Standard Model (backgrounds) at the LHC

(part 2)

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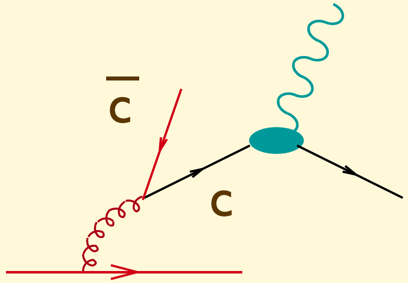
Exercise: charm in the proton



$$\frac{dc(x,Q)}{dt} = \frac{\alpha_s}{2\pi}\int_x^1 \frac{dy}{y} g(y,Q)\, P_{qg}(\frac{x}{y})$$

Assuming a typical behaviour of the gluon density: $g(x,Q) \sim A/x$

and using $P_{qg}(x) = \frac{1}{2}\left[x^2 + (1-x)^2\right]$ we get:

$$\frac{dc(x,Q)}{dt} = \frac{\alpha_s}{2\pi}\int_x^1 \frac{dy}{y} g(x/y,Q)\, P_{qg}(y) = \frac{\alpha_s}{2\pi}\int_x^1 dy \frac{A}{x}\frac{1}{2}\left[y^2 + (1-y)^2\right] = \frac{\alpha_s}{6\pi}\frac{A}{x}$$

and therefore:

$$c(x,Q) \sim \frac{\alpha_s}{6\pi} \log(\frac{Q^2}{m_c^2})\, g(x,Q)$$

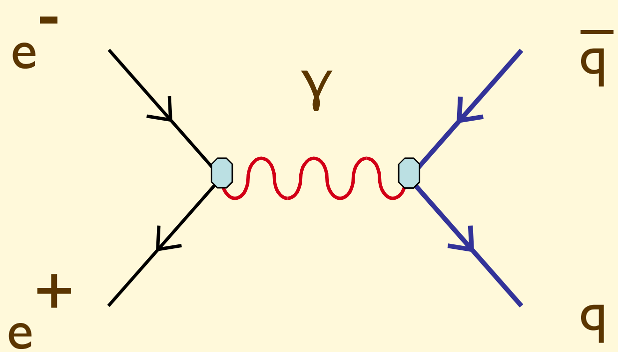
Corrections to this simple formula will arise due to the Q dependence of g(x) and of α_s

Evolution of hadronic final states

Asymptotic freedom implies that at E_{CM} >> 1 GeV

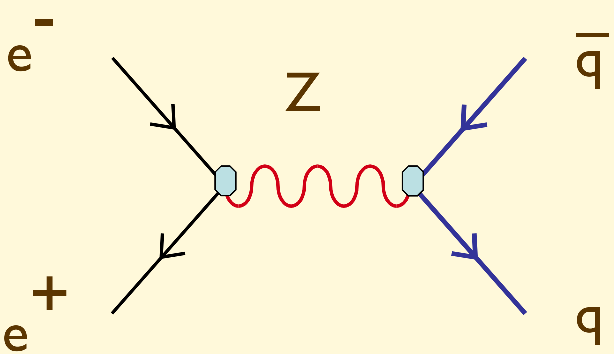
$\sigma(e^+ e^- \rightarrow \text{hadrons})$ ⟷ $\sigma(e^+ e^- \rightarrow \text{quarks/gluons})$

At the Leading Order (LO) in PT:



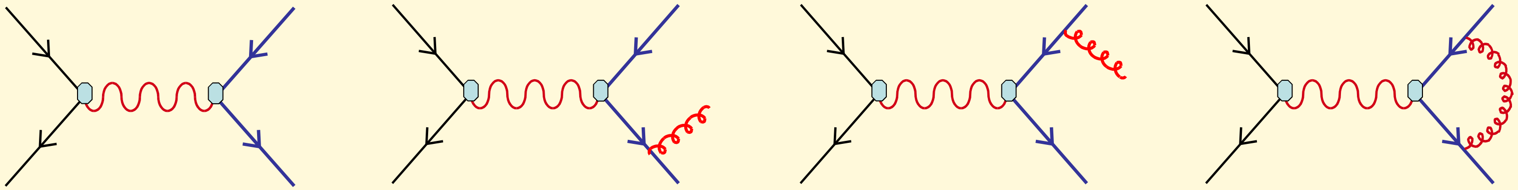
$$\sigma_0(e^+e^- \rightarrow qq) = \frac{4\pi\alpha^2}{9s} N_c \sum_{f=u,d,\ldots} e_{q_f}^2$$

$$\frac{\sigma_0(e^+e^- \rightarrow qq)}{\sigma_0(e^+e^- \rightarrow \mu^+\mu^-)} = N_c \sum_{f=u,d,\ldots} e_{q_f}^2$$



$$\frac{\sigma_0(e^+e^- \rightarrow Z \rightarrow qq)}{\sigma_0(e^+e^- \rightarrow Z \rightarrow \mu^+\mu^-)} = N_c \frac{\sum_{f=u,d,\ldots} \left(v_{q_f}^2 + a_{q_f}^2\right)}{\left(v_\mu^2 + a_\mu^2\right)}$$

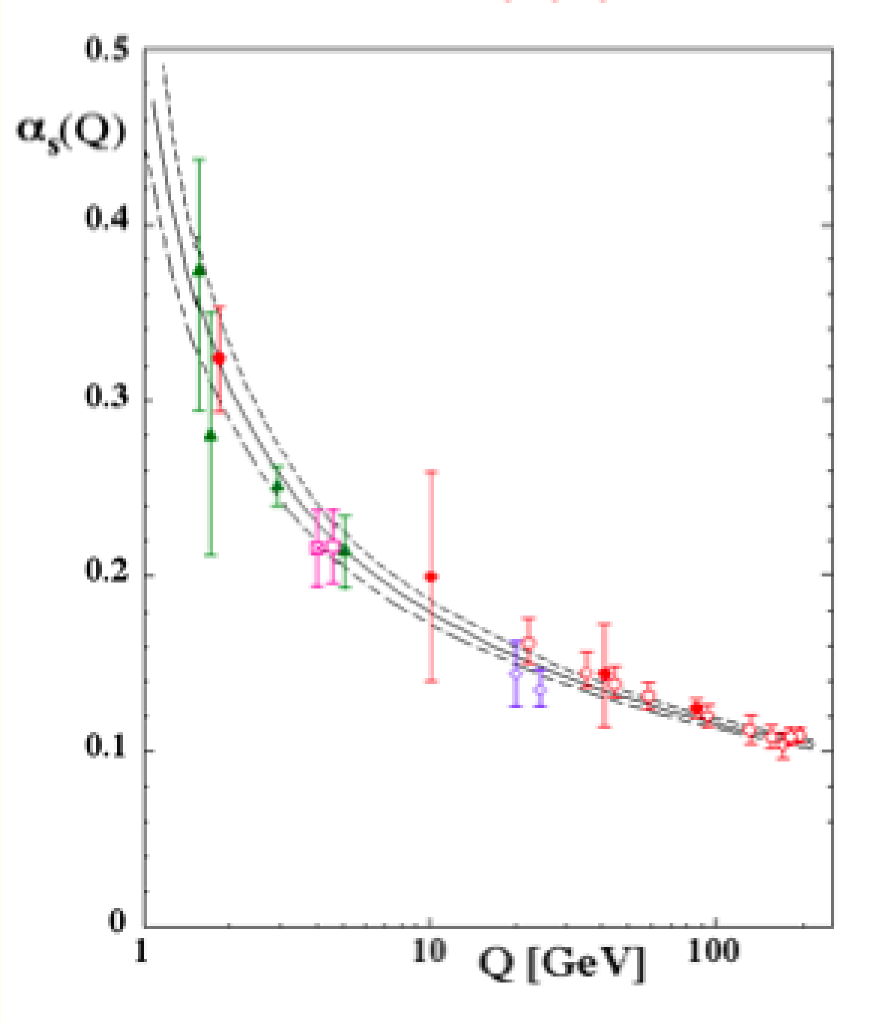
Adding higher-order perturbative terms:



+ ≥2-gluon emissions

$$\sigma_1(e^+e^- \to qq(g)) = \sigma_0(e^+e^- \to qq)\left(1 + \frac{\alpha_s(E_{CM})}{\pi} + O(\alpha_s^2)\right)$$

O(3%) at M_Z



Excellent agreement with data,
provided N_c=3
Extraction of αs consistent with the Q evolution predicted by QCD

To which extent the mapping

partons ⇔ particles

holds beyond the total cross section, e.g. for:

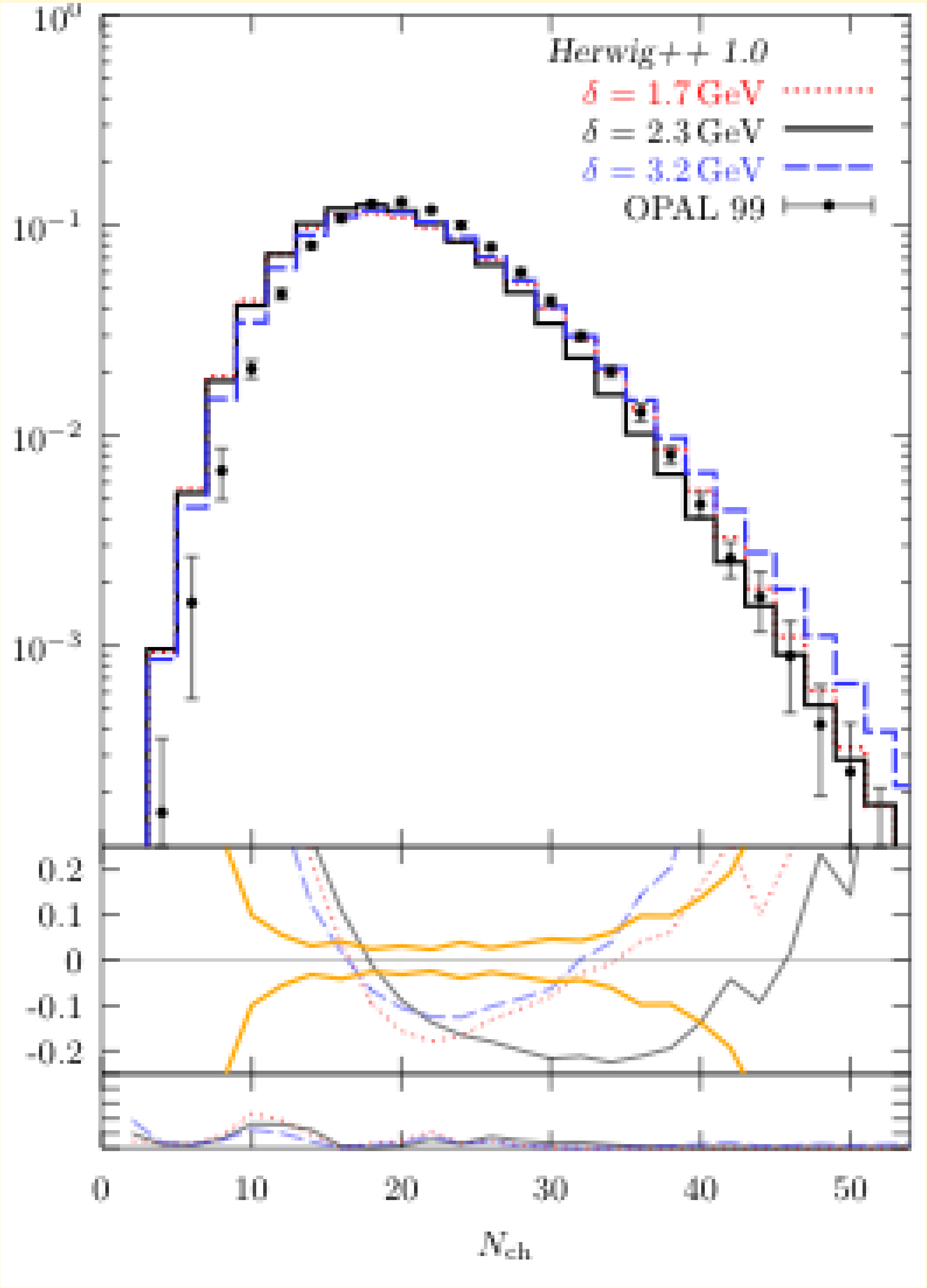
- momentum distributions
- angular distributions
- other kinematical correlations

?

