2 \right) dv = 4\pi^2 \chi(M).$$

Each term separately is not a topological invariant. However, both $\int_M Q_g dv$ and $\int_M |W_g|^2 dv$ are conformally invariant.

- Extremal metrics for linear combinations of the functionals I, II, III were useful in studying rigidity of K-E metrics in 4D ([Gursky, '98]).

Explicit constants

In the above examples

Explicit constants

In the above examples

- If $A_g = L_g$, the conformal Laplacian, then

$$(\gamma_1, \gamma_2, \gamma_3) = \left(1, -4, -\frac{2}{3}\right).$$

Explicit constants

In the above examples

- If $A_g = L_g$, the conformal Laplacian, then

$$(\gamma_1, \gamma_2, \gamma_3) = \left(1, -4, -\frac{2}{3}\right).$$

- If $A_g = P_g$, the Paneitz operator, then

$$(\gamma_1, \gamma_2, \gamma_3) = \left(-\frac{1}{4}, -14, \frac{8}{3}\right).$$

Explicit constants

In the above examples

- If $A_g = L_g$, the conformal Laplacian, then

$$(\gamma_1, \gamma_2, \gamma_3) = \left(1, -4, -\frac{2}{3}\right).$$

- If $A_g = P_g$, the Paneitz operator, then

$$(\gamma_1, \gamma_2, \gamma_3) = \left(-\frac{1}{4}, -14, \frac{8}{3}\right).$$

- If $A_g = \mathcal{D}$, the Dirac operator, then

$$(\gamma_1, \gamma_2, \gamma_3) = \left(-7, -88, -\frac{14}{3}\right).$$

Explicit constants

In the above examples

- If $A_g = L_g$, the conformal Laplacian, then

$$(\gamma_1, \gamma_2, \gamma_3) = \left(1, -4, -\frac{2}{3}\right).$$

- If $A_g = P_g$, the Paneitz operator, then

$$(\gamma_1, \gamma_2, \gamma_3) = \left(-\frac{1}{4}, -14, \frac{8}{3}\right).$$

- If $A_g = \mathcal{D}$, the Dirac operator, then

$$(\gamma_1, \gamma_2, \gamma_3) = \left(-7, -88, -\frac{14}{3}\right).$$

- Sometimes we will reverse signs to get coercivity/convexity.

Extremals of determinants in 4D

Extremals of determinants in 4D

Theorem ([Chang-Yang, '95]) For $n = 4$ assume:

$$(i) \ \gamma_2, \gamma_3 > 0,$$

Extremals of determinants in 4D

Theorem ([Chang-Yang, '95]) For $n = 4$ assume:

$$(i) \ \gamma_2, \gamma_3 > 0, \quad (ii) \ \gamma_1 \int_M |W_g|^2 \, dv + \gamma_2 \int_M Q_g \, dv < 8\gamma_2\pi^2.$$

Extremals of determinants in 4D

Theorem ([Chang-Yang, '95]) For $n = 4$ assume:

$$(i) \ \gamma_2, \gamma_3 > 0, \quad (ii) \ \gamma_1 \int_M |W_g|^2 \, dv + \gamma_2 \int_M Q_g \, dv < 8\gamma_2\pi^2.$$

Then $\inf_{w \in W^{2,2}} F_A[w]$ is attained.

Extremals of determinants in 4D

Theorem ([Chang-Yang, '95]) For $n = 4$ assume:

$$(i) \ \gamma_2, \gamma_3 > 0, \quad (ii) \ \gamma_1 \int_M |W_g|^2 \, dv + \gamma_2 \int_M Q_g \, dv < 8\gamma_2\pi^2.$$

Then $\inf_{w \in W^{2,2}} F_A[w]$ is attained.

Remarks - (ii) implies coercivity of F_A , via sharp Moser-Trudinger inequalities ([Adams, '88]): direct methods yield a maximizer for F_A .

Extremals of determinants in 4D

Theorem ([Chang-Yang, '95]) For $n = 4$ assume:

$$(i) \ \gamma_2, \gamma_3 > 0, \quad (ii) \ \gamma_1 \int_M |W_g|^2 \, dv + \gamma_2 \int_M Q_g \, dv < 8\gamma_2\pi^2.$$

Then $\inf_{w \in W^{2,2}} F_A[w]$ is attained.

Remarks - (ii) implies coercivity of F_A , via sharp Moser-Trudinger inequalities ([Adams, '88]): direct methods yield a maximizer for F_A .

- The assumptions are conformally invariant and are satisfied (roughly) in positive curvature ([Gursky, '99]).

Extremals of determinants in 4D

Theorem ([Chang-Yang, '95]) For $n = 4$ assume:

$$(i) \ \gamma_2, \gamma_3 > 0, \quad (ii) \ \gamma_1 \int_M |W_g|^2 \, dv + \gamma_2 \int_M Q_g \, dv < 8\gamma_2\pi^2.$$

Then $\inf_{w \in W^{2,2}} F_A[w]$ is attained.

Remarks - (ii) implies coercivity of F_A , via sharp Moser-Trudinger inequalities ([Adams, '88]): direct methods yield a maximizer for F_A .

- The assumptions are conformally invariant and are satisfied (roughly) in positive curvature ([Gursky, '99]). In negative curvature there are many examples of *large Gauss-Bonnet integrals*.

Extremals of determinants in 4D

Theorem ([Chang-Yang, '95]) For $n = 4$ assume:

$$(i) \ \gamma_2, \gamma_3 > 0, \quad (ii) \ \gamma_1 \int_M |W_g|^2 \, dv + \gamma_2 \int_M Q_g \, dv < 8\gamma_2\pi^2.$$

Then $\inf_{w \in W^{2,2}} F_A[w]$ is attained.

Remarks - (ii) implies coercivity of F_A , via sharp Moser-Trudinger inequalities ([Adams, '88]): direct methods yield a maximizer for F_A .

- The assumptions are conformally invariant and are satisfied (roughly) in positive curvature ([Gursky, '99]). In negative curvature there are many examples of *large Gauss-Bonnet integrals*.

Open: uniqueness

Extremals of determinants in 4D

Theorem ([Chang-Yang, '95]) For $n = 4$ assume:

$$(i) \ \gamma_2, \gamma_3 > 0, \quad (ii) \ \gamma_1 \int_M |W_g|^2 \, dv + \gamma_2 \int_M Q_g \, dv < 8\gamma_2\pi^2.$$

Then $\inf_{w \in W^{2,2}} F_A[w]$ is attained.

Remarks - (ii) implies coercivity of F_A , via sharp Moser-Trudinger inequalities ([Adams, '88]): direct methods yield a maximizer for F_A .

- The assumptions are conformally invariant and are satisfied (roughly) in positive curvature ([Gursky, '99]). In negative curvature there are many examples of *large Gauss-Bonnet integrals*.

Open: uniqueness ([Gursky-Streets, '18] for the σ_2 -equation).

Extremals of determinants in 4D

Theorem ([Chang-Yang, '95]) For $n = 4$ assume:

$$(i) \ \gamma_2, \gamma_3 > 0, \quad (ii) \ \gamma_1 \int_M |W_g|^2 \, dv + \gamma_2 \int_M Q_g \, dv < 8\gamma_2\pi^2.$$

Then $\inf_{w \in W^{2,2}} F_A[w]$ is attained.

Remarks - (ii) implies coercivity of F_A , via sharp Moser-Trudinger inequalities ([Adams, '88]): direct methods yield a maximizer for F_A .

- The assumptions are conformally invariant and are satisfied (roughly) in positive curvature ([Gursky, '99]). In negative curvature there are many examples of *large Gauss-Bonnet integrals*.

Open: uniqueness ([Gursky-Streets, '18] for the σ_2 -equation).

- We aim to discuss here the situations when either (ii) fails (e.g. in negative curvature), or when (i) fails (as for the Paneitz operator).

