

Comments on Alfven propulsion (Steve Adler, IAS)

Conversion of the Drell, Foley, Ruderman formulas to SI notation:

Their equations (6), (8), (9), (10) for  $V$ ,  $P$ ,  $I$ , and  $Z$  become respectively, with  $\mu_0$  the vacuum permeability:

$$V = v_c B_0 M \quad (6)$$

$$P = \frac{ML}{\mu_0} \frac{v_c^2}{v_a} B_0^2 \quad (8)$$

$$I = \frac{L}{\mu_0} \frac{v_c}{v_a} B_0 \quad (9)$$

$$Z = \frac{M}{L} v_a \mu_0 \quad (10)$$

I have corrected an apparent factor of 2 error in their Eqs. (8-10), since I think that the energy density should be  $h^2/8\pi$  rather than  $h^2/4\pi$ , so I replaced the  $h^2/4\pi$  in their (8) by  $h^2/2\mu_0$ . The SI formula of Eq. (10) above reproduces their  $Z = .23\Omega$  in the example on page 174 of their article, since with  $M \sim L$ ,  $v_a = 2 \times 10^5 \text{m/s}$ , and  $\mu_0 = 1.256 \times 10^{-6} \text{mkg}/(\text{s}^2 \text{A}^2)$ , one gets  $Z = 0.25 \text{m}^2 \text{kg}/(\text{s}^3 \text{A}^2) = 0.25\Omega$  (with  $\Omega = \text{ohm}$ ).

For the NEAR satellite solar panels, the ratio  $\frac{M}{L}$  will be smaller than 1 if  $M$  is the panel width, and much smaller than 1 if  $M$  is the panel thickness. Since the altitude is higher than 1600m, and since the Alfven velocity  $v_a$  increases with decreasing ionic density (see the discussion following their Eq. (4)), one will have  $v_a > 10^7 \text{m/s}$ ; scaling from their example, one has

$$Z = \frac{M}{L} \frac{v_a}{(2 \times 10^5 \text{m/s})} \times 0.25\Omega \quad ;$$

since  $v_a < c = 3 \times 10^8 \text{m/s}$ , this gives

$$12.5 \frac{M}{L} \Omega < Z < 375 \frac{M}{L} \Omega \quad (I)$$

For a solar cell of voltage  $V$ , internal series resistance  $R_S$  and insulation resistance  $R_I$  between the cell elements and the plasma, the power that can be developed in Alfven propulsion when  $N$  cells in a square array are wired in parallel will be

$$P = \frac{V^2}{(Z/N^{1/2} + (R_S + R_I)/N)} \quad (II)$$

Here  $Z$  is the impedance from Eq. (10) evaluated with the dimension  $L$  of the individual cell; since a square array will have side length  $LN^{1/2}$ , the  $Z$  value for the array will be  $Z/N^{1/2}$ . Similarly, for  $N$  cells wired in parallel, the resistances  $R_S$  and  $R_I$  of an individual cell are divided by  $N$  in an array. Note that Eq. (II) has the correct invariance property if each cell is imagined to be cut into four sub-cells wired in parallel: the cell side is decreased by a factor of 2, and so  $Z$  increases by a factor of 2, but since  $N$  is increased by a factor of 4, this increase in  $Z$  is compensated by the factor  $1/N^{1/2}$ . Similarly, the resistances of each subcell are increased by a factor of 4 (one expects the series resistance and the insulation resistance, per unit area, to be constant), but this is compensated by the factor  $1/N$ .

To make a rough estimate, from advertisements on the internet I find typical polycrystalline solar cell parameters  $V \sim 7 \text{volt}$ ,  $R_S \sim 150\Omega$ ,  $L \sim 5 \text{cm}$ , and from pictures it looks like  $M/L \sim 1/12.5$  as a rough order

of magnitude. Then Eq. (I) gives  $1\Omega < Z < 30\Omega$ . Taking the lowest  $Z$  value in this range, Eq. (II) gives an upper bound on  $P$ . Assuming that  $N \sim 10^4$ , this becomes

$$P < \frac{49}{(2.5 \times 10^{-2} + 10^{-4}R_I)} \text{watt} \quad , \quad (III)$$

with the insulation resistance  $R_I$  in ohms. So the possible power going into Alfvén propulsion depends crucially on the quality of the solar cell insulation resistance  $R_I$ . For  $R_I \sim 10^5$  ohms, one could get 5 watts of propulsion power, of the order needed for the flyby anomalies. But good glass insulation should be much better; for a  $R_I$  of 10 megohms, one gets a bound on the Alfvén propulsion power of 0.05 watt, too small to be of interest.