This is crucial to establish the connection between the physical observables, and the fundamental couplings of the underlying theory (e.g. to extract SUSY couplings, or to verify the spin of new particles, which are detected through hadronic final states)

Experimentally, the final states contain a large number of particles, not the 2 or 3 which apparently saturate the perturbative cross-section.

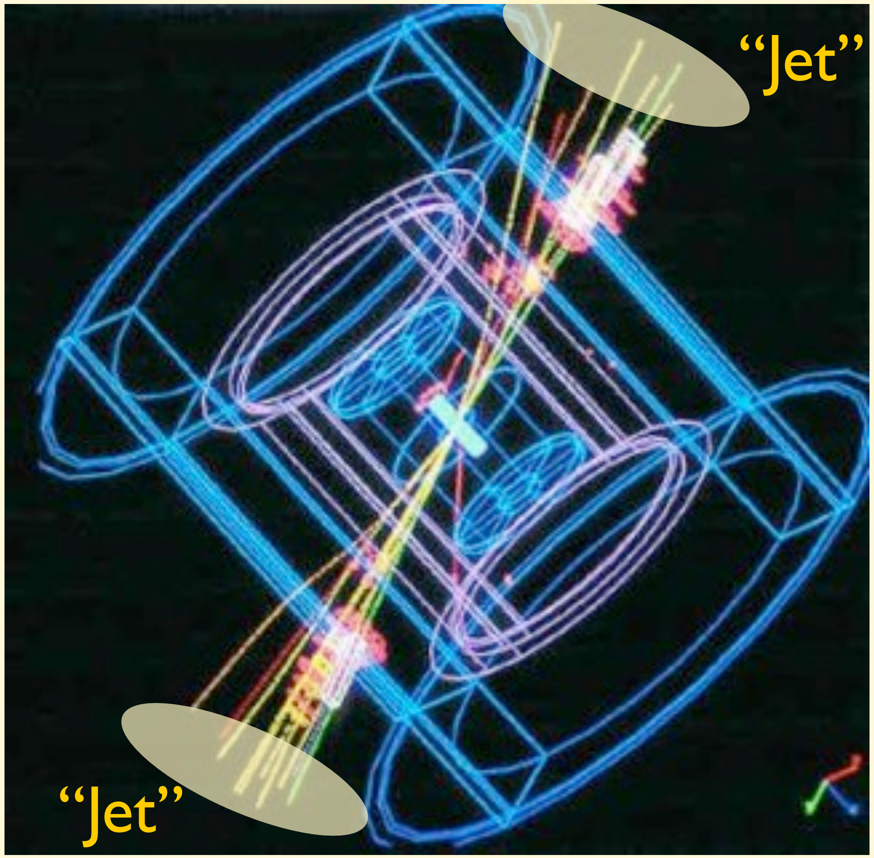
Experimental multiplicity distribution

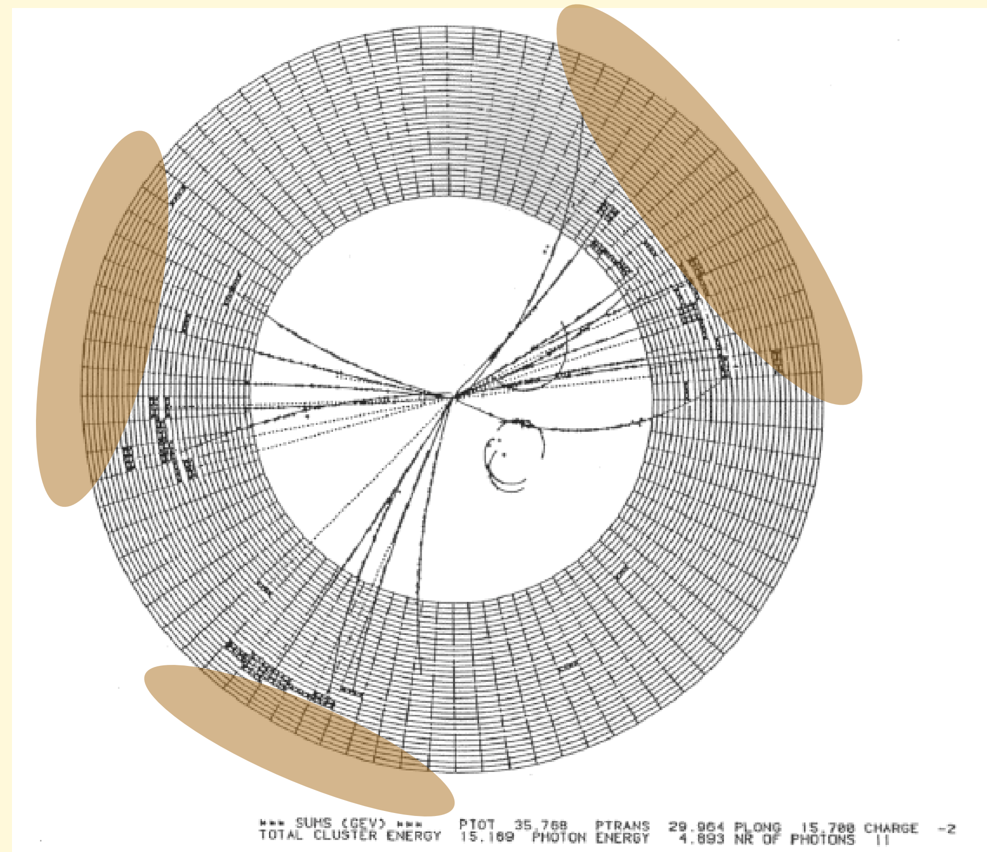


$\langle n_{charged} \rangle = 20.9$

This appears to jeopardize the hope to use particles to learn about the initial properties of partons

Look more closely at the structure of these events:



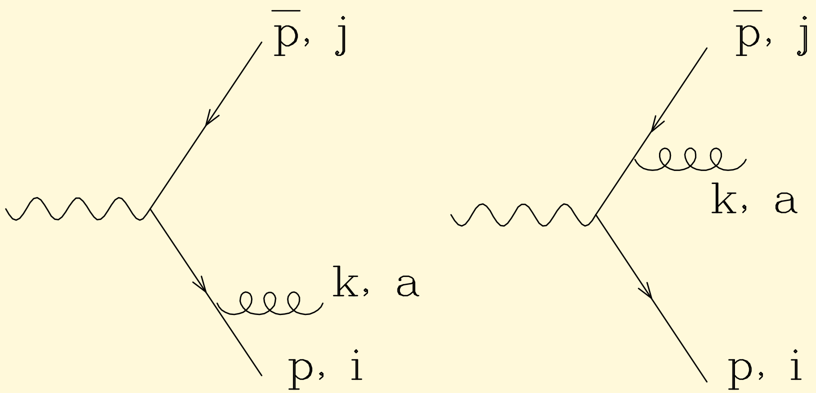


$e^+ e^- \rightarrow q\bar{q}$ ⇨ $e^+ e^- \rightarrow$ 2 jets

$e^+ e^- \rightarrow q\bar{q}g$ ⇨ $e^+ e^- \rightarrow$ 3 jets

The puzzle appears to be solved by associating partons to collimated "**jets**" of hadrons

Soft gluon emission



$$A = \bar{u}(p)\epsilon(k)(ig)\frac{-i}{\not{p}+\not{k}}\Gamma^\mu v(\bar{p})\ \lambda^a_{ij} + \bar{u}(p)\Gamma^\mu \frac{i}{\not{\bar{p}}+\not{k}}(ig)\epsilon(k)v(\bar{p})\ \lambda^a_{ij}$$

$$= \left[\frac{g}{2p\cdot k}\bar{u}(p)\epsilon(k)(\not{p}+\not{k})\Gamma^\mu v(\bar{p}) - \frac{g}{2\bar{p}\cdot k}\bar{u}(p)\Gamma^\mu(\not{\bar{p}}+\not{k})\epsilon(k)v(\bar{p})\right]\lambda^a_{ij}$$

$p\cdot k = p_0\, k_0\, (1-\cos\theta) \Rightarrow$ singularities for collinear ($\cos\theta \to 1$) or soft ($k_0 \to 0$) emission

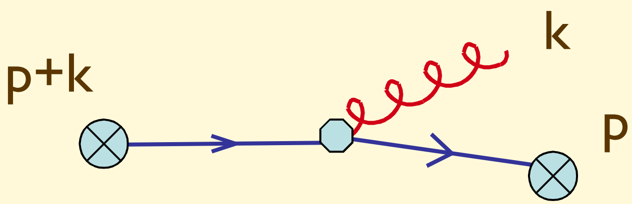
Collinear emission does not alter the global structure of the final state, since its preserves its "pencil-like-ness". **Soft emission** at large angle, however, could spoil the structure, and leads to strong interferences between emissions from different legs. So soft emission needs to be studied in more detail.

In the soft ($k_0 \to 0$) limit the amplitude simplifies and factorizes as follows:

$$A_{soft} = g\lambda^a_{ij}\left(\frac{p\cdot\varepsilon}{p\cdot k} - \frac{\bar{p}\cdot\varepsilon}{\bar{p}\cdot k}\right) A_{Born}$$

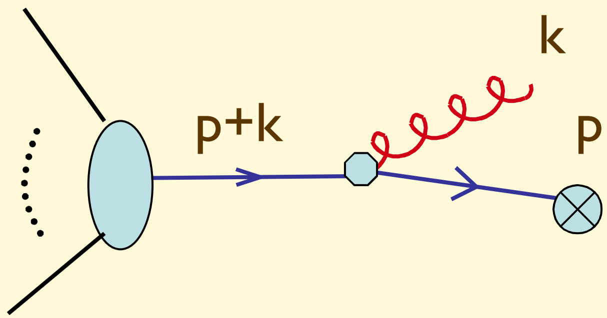
Factorization: it is the expression of the independence of long-wavelength (soft) emission on the nature of the hard (short-distance) process.