Extremals of determinants in 4D

Theorem ([Chang-Yang, '95]) For $n = 4$ assume:

$$(i) \ \gamma_2, \gamma_3 > 0, \quad (ii) \ \gamma_1 \int_M |W_g|^2 \, dv + \gamma_2 \int_M Q_g \, dv < 8\gamma_2\pi^2.$$

Then $\inf_{w \in W^{2,2}} F_A[w]$ is attained.

Remarks - (ii) implies coercivity of F_A , via sharp Moser-Trudinger inequalities ([Adams, '88]): direct methods yield a maximizer for F_A .

- The assumptions are conformally invariant and are satisfied (roughly) in positive curvature ([Gursky, '99]). In negative curvature there are many examples of *large Gauss-Bonnet integrals*.

Open: uniqueness ([Gursky-Streets, '18] for the σ_2 -equation).

- We aim to discuss here the situations when either (ii) fails (e.g. in negative curvature), or when (i) fails (as for the Paneitz operator). The latter case is indeed much harder.

On the functional II (special case: $(\gamma_1, \gamma_3) = (0, 0)$)

On the functional II (special case: $(\gamma_1, \gamma_3) = (0, 0)$)

(Conformal) extremals of II , having constant Q -curvature, solve

$$P_g u + 2Q_g = 2\overline{Q}e^{4u}; \quad \overline{Q} \in \mathbb{R}.$$

On the functional II (special case: $(\gamma_1, \gamma_3) = (0, 0)$)

(Conformal) extremals of II , having constant Q -curvature, solve

$$P_g u + 2Q_g = 2\overline{Q}e^{4u}; \quad \overline{Q} \in \mathbb{R}.$$

The result in [Chang-Yang, '95] applies when $k_Q := \int_M Q_g < 8\pi^2$.

On the functional II (special case: $(\gamma_1, \gamma_3) = (0, 0)$)

(Conformal) extremals of II , having constant Q -curvature, solve

$$P_g u + 2Q_g = 2\overline{Q}e^{4u}; \quad \overline{Q} \in \mathbb{R}.$$

The result in [Chang-Yang, '95] applies when $k_Q := \int_M Q_g < 8\pi^2$. If instead $k_Q > 8\pi^2$, then II is unbounded on both sides: k_Q *beats* the Moser-Trudinger constant.

On the functional II (special case: $(\gamma_1, \gamma_3) = (0, 0)$)

(Conformal) extremals of II , having constant Q -curvature, solve

$$P_g u + 2Q_g = 2\overline{Q}e^{4u}; \quad \overline{Q} \in \mathbb{R}.$$

The result in [Chang-Yang, '95] applies when $k_Q := \int_M Q_g < 8\pi^2$. If instead $k_Q > 8\pi^2$, then II is unbounded on both sides: k_Q *beats* the Moser-Trudinger constant.

Still, in [Djadli-M., '08] existence was found provided $k_Q \notin 8\pi^2\mathbb{N}$.

On the functional II (special case: $(\gamma_1, \gamma_3) = (0, 0)$)

(Conformal) extremals of II , having constant Q -curvature, solve

$$P_g u + 2Q_g = 2\overline{Q}e^{4u}; \quad \overline{Q} \in \mathbb{R}.$$

The result in [Chang-Yang, '95] applies when $k_Q := \int_M Q_g < 8\pi^2$. If instead $k_Q > 8\pi^2$, then II is unbounded on both sides: k_Q *beats* the Moser-Trudinger constant.

Still, in [Djadli-M., '08] existence was found provided $k_Q \notin 8\pi^2\mathbb{N}$. The main tool are improved M-T inequalities, in the spirit of [Aubin, '76]: *spreading* of conformal volume leads to better functional inequalities.

On the functional II (special case: $(\gamma_1, \gamma_3) = (0, 0)$)

(Conformal) extremals of II , having constant Q -curvature, solve

$$P_g u + 2Q_g = 2\bar{Q}e^{4u}; \quad \bar{Q} \in \mathbb{R}.$$

The result in [Chang-Yang, '95] applies when $k_Q := \int_M Q_g < 8\pi^2$. If instead $k_Q > 8\pi^2$, then II is unbounded on both sides: k_Q *beats* the Moser-Trudinger constant.

Still, in [Djadli-M., '08] existence was found provided $k_Q \notin 8\pi^2\mathbb{N}$. The main tool are improved M-T inequalities, in the spirit of [Aubin, '76]: *spreading* of conformal volume leads to better functional inequalities.

If for example, if $k_Q \in (8\pi^2, 16\pi^2)$ and if II is large negative, then the conformal volume must *concentrate* near a single point of M .

On the functional II (special case: $(\gamma_1, \gamma_3) = (0, 0)$)

(Conformal) extremals of II , having constant Q -curvature, solve

$$P_g u + 2Q_g = 2\overline{Q}e^{4u}; \quad \overline{Q} \in \mathbb{R}.$$

The result in [Chang-Yang, '95] applies when $k_Q := \int_M Q_g < 8\pi^2$. If instead $k_Q > 8\pi^2$, then II is unbounded on both sides: k_Q *beats* the Moser-Trudinger constant.

Still, in [Djadli-M., '08] existence was found provided $k_Q \notin 8\pi^2\mathbb{N}$. The main tool are improved M-T inequalities, in the spirit of [Aubin', 76]: *spreading* of conformal volume leads to better functional inequalities.

If for example, if $k_Q \in (8\pi^2, 16\pi^2)$ and if II is large negative, then the conformal volume must *concentrate* near a single point of M . One can then exploit the topology of M to find extremals of *min-max type*.

On the functional II (special case: $(\gamma_1, \gamma_3) = (0, 0)$)

(Conformal) extremals of II , having constant Q -curvature, solve

$$P_g u + 2Q_g = 2\bar{Q}e^{4u}; \quad \bar{Q} \in \mathbb{R}.$$

The result in [Chang-Yang, '95] applies when $k_Q := \int_M Q_g < 8\pi^2$. If instead $k_Q > 8\pi^2$, then II is unbounded on both sides: k_Q *beats* the Moser-Trudinger constant.

Still, in [Djadli-M., '08] existence was found provided $k_Q \notin 8\pi^2\mathbb{N}$. The main tool are improved M-T inequalities, in the spirit of [Aubin', 76]: *spreading* of conformal volume leads to better functional inequalities.

If for example, if $k_Q \in (8\pi^2, 16\pi^2)$ and if II is large negative, then the conformal volume must *concentrate* near a single point of M . One can then exploit the topology of M to find extremals of *min-max type*. This also works for more general determinant functionals, provided $\gamma_2, \gamma_3 > 0$.

Compactness and quantization for Q -curvature

Compactness and quantization for Q -curvature

The topological structure of the energy (joint with a *monotonicity argument* by Struwe) allows to produce solutions of *perturbed equations*

$$P_g u_n + 2Q_n = 2\overline{Q}_n e^{4u_n}; \quad Q_n \rightarrow Q_g, \quad \overline{Q}_n \rightarrow \overline{Q}.$$

Compactness and quantization for Q -curvature

The topological structure of the energy (joint with a *monotonicity argument* by Struwe) allows to produce solutions of *perturbed equations*

$$P_g u_n + 2Q_n = 2\overline{Q}_n e^{4u_n}; \quad Q_n \rightarrow Q_g, \quad \overline{Q}_n \rightarrow \overline{Q}.$$

We wish then to pass to the limit, but in general solutions might *blow-up*, and one tries to reach a contradiction.

Compactness and quantization for Q -curvature

The topological structure of the energy (joint with a *monotonicity argument* by Struwe) allows to produce solutions of *perturbed equations*

$$P_g u_n + 2Q_n = 2\overline{Q}_n e^{4u_n}; \quad Q_n \rightarrow Q_g, \quad \overline{Q}_n \rightarrow \overline{Q}.$$

We wish then to pass to the limit, but in general solutions might *blow-up*, and one tries to reach a contradiction.

If blow-up occurs, use Green's formula to show that e^{4u_n} accumulates at finitely-many points ([Brezis-Merle', 91]), so $u_n - \bar{u}_n \rightarrow u_s$, with u_s s.t.

$$P_g u_s + 2Q_g = \sum_{i=1}^l \beta_i \delta_{p_i}; \quad \beta_i > 0.$$

Compactness and quantization for Q -curvature

The topological structure of the energy (joint with a *monotonicity argument* by Struwe) allows to produce solutions of *perturbed equations*

$$P_g u_n + 2Q_n = 2\overline{Q}_n e^{4u_n}; \quad Q_n \rightarrow Q_g, \quad \overline{Q}_n \rightarrow \overline{Q}.$$

We wish then to pass to the limit, but in general solutions might *blow-up*, and one tries to reach a contradiction.

If blow-up occurs, use Green's formula to show that e^{4u_n} accumulates at finitely-many points ([Brezis-Merle', 91]), so $u_n - \bar{u}_n \rightarrow u_s$, with u_s s.t.

$$P_g u_s + 2Q_g = \sum_{i=1}^l \beta_i \delta_{p_i}; \quad \beta_i > 0.$$

Since the operator on the l.h.s. is linear, the singular solution is a linear combinations of (logarithmic) Green's functions.

Compactness and quantization for Q -curvature

The topological structure of the energy (joint with a *monotonicity argument* by Struwe) allows to produce solutions of *perturbed equations*

$$P_g u_n + 2Q_n = 2\overline{Q}_n e^{4u_n}; \quad Q_n \rightarrow Q_g, \quad \overline{Q}_n \rightarrow \overline{Q}.$$

We wish then to pass to the limit, but in general solutions might *blow-up*, and one tries to reach a contradiction.

If blow-up occurs, use Green's formula to show that e^{4u_n} accumulates at finitely-many points ([Brezis-Merle', 91]), so $u_n - \bar{u}_n \rightarrow u_s$, with u_s s.t.

$$P_g u_s + 2Q_g = \sum_{i=1}^l \beta_i \delta_{p_i}; \quad \beta_i > 0.$$

Since the operator on the l.h.s. is linear, the singular solution is a linear combinations of (logarithmic) Green's functions.

Finally *bubbling analysis*, shows that $\beta_i = 8\pi^2$ for all i ([Li-Shafrir, '93], [Druet-Robert, '06], [M., '06]), a contradiction to $k_Q \notin 8\pi^2\mathbb{N}$.

Compactness of extremal metrics for F_L

Compactness of extremal metrics for F_L

We focus next on the log-determinant of the conformal Laplacian L (some results apply to more general F_A 's).

Compactness of extremal metrics for F_L

We focus next on the log-determinant of the conformal Laplacian L (some results apply to more general F_A 's). Extremal metrics satisfy

$$\mathcal{N}(u) + U_g = \mu e^{4u}, \quad \text{with } \mathcal{N}(u) \text{ in divergence form.}$$

Compactness of extremal metrics for F_L

We focus next on the log-determinant of the conformal Laplacian L (some results apply to more general F_A 's). Extremal metrics satisfy

$$\mathcal{N}(u) + U_g = \mu e^{4u}, \quad \text{with } \mathcal{N}(u) \text{ in divergence form.}$$

$$\mathcal{N}_L(u) \simeq \Delta^2 u - \Delta_4 u; \quad \Delta_4 u = \operatorname{div}(|\nabla u|^2 \nabla u),$$

$$U_g = \gamma_1 |W_g|_g^2 + \gamma_2 Q_g - \gamma_3 \Delta_g R_g; \quad \mu = \int_M U_g dv.$$

Compactness of extremal metrics for F_L

We focus next on the log-determinant of the conformal Laplacian L (some results apply to more general F_A 's). Extremal metrics satisfy

$$\mathcal{N}(u) + U_g = \mu e^{4u}, \quad \text{with } \mathcal{N}(u) \text{ in divergence form.}$$

$$\mathcal{N}_L(u) \simeq \Delta^2 u - \Delta_4 u; \quad \Delta_4 u = \operatorname{div}(|\nabla u|^2 \nabla u),$$

$$U_g = \gamma_1 |W_g|_g^2 + \gamma_2 Q_g - \gamma_3 \Delta_g R_g; \quad \mu = \int_M U_g dv.$$

Theorem A ([Esposito-M., w.i.p.]

Suppose $U_n \rightarrow U_g$ and $\mu_n \rightarrow \mu$ in $C^1(M^4)$.

Compactness of extremal metrics for F_L

We focus next on the log-determinant of the conformal Laplacian L (some results apply to more general F_A 's). Extremal metrics satisfy

$$\mathcal{N}(u) + U_g = \mu e^{4u}, \quad \text{with } \mathcal{N}(u) \text{ in divergence form.}$$

$$\mathcal{N}_L(u) \simeq \Delta^2 u - \Delta_4 u; \quad \Delta_4 u = \operatorname{div}(|\nabla u|^2 \nabla u),$$

$$U_g = \gamma_1 |W_g|_g^2 + \gamma_2 Q_g - \gamma_3 \Delta_g R_g; \quad \mu = \int_M U_g dv.$$

Theorem A ([Esposito-M., w.i.p.]

Suppose $U_n \rightarrow U_g$ and $\mu_n \rightarrow \mu$ in $C^1(M^4)$. Let u_n solve

$$\mathcal{N}_L(u_n) + U_n = \mu_n e^{4u_n}, \quad \text{with } \int_M e^{4u_n} dv \leq C.$$

Compactness of extremal metrics for F_L

We focus next on the log-determinant of the conformal Laplacian L (some results apply to more general F_A 's). Extremal metrics satisfy

$$\mathcal{N}(u) + U_g = \mu e^{4u}, \quad \text{with } \mathcal{N}(u) \text{ in divergence form.}$$

$$\mathcal{N}_L(u) \simeq \Delta^2 u - \Delta_4 u; \quad \Delta_4 u = \operatorname{div}(|\nabla u|^2 \nabla u),$$

$$U_g = \gamma_1 |W_g|_g^2 + \gamma_2 Q_g - \gamma_3 \Delta_g R_g; \quad \mu = \int_M U_g dv.$$

Theorem A ([Esposito-M., w.i.p.])

Suppose $U_n \rightarrow U_g$ and $\mu_n \rightarrow \mu$ in $C^1(M^4)$. Let u_n solve

$$\mathcal{N}_L(u_n) + U_n = \mu_n e^{4u_n}, \quad \text{with } \int_M e^{4u_n} dv \leq C.$$

Then either $(u_n)_n$ stays bounded in $C^{4,\alpha}(M)$ or $\mu_n e^{4u_n} \rightharpoonup 8\pi^2 \sum_{i=1}^l \delta_{p_i}$, for distinct points $p_1, \dots, p_l$.

Compactness of extremal metrics for F_L

We focus next on the log-determinant of the conformal Laplacian L (some results apply to more general F_A 's). Extremal metrics satisfy

$$\mathcal{N}(u) + U_g = \mu e^{4u}, \quad \text{with } \mathcal{N}(u) \text{ in divergence form.}$$

$$\mathcal{N}_L(u) \simeq \Delta^2 u - \Delta_4 u; \quad \Delta_4 u = \operatorname{div}(|\nabla u|^2 \nabla u),$$

$$U_g = \gamma_1 |W_g|_g^2 + \gamma_2 Q_g - \gamma_3 \Delta_g R_g; \quad \mu = \int_M U_g dv.$$

Theorem A ([Esposito-M., w.i.p.])