Another simple derivation of soft-gluon emission rules



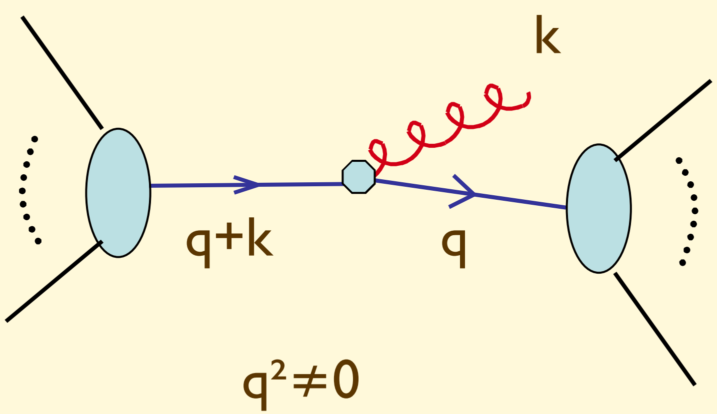
charge current of a free fermion

$$\psi(p)\gamma_\mu\psi(p+k)\,\varepsilon^\mu(k) \;\xrightarrow{k\to 0}\; \psi(p)\gamma_\mu\psi(p)\varepsilon^\mu(k) \;=\; 2p\cdot\varepsilon$$



$$\frac{1}{\not p+\not k}\gamma_\mu\psi(p)\,\varepsilon^\mu(k) \;\xrightarrow{k\to 0}$$

$$\frac{1}{2p\cdot k}\not p\gamma_\mu\psi(p)\varepsilon^\mu(k) \;=\; \frac{p\cdot\varepsilon}{p\cdot k}\psi(p)$$



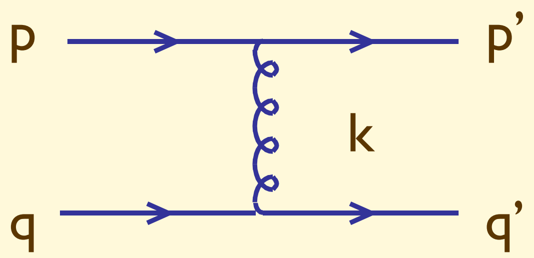
$$\frac{1}{\not q+\not k}\gamma_\mu\frac{1}{\not q}\varepsilon^\mu(k) \;\xrightarrow{q^2\neq 0,\,k\to 0}\; \frac{1}{q^2}\not q\gamma_\mu\not q\frac{1}{q^2}\varepsilon^\mu(k)$$

=> finite

<< Begin Exercise: Small-angle jet production, a useful approximation for the determination of the matrix elements and of the cross-section

At small scattering angle, $t = (p_1 - p_3)^2 \sim (1 - \cos\theta) \to 0$

and the $1/t^2$ propagators associated with t-channel gluon exchange dominate the matrix elements for all processes. In this limit it is easy to evaluate the matrix elements. For example:



$$\sim (\lambda^a)_{ij}(\lambda^a)_{kl}(2p_\mu)\frac{1}{t}(2q_\mu) = \frac{2s}{t}(\lambda^a)_{ij}(\lambda^a)_{kl}$$

where we used the fact that, for k=p-p'<<p (small angle scattering),

$$u(p')\gamma_\mu u(p) \sim u(p)\gamma_\mu u(p) = 2p_\mu$$

Using our colour algebra results, we then get: $\overline{\sum_{col,spin}} |M|^2 = \frac{1}{N_c^2}\frac{N_c^2-1}{4}\frac{4s^2}{t^2}$

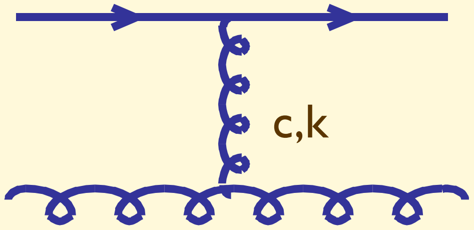
Noting that the result must be symmetric under s↔u exchange, and setting Nc=3, we finally obtain: $\overline{\sum_{col,spin}} |M|^2 = \frac{4}{9}\frac{s^2+u^2}{t^2}$

which turns out to be the exact result!

<< Continue Exercise>>

Quark-gluon and gluon-gluon scattering

We repeat the exercise in the more complex case of qg scattering, assuming the dominance of the t-channel gluon-exchange diagram:



$$\sim f^{abc}\lambda^c_{ij}\,2p_\mu\frac{1}{t}2q_\mu = 2\frac{s}{t}f^{abc}\lambda^c_{ij}$$

Using the colour algebra results, and enforcing the s↔u symmetry, we get:

$$\overline{\sum_{col,spin}}|M|^2 = \frac{s^2+u^2}{t^2}$$

which differs by only 20% from the exact result

$$\overline{\sum_{col,spin}}|M|^2 = \frac{s^2+u^2}{t^2} - \frac{4}{9}\frac{s^2+u^2}{us}$$

even in the large-angle region, at 90°

In a similar way we obtain for gg scattering (using the t↔u symmetry):

$$\overline{\sum_{col,spin}}|M(gg\to gg)|^2 = \frac{9}{2}\left(\frac{s^2}{t^2}+\frac{s^2}{u^2}\right)$$

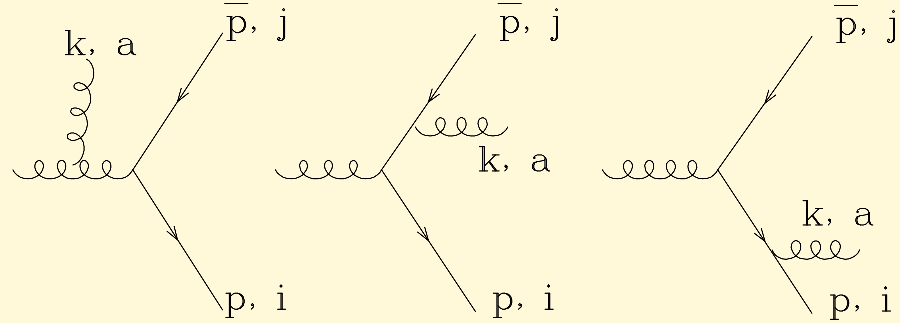
compared to the exact result

$$\overline{\sum_{col,spin}}|M(gg\to gg)|^2 = \frac{9}{2}\left(3-\frac{ut}{s^2}-\frac{us}{t^2}-\frac{st}{u^2}\right)$$

with a 20% difference at 90°

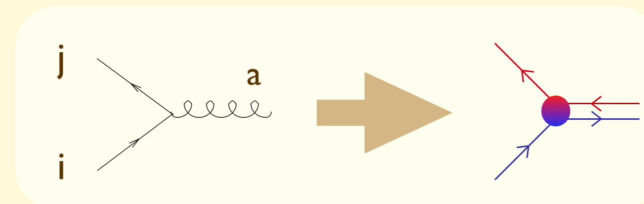
End Exercise >>

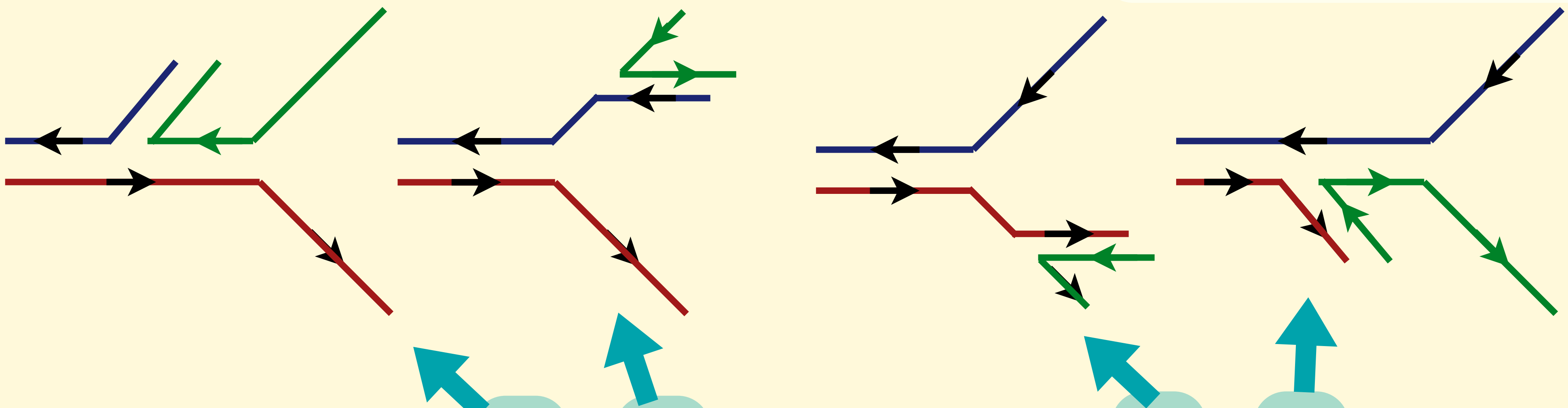
Similar, but more structured, result in the case of a fully coloured process:



$$A_{soft} = g\,(\lambda^a\lambda^b)_{ij}\left[\frac{Q\varepsilon}{Qk} - \frac{\bar{p}\varepsilon}{\bar{p}k}\right] + g\,(\lambda^b\lambda^a)_{ij}\left[\frac{p\varepsilon}{pk} - \frac{Q\varepsilon}{Qk}\right]$$

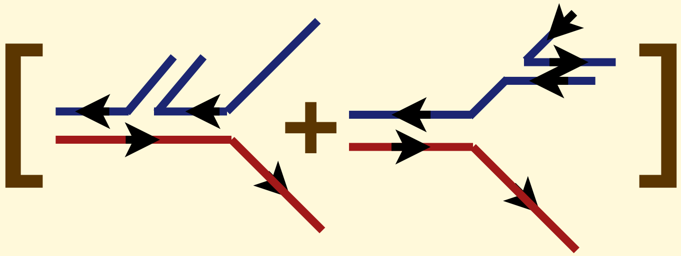
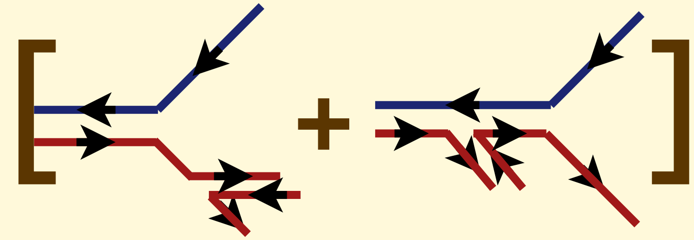
The four terms correspond to the two possible ways colour can flow, and to the two possible emissions for each colour flow:





$$A_{soft} = g\,(\lambda^a\lambda^b)_{ij}\left[\frac{Q\varepsilon}{Qk} - \frac{\bar{p}\varepsilon}{\bar{p}k}\right] + g\,(\lambda^b\lambda^a)_{ij}\left[\frac{p\varepsilon}{pk} - \frac{Q\varepsilon}{Qk}\right]$$