Suppose $U_n \rightarrow U_g$ and $\mu_n \rightarrow \mu$ in $C^1(M^4)$. Let u_n solve

$$\mathcal{N}_L(u_n) + U_n = \mu_n e^{4u_n}, \quad \text{with } \int_M e^{4u_n} dv \leq C.$$

Then either $(u_n)_n$ stays bounded in $C^{4,\alpha}(M)$ or $\mu_n e^{4u_n} \rightharpoonup 8\pi^2 \sum_{i=1}^l \delta_{p_i}$, for distinct points $p_1, \dots, p_l$. In the latter case $\mu \in 8\pi^2 \mathbb{N}$.

Uniform bounds and ε -regularity $\left(\frac{\gamma_2}{\gamma_3} > \frac{3}{2}\right)$

Uniform bounds and ε -regularity $\left(\frac{\gamma_2}{\gamma_3} > \frac{3}{2}\right)$

Solutions *blow-up* when $\sup_M u_n \rightarrow +\infty$ for $n \rightarrow +\infty$.

Uniform bounds and ε -regularity $\left(\frac{\gamma_2}{\gamma_3} > \frac{3}{2}\right)$

Solutions *blow-up* when $\sup_M u_n \rightarrow +\infty$ for $n \rightarrow +\infty$. However some quantities stay uniformly bounded (thinking of $\log |x|$, just missing $W^{2,2}$ and $W^{1,4}$).

Uniform bounds and ε -regularity $\left(\frac{\gamma_2}{\gamma_3} > \frac{3}{2}\right)$

Solutions *blow-up* when $\sup_M u_n \rightarrow +\infty$ for $n \rightarrow +\infty$. However some quantities stay uniformly bounded (thinking of $\log|x|$, just missing $W^{2,2}$ and $W^{1,4}$). It can be shown that ([Dolzmann-Hungerbühler-Müller, '00])

$$\|u_n - \bar{u}_n\|_{W^{2,q}} \leq C_q \quad \text{for } q < 2; \quad [u_n]_{BMO} \leq C.$$

Uniform bounds and ε -regularity $\left(\frac{\gamma_2}{\gamma_3} > \frac{3}{2}\right)$

Solutions *blow-up* when $\sup_M u_n \rightarrow +\infty$ for $n \rightarrow +\infty$. However some quantities stay uniformly bounded (thinking of $\log|x|$, just missing $W^{2,2}$ and $W^{1,4}$). It can be shown that ([Dolzmann-Hungerbühler-Müller, '00])

$$\|u_n - \bar{u}_n\|_{W^{2,q}} \leq C_q \quad \text{for } q < 2; \quad [u_n]_{BMO} \leq C.$$

In [Uhlenbeck-Viaclovsky, '00] an ε -regularity result was proved:

Uniform bounds and ε -regularity $\left(\frac{\gamma_2}{\gamma_3} > \frac{3}{2}\right)$

Solutions *blow-up* when $\sup_M u_n \rightarrow +\infty$ for $n \rightarrow +\infty$. However some quantities stay uniformly bounded (thinking of $\log |x|$, just missing $W^{2,2}$ and $W^{1,4}$). It can be shown that ([Dolzmann-Hungerbühler-Müller, '00])

$$\|u_n - \bar{u}_n\|_{W^{2,q}} \leq C_q \quad \text{for } q < 2; \quad [u_n]_{BMO} \leq C.$$

In [Uhlenbeck-Viaclovsky, '00] an ε -regularity result was proved:

$$\int_{B_{2r}(p)} e^{4u} dv < \varepsilon_0 \quad \implies \quad \int_{B_r(p)} (|\nabla^2 u|^2 + |\nabla u|^4) dv \leq C.$$

Uniform bounds and ε -regularity $\left(\frac{\gamma_2}{\gamma_3} > \frac{3}{2}\right)$

Solutions *blow-up* when $\sup_M u_n \rightarrow +\infty$ for $n \rightarrow +\infty$. However some quantities stay uniformly bounded (thinking of $\log|x|$, just missing $W^{2,2}$ and $W^{1,4}$). It can be shown that ([Dolzmann-Hungerbühler-Müller, '00])

$$\|u_n - \bar{u}_n\|_{W^{2,q}} \leq C_q \quad \text{for } q < 2; \quad [u_n]_{BMO} \leq C.$$

In [Uhlenbeck-Viaclovsky, '00] an ε -regularity result was proved:

$$\int_{B_{2r}(p)} e^{4u} dv < \varepsilon_0 \quad \implies \quad \int_{B_r(p)} (|\nabla^2 u|^2 + |\nabla u|^4) dv \leq C.$$

The M-T inequality then implies $e^{4u} \in L^q(B_r(p))$ near p for some $q > 1$.

Uniform bounds and ε -regularity $\left(\frac{\gamma_2}{\gamma_3} > \frac{3}{2}\right)$

Solutions *blow-up* when $\sup_M u_n \rightarrow +\infty$ for $n \rightarrow +\infty$. However some quantities stay uniformly bounded (thinking of $\log |x|$, just missing $W^{2,2}$ and $W^{1,4}$). It can be shown that ([Dolzmann-Hungerbühler-Müller, '00])

$$\|u_n - \bar{u}_n\|_{W^{2,q}} \leq C_q \quad \text{for } q < 2; \quad [u_n]_{BMO} \leq C.$$

In [Uhlenbeck-Viaclovsky, '00] an ε -regularity result was proved:

$$\int_{B_{2r}(p)} e^{4u} dv < \varepsilon_0 \quad \implies \quad \int_{B_r(p)} (|\nabla^2 u|^2 + |\nabla u|^4) dv \leq C.$$

The M-T inequality then implies $e^{4u} \in L^q(B_r(p))$ near p for some $q > 1$.

Consequence. At blow-up points concentrates at least ε_0 volume, so the set of blow-up points is finite.

Existence of *fundamental solutions* $(\gamma_2, \gamma_3 > 0)$

Existence of *fundamental solutions* $(\gamma_2, \gamma_3 > 0)$

Proposition 1

There exists a distributional solution of $\mathcal{N}(u_s) + U_g = \sum_{i=1}^l \beta_i \delta_{p_i}$ such that $u_s = \alpha_i \log d(x, p_i) + w$ near p_i , with $\alpha_i = \alpha_i(\beta_i) < 0$ (explicit), and

$$\lim_{x \rightarrow 0} |x|^k |\nabla^{(k)} w| = 0 \quad \forall k = 1, 2, 3.$$

Existence of *fundamental solutions* $(\gamma_2, \gamma_3 > 0)$

Proposition 1

There exists a distributional solution of $\mathcal{N}(u_s) + U_g = \sum_{i=1}^l \beta_i \delta_{p_i}$ such that $u_s = \alpha_i \log d(x, p_i) + w$ near p_i , with $\alpha_i = \alpha_i(\beta_i) < 0$ (explicit), and

$$\lim_{x \rightarrow 0} |x|^k |\nabla^{(k)} w| = 0 \quad \forall k = 1, 2, 3.$$

To prove existence, one can use an approximate solution u_{app} of the form

$$u_{\text{app}}(x) \simeq \sum_{i=1}^l \alpha_i \log d(x, p_i) \quad \alpha_i = \alpha_i(\beta_i).$$

Existence of *fundamental solutions* $(\gamma_2, \gamma_3 > 0)$

Proposition 1

There exists a distributional solution of $\mathcal{N}(u_s) + U_g = \sum_{i=1}^l \beta_i \delta_{p_i}$ such that $u_s = \alpha_i \log d(x, p_i) + w$ near p_i , with $\alpha_i = \alpha_i(\beta_i) < 0$ (explicit), and

$$\lim_{x \rightarrow 0} |x|^k |\nabla^{(k)} w| = 0 \quad \forall k = 1, 2, 3.$$

To prove existence, one can use an approximate solution u_{app} of the form

$$u_{\text{app}}(x) \simeq \sum_{i=1}^l \alpha_i \log d(x, p_i) \quad \alpha_i = \alpha_i(\beta_i).$$

$(M \setminus \{p_1, \dots, p_l\}, e^{4u_{\text{app}}})$ has conical points and/or conical/cylindrical ends.

Existence of *fundamental solutions* $(\gamma_2, \gamma_3 > 0)$

Proposition 1

There exists a distributional solution of $\mathcal{N}(u_s) + U_g = \sum_{i=1}^l \beta_i \delta_{p_i}$ such that $u_s = \alpha_i \log d(x, p_i) + w$ near p_i , with $\alpha_i = \alpha_i(\beta_i) < 0$ (explicit), and

$$\lim_{x \rightarrow 0} |x|^k |\nabla^{(k)} w| = 0 \quad \forall k = 1, 2, 3.$$

To prove existence, one can use an approximate solution u_{app} of the form

$$u_{\text{app}}(x) \simeq \sum_{i=1}^l \alpha_i \log d(x, p_i) \quad \alpha_i = \alpha_i(\beta_i).$$

$(M \setminus \{p_1, \dots, p_l\}, e^{4u_{\text{app}}})$ has conical points and/or conical/cylindrical ends. Here one gains the variational structure, obtaining existence with exponential $W^{2,2}$ -decay along the ends.

Existence of *fundamental solutions* $(\gamma_2, \gamma_3 > 0)$

Proposition 1

There exists a distributional solution of $\mathcal{N}(u_s) + U_g = \sum_{i=1}^l \beta_i \delta_{p_i}$ such that $u_s = \alpha_i \log d(x, p_i) + w$ near p_i , with $\alpha_i = \alpha_i(\beta_i) < 0$ (explicit), and

$$\lim_{x \rightarrow 0} |x|^k |\nabla^{(k)} w| = 0 \quad \forall k = 1, 2, 3.$$

To prove existence, one can use an approximate solution u_{app} of the form

$$u_{\text{app}}(x) \simeq \sum_{i=1}^l \alpha_i \log d(x, p_i) \quad \alpha_i = \alpha_i(\beta_i).$$

$(M \setminus \{p_1, \dots, p_l\}, e^{4u_{\text{app}}})$ has conical points and/or conical/cylindrical ends. Here one gains the variational structure, obtaining existence with exponential $W^{2,2}$ -decay along the ends.

For the p -Laplacian see [Serrin, '64], [Kichenassamy-Veron, '86]: in this case one has homogeneity of the operator and the maximum principle.

Uniqueness of fundamental solutions $(\gamma_2 = 6\gamma_3)$

Uniqueness of fundamental solutions $(\gamma_2 = 6\gamma_3)$

In [Chang-Yang,'95] it was shown by Bochner's identity and $\gamma_2 = 6\gamma_3$ that (the differential part of) F_L is convex on $W^{2,2}(M)$.

Uniqueness of fundamental solutions $(\gamma_2 = 6\gamma_3)$

In [Chang-Yang,'95] it was shown by Bochner's identity and $\gamma_2 = 6\gamma_3$ that (the differential part of) F_L is convex on $W^{2,2}(M)$. This gives uniqueness of solutions for $\mathcal{N}(u) = f$ in $W^{2,2}(M)$, but $u_s \notin W^{2,2}(M)$.

Uniqueness of fundamental solutions $(\gamma_2 = 6\gamma_3)$

In [Chang-Yang, '95] it was shown by Bochner's identity and $\gamma_2 = 6\gamma_3$ that (the differential part of) F_L is convex on $W^{2,2}(M)$. This gives uniqueness of solutions for $\mathcal{N}(u) = f$ in $W^{2,2}(M)$, but $u_s \notin W^{2,2}(M)$.

In [Boccardo-Gallouët, '92] solutions to (2nd-order) PDEs with measure data were found as limits of solutions with mollified right-hand sides.

Uniqueness of fundamental solutions $(\gamma_2 = 6\gamma_3)$

In [Chang-Yang, '95] it was shown by Bochner's identity and $\gamma_2 = 6\gamma_3$ that (the differential part of) F_L is convex on $W^{2,2}(M)$. This gives uniqueness of solutions for $\mathcal{N}(u) = f$ in $W^{2,2}(M)$, but $u_s \notin W^{2,2}(M)$.

In [Boccardo-Gallouët, '92] solutions to (2nd-order) PDEs with measure data were found as limits of solutions with mollified right-hand sides.

The *grand L^p space* ([Iwaniec-Sbordone, '92]) are the functions u s.t.

$$\|u\|_{L^{\theta,p)} := \sup_{\varepsilon \in (0,1]} \varepsilon^{\frac{\theta}{p}} \|u\|_{L^{p(1-\varepsilon)}} < +\infty.$$

Uniqueness of fundamental solutions $(\gamma_2 = 6\gamma_3)$

In [Chang-Yang, '95] it was shown by Bochner's identity and $\gamma_2 = 6\gamma_3$ that (the differential part of) F_L is convex on $W^{2,2}(M)$. This gives uniqueness of solutions for $\mathcal{N}(u) = f$ in $W^{2,2}(M)$, but $u_s \notin W^{2,2}(M)$.

In [Boccardo-Gallouët, '92] solutions to (2nd-order) PDEs with measure data were found as limits of solutions with mollified right-hand sides.

The *grand L^p space* ([Iwaniec-Sbordone, '92]) are the functions u s.t.

$$\|u\|_{L^{\theta,p})} := \sup_{\varepsilon \in (0,1]} \varepsilon^{\frac{\theta}{p}} \|u\|_{L^{p(1-\varepsilon)}} < +\infty.$$

It satisfies $L^{2,\infty} \hookrightarrow L^{1,2})$ and used to study compensated compactness.

Uniqueness of fundamental solutions $(\gamma_2 = 6\gamma_3)$

In [Chang-Yang, '95] it was shown by Bochner's identity and $\gamma_2 = 6\gamma_3$ that (the differential part of) F_L is convex on $W^{2,2}(M)$. This gives uniqueness of solutions for $\mathcal{N}(u) = f$ in $W^{2,2}(M)$, but $u_s \notin W^{2,2}(M)$.

In [Boccardo-Gallouët, '92] solutions to (2nd-order) PDEs with measure data were found as limits of solutions with mollified right-hand sides.

The *grand L^p space* ([Iwaniec-Sbordone, '92]) are the functions u s.t.

$$\|u\|_{L^{\theta,p})} := \sup_{\varepsilon \in (0,1]} \varepsilon^{\frac{\theta}{p}} \|u\|_{L^{p(1-\varepsilon)}} < +\infty.$$

It satisfies $L^{2,\infty} \hookrightarrow L^{1,2})$ and used to study compensated compactness.

Using arguments in [Iwaniec, '92], [Iwaniec-Greco-Sbordone, '97] one can show that u_s coincides all every solutions both in $W^{1,2,2})$ and *mollifiable*.

Uniqueness of fundamental solutions $(\gamma_2 = 6\gamma_3)$

In [Chang-Yang, '95] it was shown by Bochner's identity and $\gamma_2 = 6\gamma_3$ that (the differential part of) F_L is convex on $W^{2,2}(M)$. This gives uniqueness of solutions for $\mathcal{N}(u) = f$ in $W^{2,2}(M)$, but $u_s \notin W^{2,2}(M)$.

In [Boccardo-Gallouët, '92] solutions to (2nd-order) PDEs with measure data were found as limits of solutions with mollified right-hand sides.

The *grand L^p space* ([Iwaniec-Sbordone, '92]) are the functions u s.t.

$$\|u\|_{L^{\theta,p)} := \sup_{\varepsilon \in (0,1]} \varepsilon^{\frac{\theta}{p}} \|u\|_{L^{p(1-\varepsilon)}} < +\infty.$$

It satisfies $L^{2,\infty} \hookrightarrow L^{1,2)}$ and used to study compensated compactness.