The interference between the two colour structures

 $\propto \; (\lambda^a \lambda^b)_{ij}$  $\propto \; (\lambda^b \lambda^a)_{ij}$

is suppressed by $1/N_c^2$:

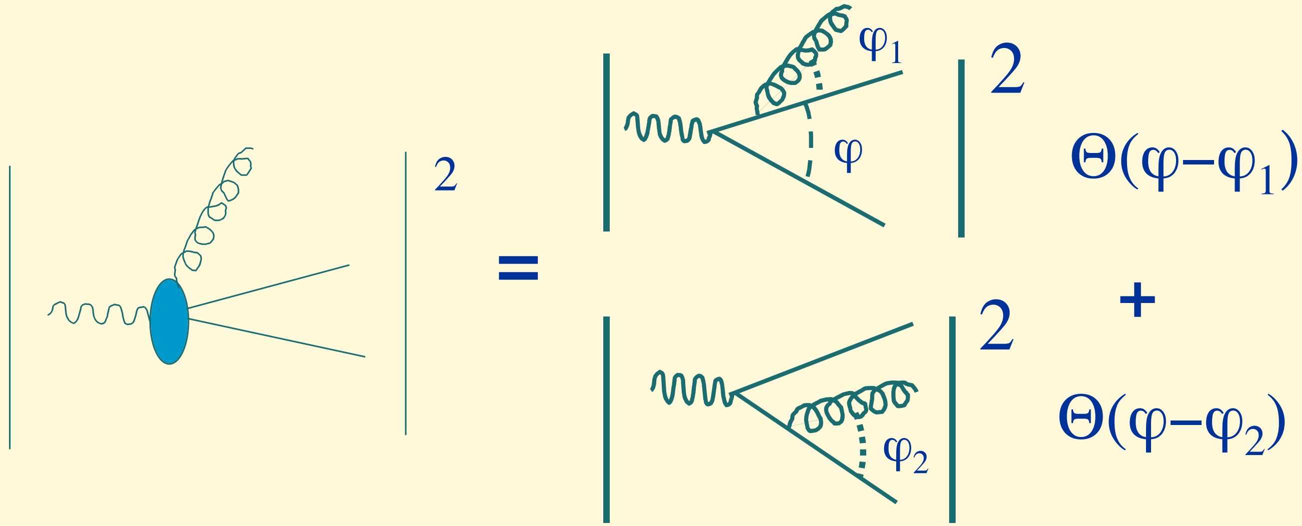
$$\sum_{a,b,i,j} |(\lambda^a \lambda^b)_{ij}|^2 = \sum_{a,b} tr \; (\lambda^a \lambda^b \lambda^b \lambda^a) = \frac{N^2-1}{2} C_F = O(N^3)$$

$$\sum_{a,b,i,j} (\lambda^a \lambda^b)_{ij} [(\lambda^b \lambda^a)_{ij}]^* = \sum_{a,b} tr(\lambda^a \lambda^b \lambda^a \lambda^b) = \frac{N^2-1}{2} \underbrace{(C_F - \frac{C_A}{2})}_{-\frac{1}{2N}} = O(N)$$

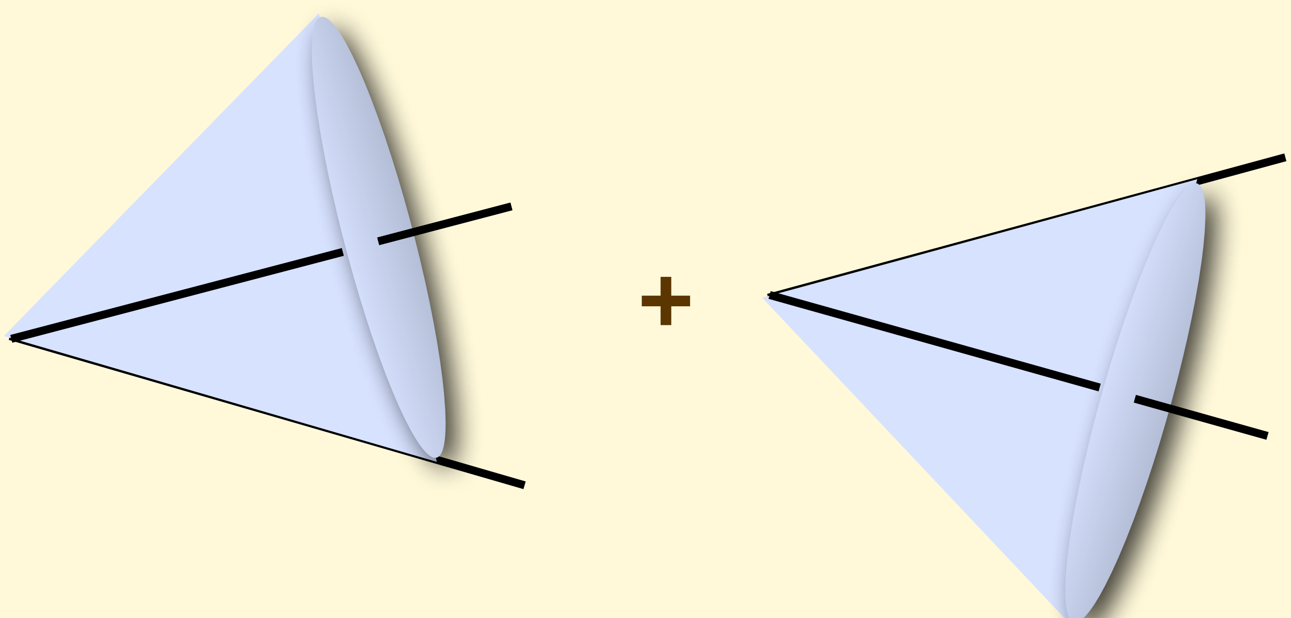
As a result, the emission of a soft gluon can be described, to the leading order in $1/N_c^2$, as the incoherent sum of the emission from the two colour currents

What about the interference between the two diagrams corresponding to the same colour flow? ➥

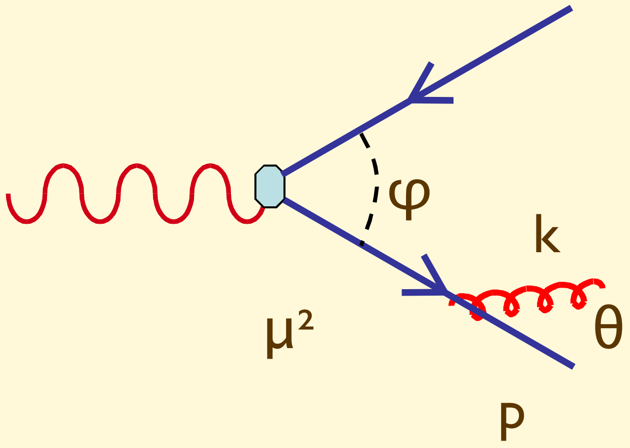
Angular ordering



Radiation inside the cones is allowed, and described by the eikonal probability, radiation outside the cones is suppressed and averages to 0 when integrated over the full azimuth



An intuitive explanation of angular ordering



Lifetime of the virtual intermediate state:

$$\tau < \gamma/\mu = E/\mu^2 = 1 / (k_0\theta^2) = 1/(k_\perp\theta)$$

$\mu^2 = (p+k)^2 = 2E\, k_0\, (1-\cos\theta)$
$\sim E\, k_0\, \theta^2 \sim E\, k_\perp\, \theta$

Distance between q and qbar after τ:

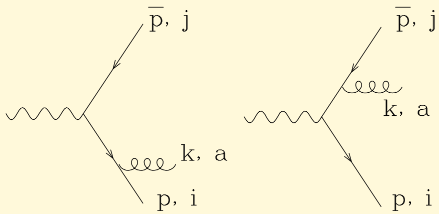
$$d = \varphi\tau = (\varphi/\theta)\ 1/k_\perp$$

If the transverse wavelength of the emitted gluon is longer than the separation between q and qbar, the gluon emission is suppressed, because the q qbar system will appear as colour neutral (=> dipole-like emission, suppressed)

Therefore $d > 1/k_\perp$, which implies

$$\theta < \varphi$$

The formal proof of angular ordering



$$d\sigma_g = \sum |A_{soft}|^2 \frac{d^3k}{(2\pi)^3 2k^0} \sum |A_0|^2 \frac{-2p^\mu p^\nu}{(pk)(pk)} g^2 \sum \varepsilon_\mu \varepsilon_\nu^* \frac{d^3k}{(2\pi)^3 2k^0}$$

$$= d\sigma_0 \frac{\alpha_s C_F}{\pi} \frac{dk^0}{k^0} \frac{d\phi}{2\pi} \frac{1-\cos\theta_{ij}}{(1-\cos\theta_{ik})(1-\cos\theta_{jk})} d\cos\theta$$

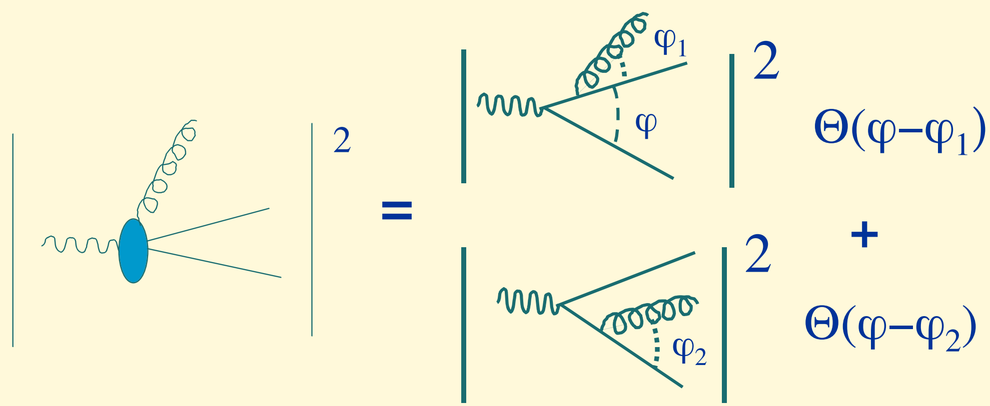
You can easily prove that:

$$\frac{1-\cos\theta_{ij}}{(1-\cos\theta_{ik})(1-\cos\theta_{jk})} = \frac{1}{2}\left[\frac{\cos\theta_{jk}-\cos\theta_{ij}}{(1-\cos\theta_{ik})(1-\cos\theta_{jk})} + \frac{1}{1-\cos\theta_{ik}}\right] + \frac{1}{2}[i \leftrightarrow j] \equiv W_{(i)} + W_{(j)}$$

where:

$$W_{(i)} \to finite\ if\ k \parallel j\ (\cos\theta_{jk} \to 1)$$
$$W_{(j)} \to finite\ if\ k \parallel i\ (\cos\theta_{ik} \to 1)$$