Using arguments in [Iwaniec, '92], [Iwaniec-Greco-Sbordone, '97] one can show that u_s coincides all every solutions both in $W^{1,2,2)}$ and *mollifiable*.

- The argument works for any (finite) measure data.

Proof of Theorem A

Proof of Theorem A

Let u_n solve $\mathcal{N}_L(u_n) + U_n = \mu_n e^{4u_n}$.

Proof of Theorem A

Let u_n solve $\mathcal{N}_L(u_n) + U_n = \mu_n e^{4u_n}$. We saw that at each blow-up point must accumulate at least ε_0 in conformal volume.

Proof of Theorem A

Let u_n solve $\mathcal{N}_L(u_n) + U_n = \mu_n e^{4u_n}$. We saw that at each blow-up point must accumulate at least ε_0 in conformal volume. Hence we have that

$$\mu_n e^{4u_n} \rightharpoonup \sum_{i=1}^l \beta_i \delta_{p_i} + h; \quad \beta_i \geq \varepsilon_0,$$

where $h = \mu e^{4u_\infty}$ is the continuous part of the limit measure (smooth).

Proof of Theorem A

Let u_n solve $\mathcal{N}_L(u_n) + U_n = \mu_n e^{4u_n}$. We saw that at each blow-up point must accumulate at least ε_0 in conformal volume. Hence we have that

$$\mu_n e^{4u_n} \rightharpoonup \sum_{i=1}^l \beta_i \delta_{p_i} + h; \quad \beta_i \geq \varepsilon_0,$$

where $h = \mu e^{4u_\infty}$ is the continuous part of the limit measure (smooth).

Step 1. Using a Pohozaev identity and the above uniqueness property it is possible to show that $\beta_i \geq 8\pi^2 \gamma_2$ ($= \int_{S^4} U_{S^4} dv$).

Proof of Theorem A

Let u_n solve $\mathcal{N}_L(u_n) + U_n = \mu_n e^{4u_n}$. We saw that at each blow-up point must accumulate at least ε_0 in conformal volume. Hence we have that

$$\mu_n e^{4u_n} \rightharpoonup \sum_{i=1}^l \beta_i \delta_{p_i} + h; \quad \beta_i \geq \varepsilon_0,$$

where $h = \mu e^{4u_\infty}$ is the continuous part of the limit measure (smooth).

Step 1. Using a Pohozaev identity and the above uniqueness property it is possible to show that $\beta_i \geq 8\pi^2 \gamma_2$ ($= \int_{S^4} U_{S^4} dv$).

Step 2. From the uniqueness of fundamental solutions, one finds that $\lim_n u_n \simeq \alpha_i \log d(x, p_i)$ near p_i , with $\alpha_i \leq -2$.

Proof of Theorem A

Let u_n solve $\mathcal{N}_L(u_n) + U_n = \mu_n e^{4u_n}$. We saw that at each blow-up point must accumulate at least ε_0 in conformal volume. Hence we have that

$$\mu_n e^{4u_n} \rightharpoonup \sum_{i=1}^l \beta_i \delta_{p_i} + h; \quad \beta_i \geq \varepsilon_0,$$

where $h = \mu e^{4u_\infty}$ is the continuous part of the limit measure (smooth).

Step 1. Using a Pohozaev identity and the above uniqueness property it is possible to show that $\beta_i \geq 8\pi^2 \gamma_2$ ($= \int_{S^4} U_{S^4} dv$).

Step 2. From the uniqueness of fundamental solutions, one finds that $\lim_n u_n \simeq \alpha_i \log d(x, p_i)$ near p_i , with $\alpha_i \leq -2$. If the weak limit u_∞ is non zero, the conformal volume would diverge.

Proof of Theorem A

Let u_n solve $\mathcal{N}_L(u_n) + U_n = \mu_n e^{4u_n}$. We saw that at each blow-up point must accumulate at least ε_0 in conformal volume. Hence we have that

$$\mu_n e^{4u_n} \rightharpoonup \sum_{i=1}^l \beta_i \delta_{p_i} + h; \quad \beta_i \geq \varepsilon_0,$$

where $h = \mu e^{4u_\infty}$ is the continuous part of the limit measure (smooth).

Step 1. Using a Pohozaev identity and the above uniqueness property it is possible to show that $\beta_i \geq 8\pi^2 \gamma_2$ ($= \int_{S^4} U_{S^4} dv$).

Step 2. From the uniqueness of fundamental solutions, one finds that $\lim_n u_n \simeq \alpha_i \log d(x, p_i)$ near p_i , with $\alpha_i \leq -2$. If the weak limit u_∞ is non zero, the conformal volume would diverge. So $h \equiv 0$.

Proof of Theorem A

Let u_n solve $\mathcal{N}_L(u_n) + U_n = \mu_n e^{4u_n}$. We saw that at each blow-up point must accumulate at least ε_0 in conformal volume. Hence we have that

$$\mu_n e^{4u_n} \rightharpoonup \sum_{i=1}^l \beta_i \delta_{p_i} + h; \quad \beta_i \geq \varepsilon_0,$$

where $h = \mu e^{4u_\infty}$ is the continuous part of the limit measure (smooth).

Step 1. Using a Pohozaev identity and the above uniqueness property it is possible to show that $\beta_i \geq 8\pi^2 \gamma_2$ ($= \int_{S^4} U_{S^4} dv$).

Step 2. From the uniqueness of fundamental solutions, one finds that $\lim_n u_n \simeq \alpha_i \log d(x, p_i)$ near p_i , with $\alpha_i \leq -2$. If the weak limit u_∞ is non zero, the conformal volume would diverge. So $h \equiv 0$.

Step 3. Use Pohozaev's identity again to show that $\beta_i = 8\pi^2 \gamma_2$ for all i .

Proof of Theorem A

Let u_n solve $\mathcal{N}_L(u_n) + U_n = \mu_n e^{4u_n}$. We saw that at each blow-up point must accumulate at least ε_0 in conformal volume. Hence we have that

$$\mu_n e^{4u_n} \rightharpoonup \sum_{i=1}^l \beta_i \delta_{p_i} + h; \quad \beta_i \geq \varepsilon_0,$$

where $h = \mu e^{4u_\infty}$ is the continuous part of the limit measure (smooth).

Step 1. Using a Pohozaev identity and the above uniqueness property it is possible to show that $\beta_i \geq 8\pi^2 \gamma_2$ ($= \int_{S^4} U_{S^4} dv$).

Step 2. From the uniqueness of fundamental solutions, one finds that $\lim_n u_n \simeq \alpha_i \log d(x, p_i)$ near p_i , with $\alpha_i \leq -2$. If the weak limit u_∞ is non zero, the conformal volume would diverge. So $h \equiv 0$.

Step 3. Use Pohozaev's identity again to show that $\beta_i = 8\pi^2 \gamma_2$ for all i .

• For general coefficients, it would be enough to know the uniqueness of the *singular profile* of u_s , without knowing global uniqueness.

Existence of extremal metrics for F_L

Existence of extremal metrics for F_L

Via min-max theory, we then obtain the following result.

Existence of extremal metrics for F_L

Via min-max theory, we then obtain the following result.

Theorem B

Assume $\gamma_2 = 6\gamma_3 \neq 0$. Suppose (M^4, g) satisfies $\int_M U_g dv \notin 8\pi^2\gamma_2\mathbb{N}$. Then there exists an extremal metric.

Existence of extremal metrics for F_L

Via min-max theory, we then obtain the following result.

Theorem B

Assume $\gamma_2 = 6\gamma_3 \neq 0$. Suppose (M^4, g) satisfies $\int_M U_g dv \notin 8\pi^2\gamma_2\mathbb{N}$. Then there exists an extremal metric.

Open problem. Understand the cases $\int_M U_g dv \in 8\pi^2\gamma_2\mathbb{N}$.

Existence of extremal metrics for F_L

Via min-max theory, we then obtain the following result.

Theorem B

Assume $\gamma_2 = 6\gamma_3 \neq 0$. Suppose (M^4, g) satisfies $\int_M U_g dv \notin 8\pi^2\gamma_2\mathbb{N}$. Then there exists an extremal metric.

Open problem. Understand the cases $\int_M U_g dv \in 8\pi^2\gamma_2\mathbb{N}$. Here the concentration/compactness dichotomy fails, and the determinant functional has *asymptotes*.

Existence of extremal metrics for F_L

Via min-max theory, we then obtain the following result.

Theorem B

Assume $\gamma_2 = 6\gamma_3 \neq 0$. Suppose (M^4, g) satisfies $\int_M U_g dv \notin 8\pi^2\gamma_2\mathbb{N}$. Then there exists an extremal metric.

Open problem. Understand the cases $\int_M U_g dv \in 8\pi^2\gamma_2\mathbb{N}$. Here the concentration/compactness dichotomy fails, and the determinant functional has *asymptotes*.

One could try to understand them defining and studying a suitable *mass* for the blown-up manifold via the fundamental solution.

The determinant of the Paneitz operator

The determinant of the Paneitz operator

It is mentioned in Connes' book on non-commutative geometry as a relevant tool for conformal theories in 4D.

The determinant of the Paneitz operator

It is mentioned in Connes' book on non-commutative geometry as a relevant tool for conformal theories in 4D. Analytically, it is also quite interesting.

The determinant of the Paneitz operator

It is mentioned in Connes' book on non-commutative geometry as a relevant tool for conformal theories in 4D. Analytically, it is also quite interesting.

In flat tori, the determinant of P_g is

$$F_P[w] = \int_{\mathbb{T}^4} [18(\Delta w)^2 + 64|\nabla w|^2 \Delta w + 32|\nabla w|^4] dx.$$

The determinant of the Paneitz operator

It is mentioned in Connes' book on non-commutative geometry as a relevant tool for conformal theories in 4D. Analytically, it is also quite interesting.

In flat tori, the determinant of P_g is

$$F_P[w] = \int_{\mathbb{T}^4} [18(\Delta w)^2 + 64|\nabla w|^2 \Delta w + 32|\nabla w|^4] dx.$$

This functional has a triple homogeneity and is again doubly critical.

The determinant of the Paneitz operator

It is mentioned in Connes' book on non-commutative geometry as a relevant tool for conformal theories in 4D. Analytically, it is also quite interesting.

In flat tori, the determinant of P_g is

$$F_P[w] = \int_{\mathbb{T}^4} [18(\Delta w)^2 + 64|\nabla w|^2 \Delta w + 32|\nabla w|^4] dx.$$

This functional has a triple homogeneity and is again doubly critical.

On S^4 instead one has

$$\begin{aligned} F_P[w] &= \int_{S^4} [18(\Delta w)^2 + 64|\nabla w|^2 \Delta w + 32|\nabla w|^4 - 60|\nabla w|^2] dv \\ &+ 112\pi^2 \log \left(\int_{S^4} e^{4(w-\bar{w})} dv \right). \end{aligned}$$

Mountain Pass structure

Proposition 2

Mountain Pass structure

Proposition 2 For both $\mathbb{T}^4$ and S^4 , F_P has a local minimum at $w \equiv 0$ (standard metrics).

Mountain Pass structure

Proposition 2 For both $\mathbb{T}^4$ and S^4 , F_P has a local minimum at $w \equiv 0$ (standard metrics). Moreover, F_P is unbounded above and below.

Mountain Pass structure

Proposition 2 For both $\mathbb{T}^4$ and S^4 , F_P has a local minimum at $w \equiv 0$ (standard metrics). Moreover, F_P is unbounded above and below.

The local minimality at $w = 0$ was noticed in [Branson, '96], computing the second variation.

Mountain Pass structure

Proposition 2 For both $\mathbb{T}^4$ and S^4 , F_P has a local minimum at $w \equiv 0$ (standard metrics). Moreover, F_P is unbounded above and below.

The local minimality at $w = 0$ was noticed in [Branson, '96], computing the second variation. To check unboundedness from below, insert into F_P the function

Mountain Pass structure

Proposition 2 For both $\mathbb{T}^4$ and S^4 , F_P has a local minimum at $w \equiv 0$ (standard metrics). Moreover, F_P is unbounded above and below.

The local minimality at $w = 0$ was noticed in [Branson, '96], computing the second variation. To check unboundedness from below, insert into F_P the function

$$w(x) \simeq -\frac{1}{2} \log(\varepsilon^2 + |x|^2); \quad \varepsilon \rightarrow 0.$$

Proposition 2 For both $\mathbb{T}^4$ and S^4 , F_P has a local minimum at $w \equiv 0$ (standard metrics). Moreover, F_P is unbounded above and below.

The local minimality at $w = 0$ was noticed in [Branson, '96], computing the second variation. To check unboundedness from below, insert into F_P the function

$$w(x) \simeq -\frac{1}{2} \log(\varepsilon^2 + |x|^2); \quad \varepsilon \rightarrow 0.$$

- Geometrically, this conformal factor generates a *cigar* (not a bubble).

Proposition 2 For both $\mathbb{T}^4$ and S^4 , F_P has a local minimum at $w \equiv 0$ (standard metrics). Moreover, F_P is unbounded above and below.

The local minimality at $w = 0$ was noticed in [Branson, '96], computing the second variation. To check unboundedness from below, insert into F_P the function

$$w(x) \simeq -\frac{1}{2} \log(\varepsilon^2 + |x|^2); \quad \varepsilon \rightarrow 0.$$

- Geometrically, this conformal factor generates a *cigar* (not a bubble).
- Loss of coercivity may happen in *different ways* (e.g., at many points)

Proposition 2 For both $\mathbb{T}^4$ and S^4 , F_P has a local minimum at $w \equiv 0$ (standard metrics). Moreover, F_P is unbounded above and below.

The local minimality at $w = 0$ was noticed in [Branson, '96], computing the second variation. To check unboundedness from below, insert into F_P the function

$$w(x) \simeq -\frac{1}{2} \log(\varepsilon^2 + |x|^2); \quad \varepsilon \rightarrow 0.$$

- Geometrically, this conformal factor generates a *cigar* (not a bubble).
- Loss of coercivity may happen in *different ways* (e.g., at many points), differently e.g. from the Q-curvature equation.

Mountain Pass structure

Proposition 2 For both $\mathbb{T}^4$ and S^4 , F_P has a local minimum at $w \equiv 0$ (standard metrics). Moreover, F_P is unbounded above and below.

The local minimality at $w = 0$ was noticed in [Branson, '96], computing the second variation. To check unboundedness from below, insert into F_P the function

$$w(x) \simeq -\frac{1}{2} \log(\varepsilon^2 + |x|^2); \quad \varepsilon \rightarrow 0.$$

- Geometrically, this conformal factor generates a *cigar* (not a bubble).
- Loss of coercivity may happen in *different ways* (e.g., at many points), differently e.g. from the Q-curvature equation.
- It goes similarly for compact hyperbolic manifolds.

A second solution on S^4

A second solution on S^4

Theorem C ([Gursky-M., '12])

Let (S^4, g_0) be the standard 4-sphere. Then F_P admits a non-trivial axially symmetric solution.

A second solution on S^4

Theorem C ([Gursky-M., '12])

Let (S^4, g_0) be the standard 4-sphere. Then F_P admits a non-trivial axially symmetric solution.

Remarks (a) For most geometric problems the round metric is *the only critical point*. One has indeed uniqueness of the round metric for constant mean curvature, Gaussian curvature, scalar curvature and Q-curvature.

A second solution on S^4

Theorem C ([Gursky-M., '12])

Let (S^4, g_0) be the standard 4-sphere. Then F_P admits a non-trivial axially symmetric solution.