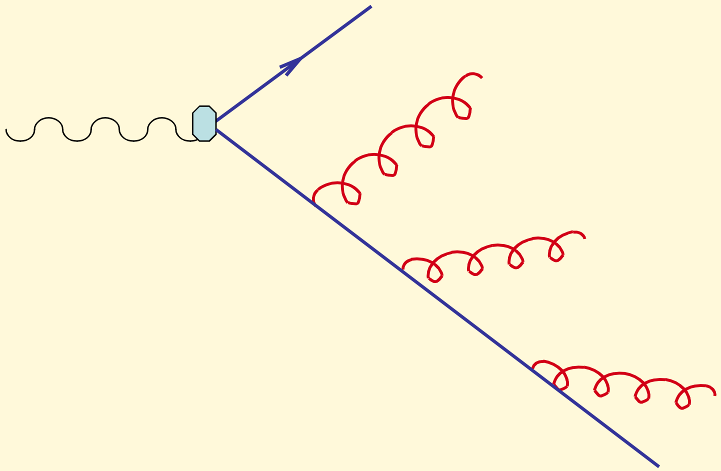
The probabilistic interpretation of W(i) and W(j) is a priori spoiled by their non-positivity. However, you can prove that after azimuthal averaging:



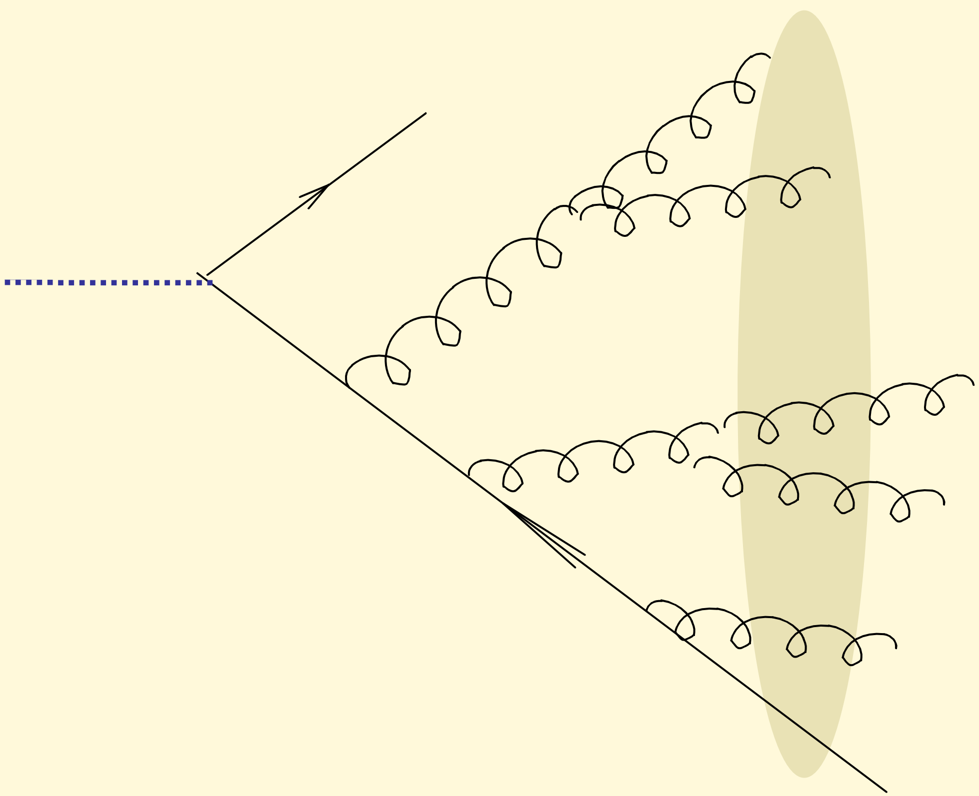
$$\int \frac{d\phi}{2\pi} W_{(i)} = \frac{1}{1-\cos\theta_{ik}} \text{ if } \theta_{ik} < \theta_{ij}\,, \quad 0 \text{ otherwise}$$

$$\int \frac{d\phi}{2\pi} W_{(j)} = \frac{1}{1-\cos\theta_{jk}} \text{ if } \theta_{jk} < \theta_{ij}\,, \quad 0 \text{ otherwise}$$

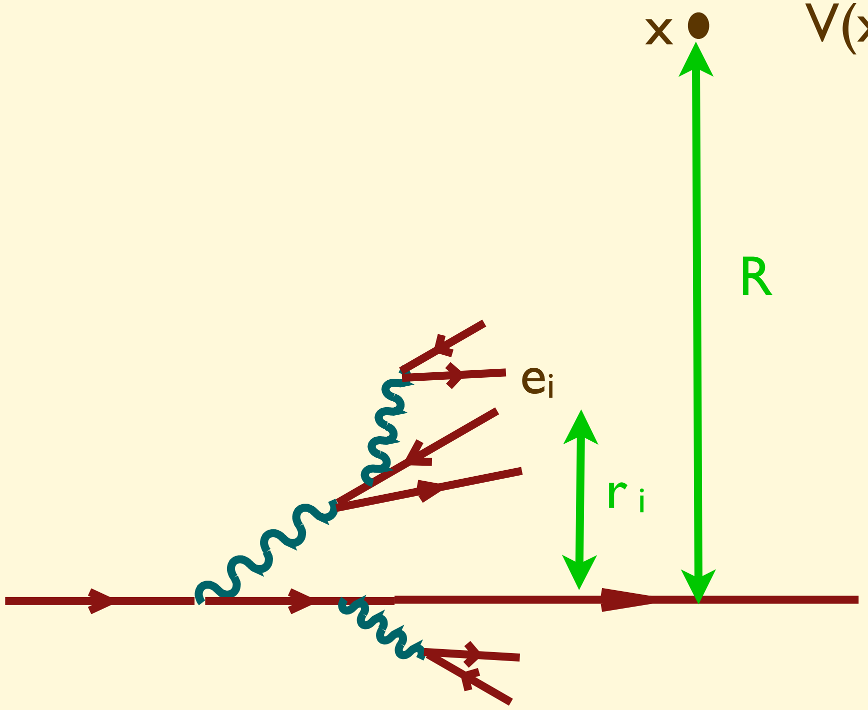
Further branchings will obey angular ordering relative to the new angles. As a result emission angles get smaller and smaller, squeezing the jet



The construction can be iterated to the next emission, with the result that emission angles keep getting smaller and smaller => **jet structure**

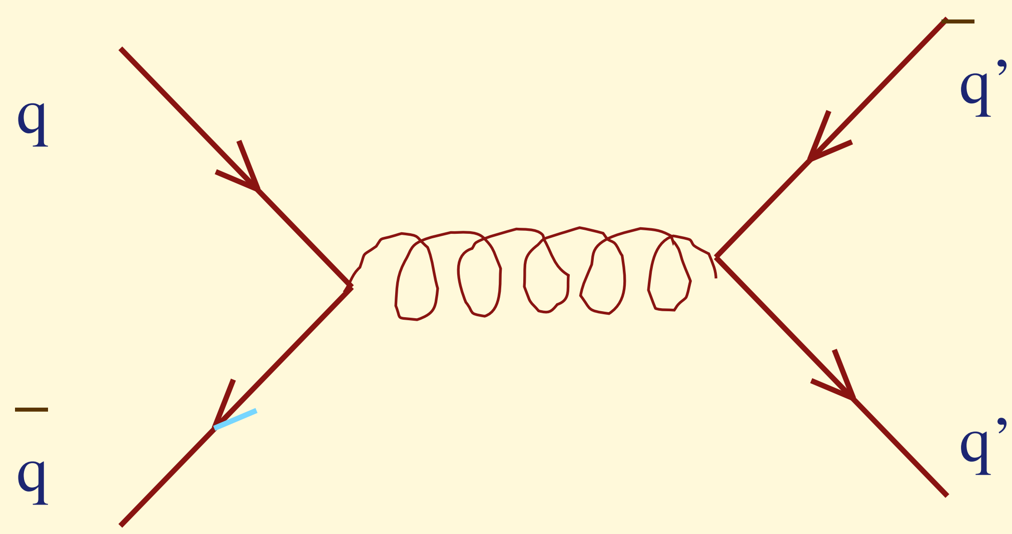


Total colour charge of the system is equal to the quark colour charge. Treating the system as the incoherent superposition of N gluons would lead to artificial growth of gluon multiplicity. Angular ordering enforces coherence, and leads to the proper evolution with energy of particle multiplicities.