Remarks (a) For most geometric problems the round metric is *the only critical point*. One has indeed uniqueness of the round metric for constant mean curvature, Gaussian curvature, scalar curvature and Q-curvature.

(b) Uniqueness also holds for critical points of $\det L_g$ ([Gursky, '97]).

A second solution on S^4

Theorem C ([Gursky-M., '12])

Let (S^4, g_0) be the standard 4-sphere. Then F_P admits a non-trivial axially symmetric solution.

Remarks (a) For most geometric problems the round metric is *the only critical point*. One has indeed uniqueness of the round metric for constant mean curvature, Gaussian curvature, scalar curvature and Q-curvature.

(b) Uniqueness also holds for critical points of $\det L_g$ ([Gursky, '97]). From the positive second variation at $w = 0$, Branson speculated uniqueness for critical points of F_P as well (false).

A second solution on S^4

Theorem C ([Gursky-M., '12])

Let (S^4, g_0) be the standard 4-sphere. Then F_P admits a non-trivial axially symmetric solution.

Remarks (a) For most geometric problems the round metric is *the only critical point*. One has indeed uniqueness of the round metric for constant mean curvature, Gaussian curvature, scalar curvature and Q-curvature.

(b) Uniqueness also holds for critical points of $\det L_g$ ([Gursky, '97]). From the positive second variation at $w = 0$, Branson speculated uniqueness for critical points of F_P as well (false).

(c) The mountain pass structure suggests to use a variational approach.

A second solution on S^4

Theorem C ([Gursky-M., '12])

Let (S^4, g_0) be the standard 4-sphere. Then F_P admits a non-trivial axially symmetric solution.

Remarks (a) For most geometric problems the round metric is *the only critical point*. One has indeed uniqueness of the round metric for constant mean curvature, Gaussian curvature, scalar curvature and Q-curvature.

(b) Uniqueness also holds for critical points of $\det L_g$ ([Gursky, '97]). From the positive second variation at $w = 0$, Branson speculated uniqueness for critical points of F_P as well (false).

(c) The mountain pass structure suggests to use a variational approach. However this strategy is now out of reach: we used ODEs instead.

A second solution on S^4

Theorem C ([Gursky-M., '12])

Let (S^4, g_0) be the standard 4-sphere. Then F_P admits a non-trivial axially symmetric solution.

Remarks (a) For most geometric problems the round metric is *the only critical point*. One has indeed uniqueness of the round metric for constant mean curvature, Gaussian curvature, scalar curvature and Q-curvature.

(b) Uniqueness also holds for critical points of $\det L_g$ ([Gursky, '97]). From the positive second variation at $w = 0$, Branson speculated uniqueness for critical points of F_P as well (false).

(c) The mountain pass structure suggests to use a variational approach. However this strategy is now out of reach: we used ODEs instead.

(d) A similar result holds in $\mathbb{R}^4$, much easier to prove.

Comments and open problems

Comments and open problems

Our proof is very specific and does not exploit the structure of the problem.

Comments and open problems

Our proof is very specific and does not exploit the structure of the problem. Recall that in $\mathbb{T}^4$ the determinant is

$$F_P[w] = \int_{\mathbb{T}^4} [18(\Delta w)^2 + 64|\nabla w|^2 \Delta w + 32|\nabla w|^4] \, dx.$$

Comments and open problems

Our proof is very specific and does not exploit the structure of the problem. Recall that in $\mathbb{T}^4$ the determinant is

$$F_P[w] = \int_{\mathbb{T}^4} [18(\Delta w)^2 + 64|\nabla w|^2\Delta w + 32|\nabla w|^4] dx.$$

It is difficult to find a priori bounds on solutions or P-S sequences

Comments and open problems

Our proof is very specific and does not exploit the structure of the problem. Recall that in $\mathbb{T}^4$ the determinant is

$$F_P[w] = \int_{\mathbb{T}^4} [18(\Delta w)^2 + 64|\nabla w|^2\Delta w + 32|\nabla w|^4] dx.$$

It is difficult to find a priori bounds on solutions or P-S sequences.

Notice that by Bochner's identity $\int_{\mathbb{T}^4} (\Delta u)^2 dx = \int_{\mathbb{T}^4} |\nabla^2 u|^2 dx$

Comments and open problems

Our proof is very specific and does not exploit the structure of the problem. Recall that in $\mathbb{T}^4$ the determinant is

$$F_P[w] = \int_{\mathbb{T}^4} [18(\Delta w)^2 + 64|\nabla w|^2 \Delta w + 32|\nabla w|^4] dx.$$

It is difficult to find a priori bounds on solutions or P-S sequences.

Notice that by Bochner's identity $\int_{\mathbb{T}^4} (\Delta u)^2 dx = \int_{\mathbb{T}^4} |\nabla^2 u|^2 dx$, so there is a positive lower bound for the Sobolev-type quotient

$$\inf_{u \neq 0} \frac{\int_{\mathbb{T}^4} (\Delta u)^2 dx}{\left(\int_{\mathbb{T}^4} |\nabla u|^4 dx \right)^{\frac{1}{2}}}.$$

Comments and open problems

Our proof is very specific and does not exploit the structure of the problem. Recall that in $\mathbb{T}^4$ the determinant is

$$F_P[w] = \int_{\mathbb{T}^4} [18(\Delta w)^2 + 64|\nabla w|^2 \Delta w + 32|\nabla w|^4] dx.$$

It is difficult to find a priori bounds on solutions or P-S sequences.

Notice that by Bochner's identity $\int_{\mathbb{T}^4} (\Delta u)^2 dx = \int_{\mathbb{T}^4} |\nabla^2 u|^2 dx$, so there is a positive lower bound for the Sobolev-type quotient

$$\inf_{u \neq 0} \frac{\int_{\mathbb{T}^4} (\Delta u)^2 dx}{\left(\int_{\mathbb{T}^4} |\nabla u|^4 dx \right)^{\frac{1}{2}}}.$$

It is an interesting question to characterize extremals of this quotient in $\mathbb{R}^4$, vaguely related to the above problem.

The Euler equation

The Euler equation

On $\mathbb{R}^4$ critical points satisfy

$$9\Delta^2 w + 32|\nabla^2 w|^2 - 32(\Delta w)^2 - 32\Delta u |\nabla u|^2 - 32\langle \nabla w, \nabla |\nabla w|^2 \rangle = 0.$$

The Euler equation

On $\mathbb{R}^4$ critical points satisfy

$$9\Delta^2 w + 32|\nabla^2 w|^2 - 32(\Delta w)^2 - 32\Delta u |\nabla u|^2 - 32\langle \nabla w, \nabla |\nabla w|^2 \rangle = 0.$$

The main-order term is Δ^2 : typically, decay of solutions is logarithmic.

The Euler equation

On $\mathbb{R}^4$ critical points satisfy

$$9\Delta^2 w + 32|\nabla^2 w|^2 - 32(\Delta w)^2 - 32\Delta u |\nabla u|^2 - 32\langle \nabla w, \nabla |\nabla w|^2 \rangle = 0.$$

The main-order term is Δ^2 : typically, decay of solutions is logarithmic. However solutions with finite energy have inverse-quadratic decay: some degeneracy is present.

The Euler equation

On $\mathbb{R}^4$ critical points satisfy

$$9\Delta^2 w + 32|\nabla^2 w|^2 - 32(\Delta w)^2 - 32\Delta u |\nabla u|^2 - 32\langle \nabla w, \nabla |\nabla w|^2 \rangle = 0.$$

The main-order term is Δ^2 : typically, decay of solutions is logarithmic. However solutions with finite energy have inverse-quadratic decay: some degeneracy is present.

Apart from the compactness issues, new sharp Moser-Trudinger inequalities would be expected.

Thanks for your attention