$$V(x) = \frac{e}{R} + \frac{1}{R^2} \left| \sum e_i \, \mathbf{r}_i \right|$$

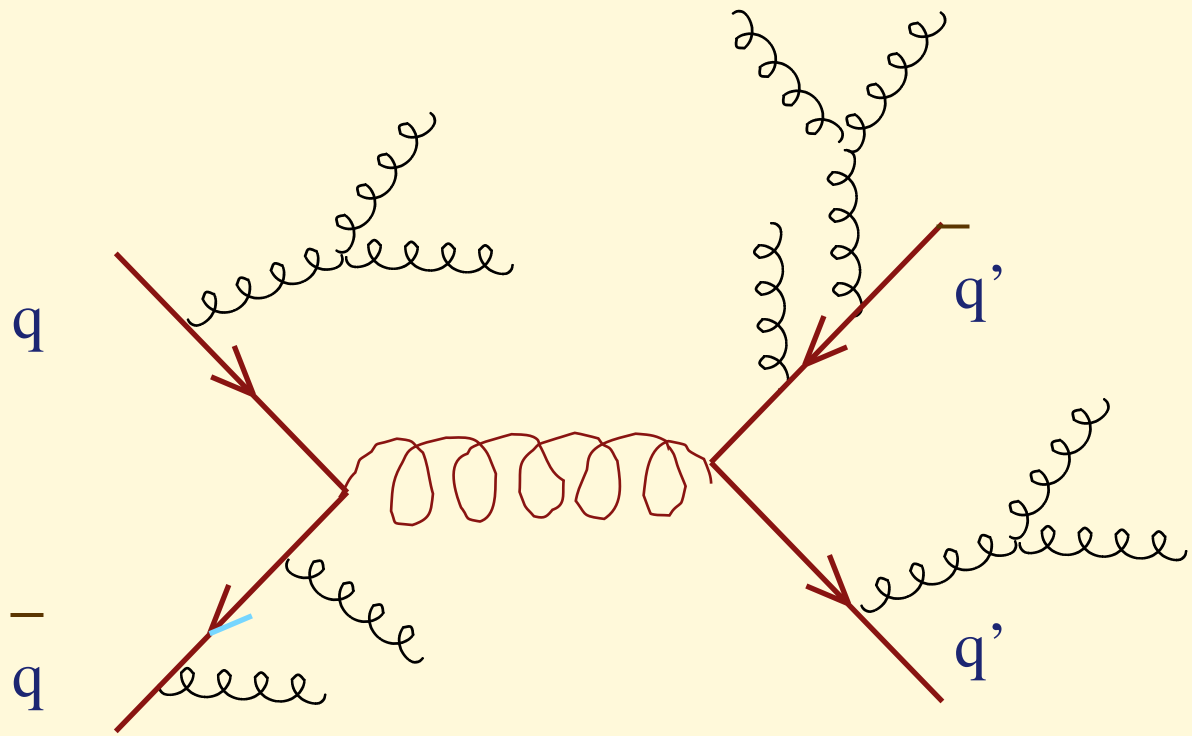
I: Generate the parton-level hard event



II: Develop the parton shower

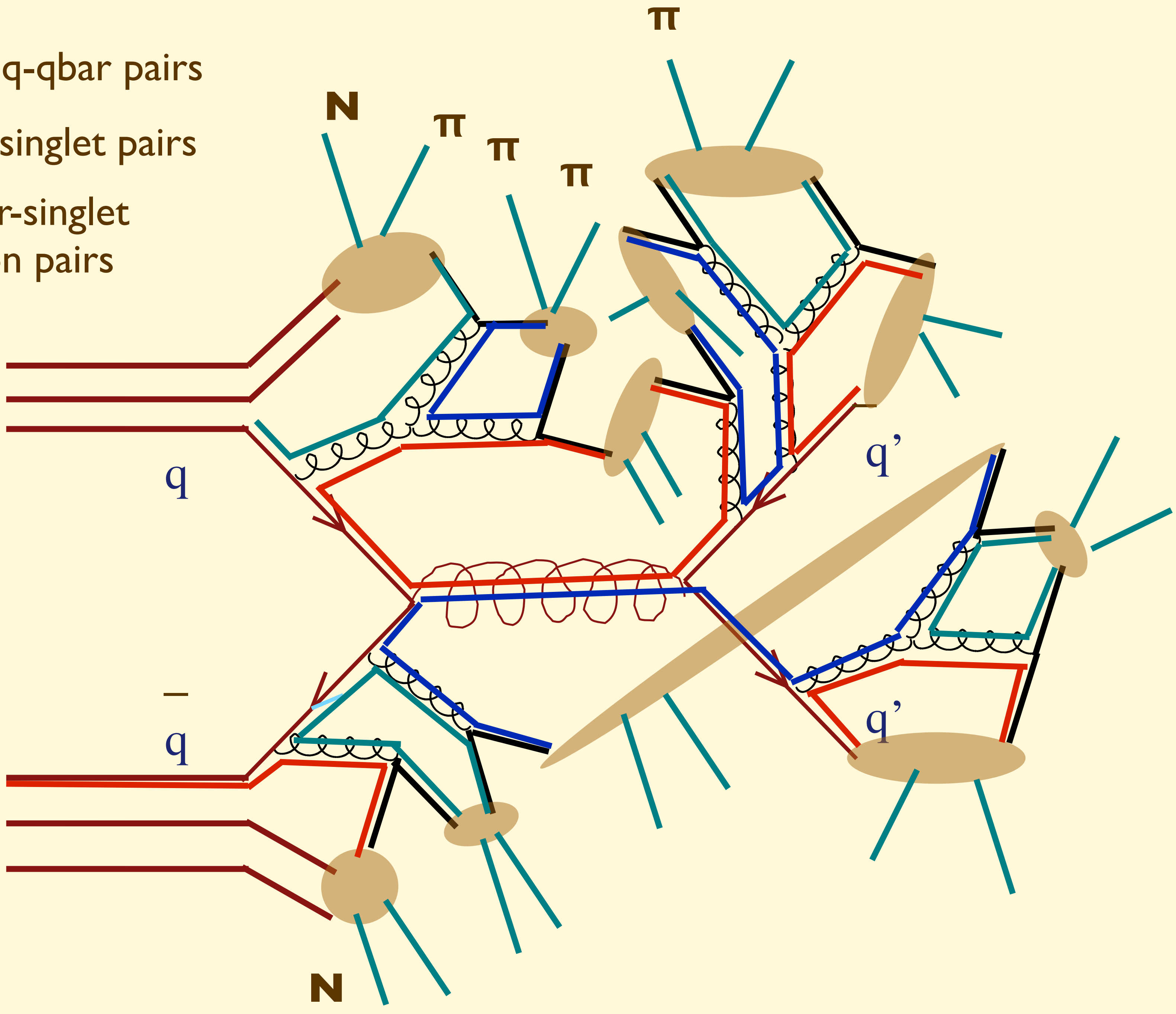
1. Final state

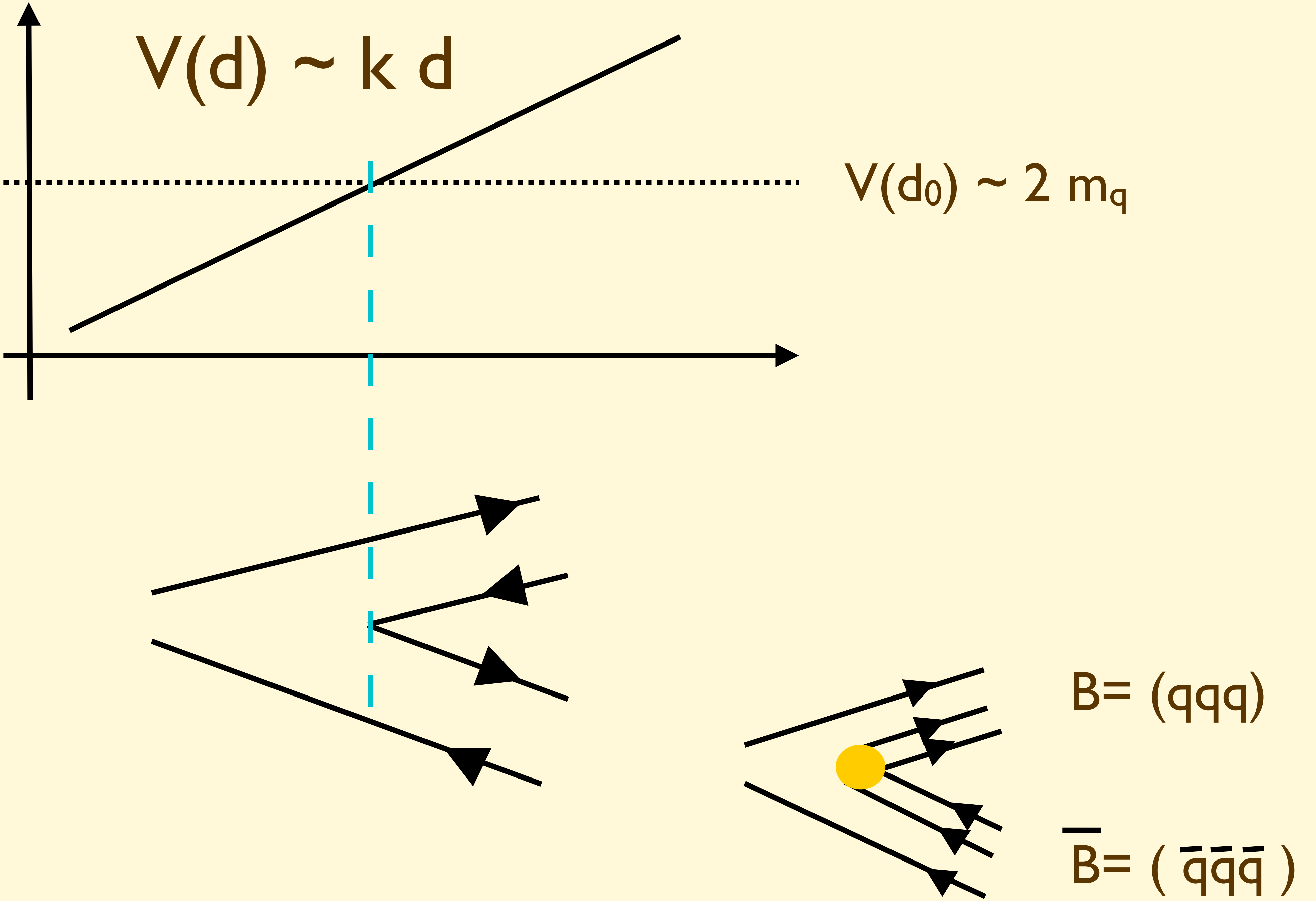
2. Initial state



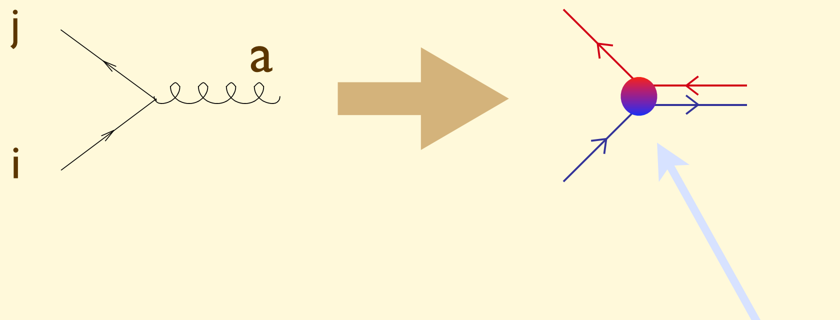
III: Partons Hadronize

1. Split gluons into q-qbar pairs
2. Connect colour-singlet pairs
3. Decay the colour-singlet clusters into hadron pairs

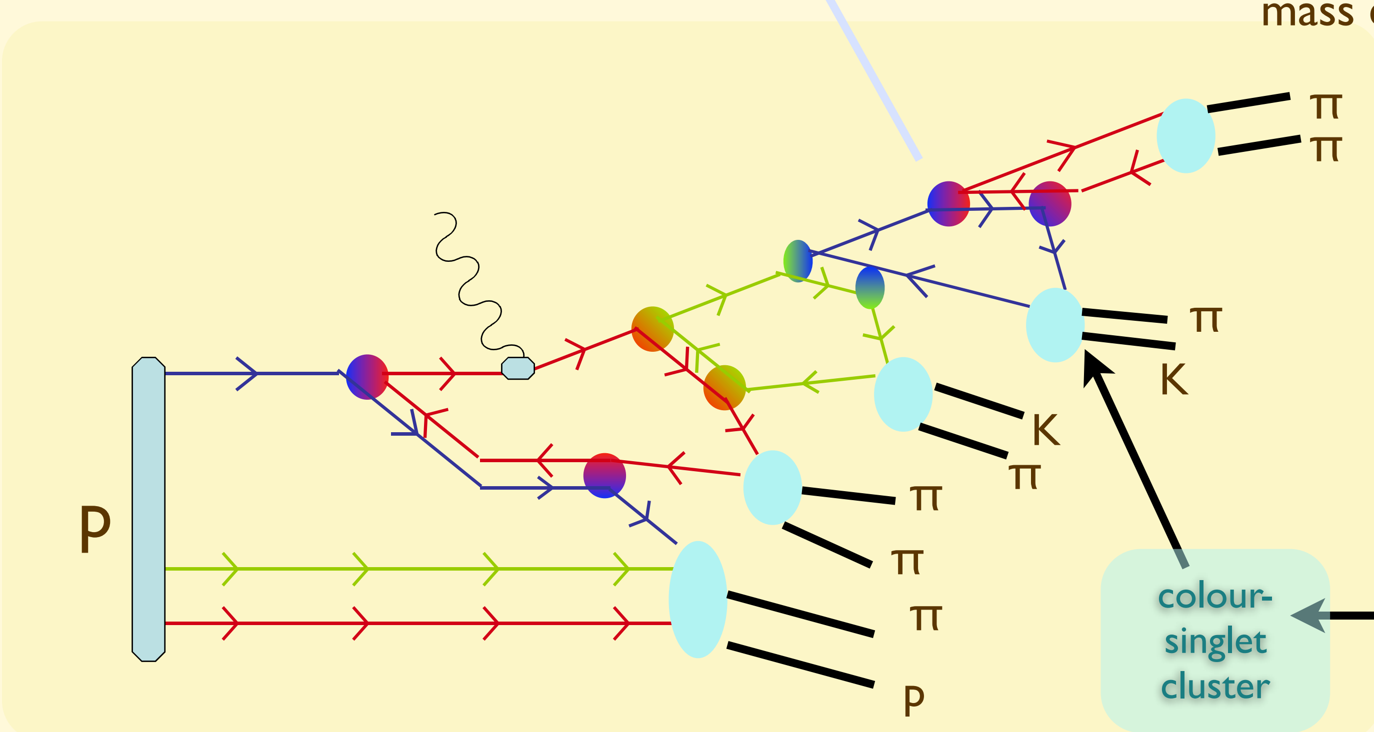


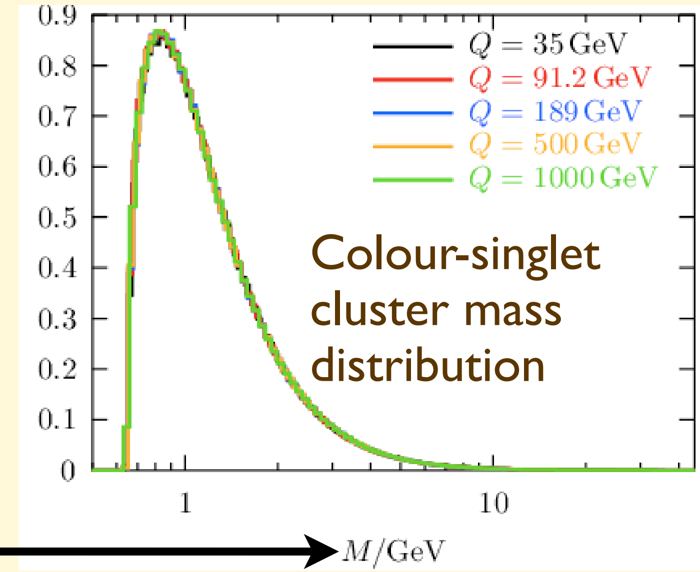


V(d) ~ k d
V(d0) ~ 2 mq
B= (qqq)
B̄= (q̄q̄q̄)

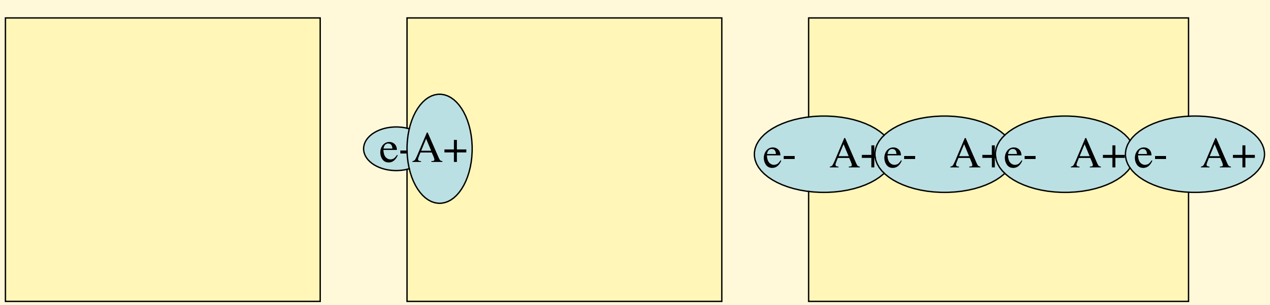


The structure of the perturbative evolution leads naturally to the clustering in phase-space of colour-singlet parton pairs ("preconfinement"). Long-range correlations are strongly suppressed. Hadronization will only act locally, on low-mass colour-singlet clusters.



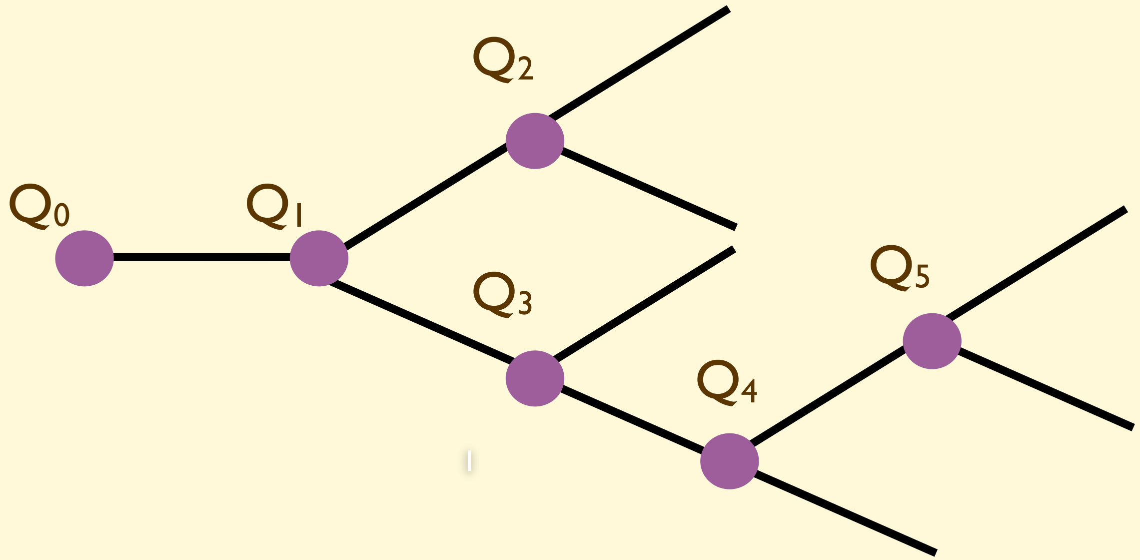


Colour is left "behind" by the struck quark. The first soft gluon emitted at large angle will connect to the beam fragments, ensuring that the beam fragments can recombine to form hadrons, and will allow the struck quark to evolve without having to worry about what happens to the proton fragments.



The shower algorithm

Sequential probabilistic evolution (Markov chain)

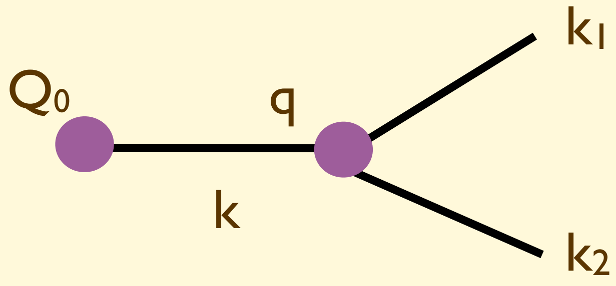


The probability of each emission only depends on the state of the splitting parton, and of the daughters. The QCD dynamics is encoded in these splitting probabilities.

The total probability of all possible evolutions is 1 (unitary evolution).

- The shower evolution does not change the event rate inherited from the parton level, matrix element computation.
- No K-factors from the shower, even though the shower describes higher-order corrections to the leading-order process

Single emission



$$\frac{d\mathrm{Prob}(Q_0 \to q^2)}{dq^2\, dz\, d\phi} = P_0 \frac{\alpha_s(\mu)}{2\pi} \frac{1}{q^2} P(z)$$

$P_0 \Rightarrow \int d\,\mathrm{Prob} = 1$

q^2 ≈ virtuality scale of the branching:

- **$(k_1 + k_2)^2$**
- **$k_1 \cdot k_2$**
- **$k_\perp{}^2$**
- **....**

ϕ = azimuth

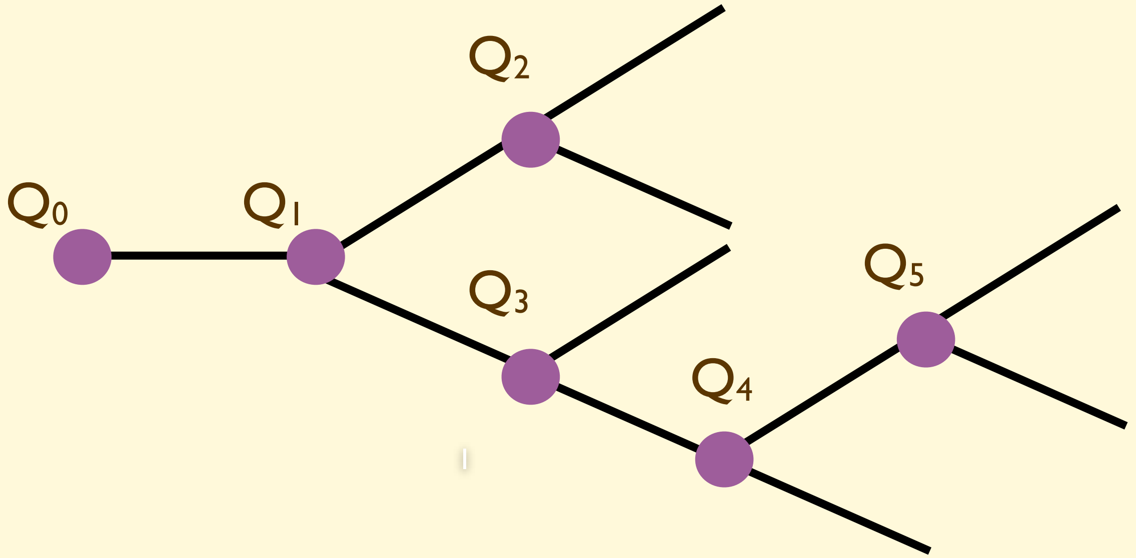
$z=P(k_2)/P(k)$≈ energy/momentum fraction carried by one of the two partons after splitting

- **$P=k^0$**
- **$P=k_{//}$**
- **$P=k_{//} + k^0$**
- **...**

$\mu = f(z,q)$

While at leading-logarithmic order (LL) all choices of evolution variables and of scale for α_s are equivalent, specific choices can lead to improved description of NLL effects and allow a more accurate and easy-to-implement inclusion of angular-ordering constraints and mass effects, as well as to a better merging of multijet ME's with the shower

Multiple emission



$$\text{Prob}(Q_0 \to Q_1) = P_0 \frac{\alpha_s}{2\pi} \int_{Q_1}^{Q_0} \frac{dq^2}{q^2} dz P(z) d\phi$$

$$\text{Prob}(Q_0 \to Q_1 \to Q_2) = P_0 \frac{\alpha_s}{2\pi} \int_{Q_1}^{Q_0} \frac{dq^2}{q^2} dz P(z) d\phi \frac{\alpha_s}{2\pi} \int_{Q_2}^{Q_1} \frac{dq^2}{q^2} dz P(z) d\phi$$

$$\sim P_0 \frac{1}{2!} [\frac{\alpha_s}{2\pi} \int_{Q_2}^{Q_0} \frac{dq^2}{q^2} dz P(z) d\phi]^2$$

$$\text{Prob}(Q_0 \to X) = P_0 \times \sum \frac{1}{n!} [\frac{\alpha_s}{2\pi} \int_{\Lambda}^{Q_0} \frac{dq^2}{q^2} dz P(z) d\phi]^n = 1$$

Λ=infrared cutoff

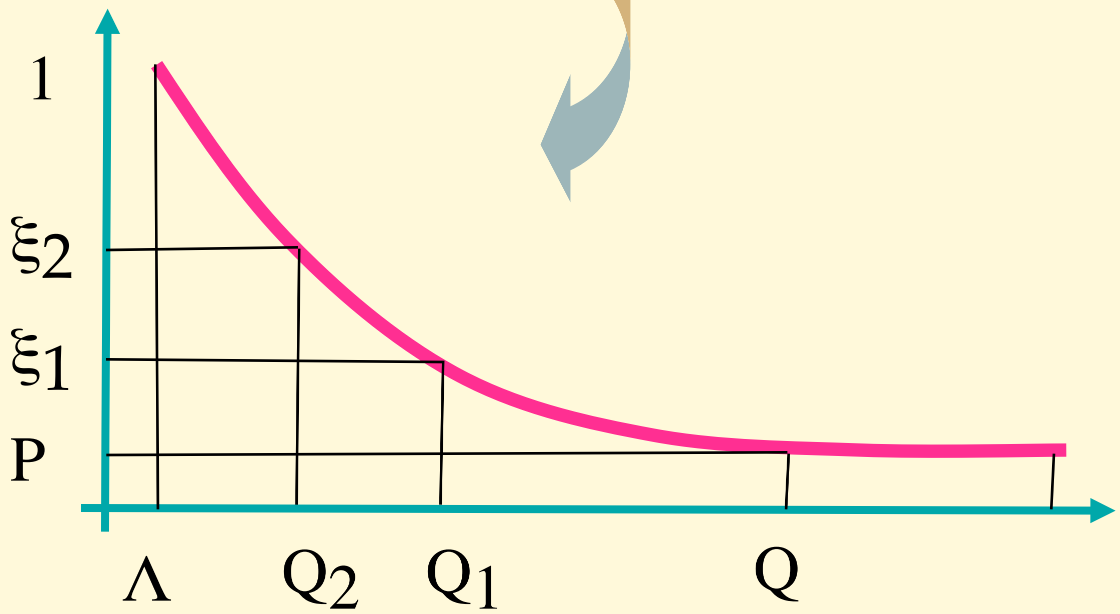
$$P_0 = \exp\{-\frac{\alpha_s}{2\pi} \int_{\Lambda}^{Q_0} \frac{dq^2}{q^2} dz P(z) d\phi\}$$

P_0 = Sudakov form factor
~ probability of no emission between the scale Q_0 and Λ

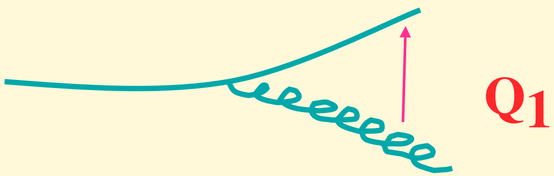
Generation of splittings

$$P(Q,\Lambda) = exp\left[-\int_\Lambda^Q \frac{dq^2}{q^2}\frac{\alpha_s(q)}{2\pi}P(z)dz\right]$$

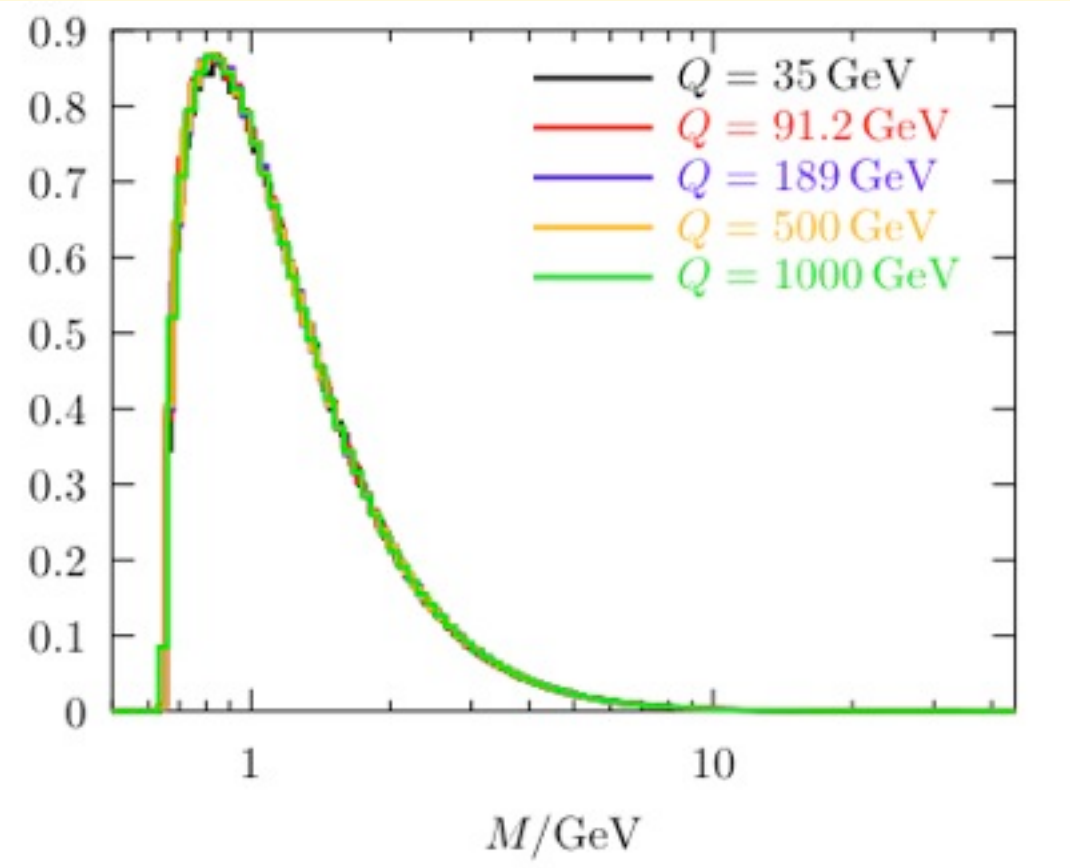
prob. of no radiation between Q and Λ



1. Generate $0< \xi_1 <1$
2. If $\xi_1 < P(Q, \Lambda) \Rightarrow$ no radiation, q' goes directly on-shell at scale Λ≈GeV
3. Else
   1. calculate Q_1 such that $P(Q_1,\Lambda)= \xi_1$
   2. emission at scale Q_1:



4. Select z according to P(z)
5. Reconstruct the full kinematics of the splitting
6. Go back to 1) and reiterate, until shower stops in 2). At each step the probability of emission gets smaller and smaller



The existence of high-mass clusters, however rare, is unavoidable, due to IR cutoff which leads to a non-zero probability that no emission takes place. This is particularly true for evolution of massive quarks (as in, e.g. Z→bb or cc). Prescriptions have to be defined to deal with the "evolution" of these clusters. **This has an impact on the z→1 behaviour of fragmentation functions.**

Phenomenologically, this leads to uncertainties, for example, in the background rates for H→γγ (jet→γ).

This approach is extremely successful in describing the properties of hadronic final states!

Ex:
Particle multiplicities:

Particle	Experiment	Measured	Old model	Herwig (C^{++})	Herwig (Fortran)
All charged	M, A, D, L, O	20.924 ± 0.117	20.22*	20.814	20.532*
γ	A, O	21.27 ± 0.6	23.03	22.67	20.74
π^0	A, D, L, O	9.59 ± 0.33	10.27	10.08	9.88
$\rho(770)^0$	A, D	1.295 ± 0.125	1.235	1.316	1.07
$\pi^\pm$	A, O	17.04 ± 0.25	16.30	16.95	16.74
$\rho(770)^\pm$	O	2.4 ± 0.43	1.99	2.14	2.06
η	A, L, O	0.956 ± 0.049	0.886	0.893	0.669*
$\omega(782)$	A, L, O	1.083 ± 0.088	0.859	0.916	1.044
$\eta(958)$	A, L, O	0.152 ± 0.03	0.13	0.136	0.106
K^0	S, A, D, L, O	2.027 ± 0.025	2.121*	2.062	2.026
$K^*(892)^0$	A, D, O	0.761 ± 0.032	0.667	0.681	0.583*
$K^*(1430)^0$	D, O	0.106 ± 0.06	0.065	0.079	0.072
$K^\pm$	A, D, O	2.319 ± 0.079	2.335	2.286	2.250
$K^*(892)^\pm$	A, D, O	0.731 ± 0.058	0.637	0.657	0.578
$\phi(1020)$	A, D, O	0.097 ± 0.007	0.107	0.114	0.134*
p	A, D, O	0.991 ± 0.054	0.981	0.947	1.027
Δ^{++}	D, O	0.088 ± 0.034	0.185	0.092	0.209*
Σ^-	O	0.083 ± 0.011	0.063	0.071	0.071
Λ	A, D, L, O	0.373 ± 0.008	0.325*	0.384	0.347*
Σ^0	A, D, O	0.074 ± 0.009	0.078	0.091	0.063
Σ^+	O	0.099 ± 0.015	0.067	0.077	0.088
$\Sigma(1385)^\pm$	A, D, O	0.0471 ± 0.0046	0.057	0.0312*	0.061*
Ξ^-	A, D, O	0.0262 ± 0.001	0.024	0.0286	0.029
$\Xi(1530)^0$	A, D, O	0.0058 ± 0.001	0.026*	0.0288*	0.009*
Ω^-	A, D, O	0.00125 ± 0.00024	0.001	0.00144	0.0009*
$f_2(1270)$	D, L, O	0.168 ± 0.021	0.113	0.150	0.173
$f_2'(1525)$	D	0.02 ± 0.008	0.003	0.012	0.012
$D^\pm$	A, D, O	0.184 ± 0.018	0.322*	0.319*	0.283*
$D^*(2010)^\pm$	A, D, O	0.182 ± 0.009	0.168	0.180	0.151*
D^0	A, D, O	0.473 ± 0.026	0.625*	0.570*	0.501
$D_s^\pm$	A, O	0.129 ± 0.013	0.218*	0.195*	0.127
$D_s^{*\pm}$	O	0.096 ± 0.046	0.082	0.066	0.043
J/Ψ	A, D, L, O	0.00544 ± 0.00029	0.006	0.00361*	0.002*
Λ_c^+	D, O	0.077 ± 0.016	0.006*	0.023*	0.001*
$\Psi'(3685)$	D, L, O	0.00229 ± 0.00041	0.001*	0.00178	0.0008*

* Indicates a prediction that differs from the measured value by more than three standard deviations.

Tabl 1. Average particle multiplicities per event in e^+e^- collisions at 91.2 GeV. Experimental data were measured by the following collaborations at LEP and at SLC: ALEPH(A), DELPHI(D), L3(L), OPAL(O), MARK2(M), and SLD(S). The theoretical predictions in the last three columns, taken from Ref. [30], correspond to various implementations of the cluster hadronization model (see Ref. [30] for details).

30. Gieseke S et al. *JHEP* (02) 005 (2004)

Ex: Energy distributions

(Winter, Krauss, Soff,
hep-ph/0311085)