

# Symmetries, Duality, and the Unity of Physics

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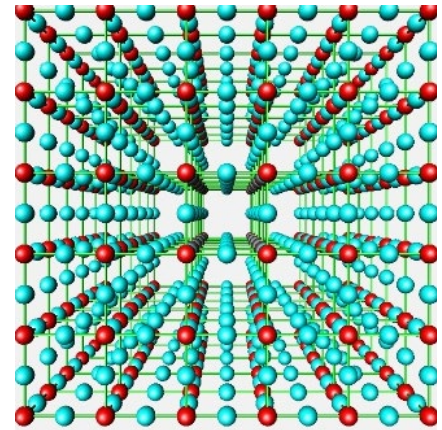
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# Symmetries



# Physicists love symmetries

- Galileo, Lorentz, Poincare
- Crystallography
- Global symmetry (flavor)
  - e.g. Isospin  $SU(2)$ ,  $SU(3)$
- Gauge (local) symmetry
  - $A_\mu \rightarrow A_\mu + \partial_\mu \Lambda$  with  $\Lambda(x)$  in Maxwell theory
  - $g_{\mu\nu} \rightarrow g_{\mu\nu} + \partial_\mu \xi^\rho g_{\rho\nu} + \partial_\nu \xi^\rho g_{\mu\rho} + \xi^\rho \partial_\rho g_{\mu\nu}$  with  $\xi^\rho(x)$  in General Relativity
  - Color  $SU(3)$ , the standard model  $SU(3) \times SU(2) \times U(1)$



# Gauge symmetry is deep

- Largest symmetry (a group for each point in spacetime)
- Useful in making the theory manifestly Lorentz invariant, unitary, and local (and hence causal)
- Appears in
  - Maxwell theory, the Standard Model of particle physics
  - General Relativity
  - Many condensed matter systems
  - Deep mathematics (fiber bundles)

# But

- Because of Gauss law the Hilbert space is gauge invariant.
- Hence: gauge symmetry is not a symmetry.
  - It does not act on anything.
  - All the operators are gauge invariant.
- A better phrase is gauge redundancy.

# Gauge symmetry is not a symmetry

- Not surprising to mathematicians and some physicists
  - Manifestly true in many lattice constructions in condensed matter physics
  - Manifestly true in some low-dimensional continuum quantum field theories
- In particle physics, perhaps we should not look for a more complete theory with a larger gauge group
- Of course, this does not mean that gauge symmetry is not an extremely powerful concept and a useful tool.

# Emergent symmetries

Global symmetries can emerge as accidental symmetries at long distances.

- Then they are approximate.
  - Continuous translations and rotations even though the microscopic theory is on a lattice.
  - Parity and time reversal symmetries are approximate symmetries at low energy, even though they are not symmetries of the standard model of particle physics
  - Approximate baryon number symmetry (and the related stability of the proton)

# Emergent symmetries

Gauge symmetries (redundancies) can be emergent.

- They must be exact because they are not symmetries.
- First examples
  - lattice systems (mostly in condensed matter)
  - FQHE
  - continuum quantum field theory in low dimensions
    - These systems do not have a massless spin one particle reflecting the emergent gauge symmetry.
- We will soon discuss examples in higher dimensions



# Symmetries in a theory of gravity

General considerations based on black holes strongly suggest that in any theory of quantum gravity

- There cannot be any exact **global symmetry**
  - Approximate, accidental global symmetries are possible.
- **Gauge symmetries** are possible.
  - The spectrum must include excitations with all electric and magnetic charges that are compatible with the symmetry.
    - Actually, this follows from the previous point.

# Global vs. Local Symmetries

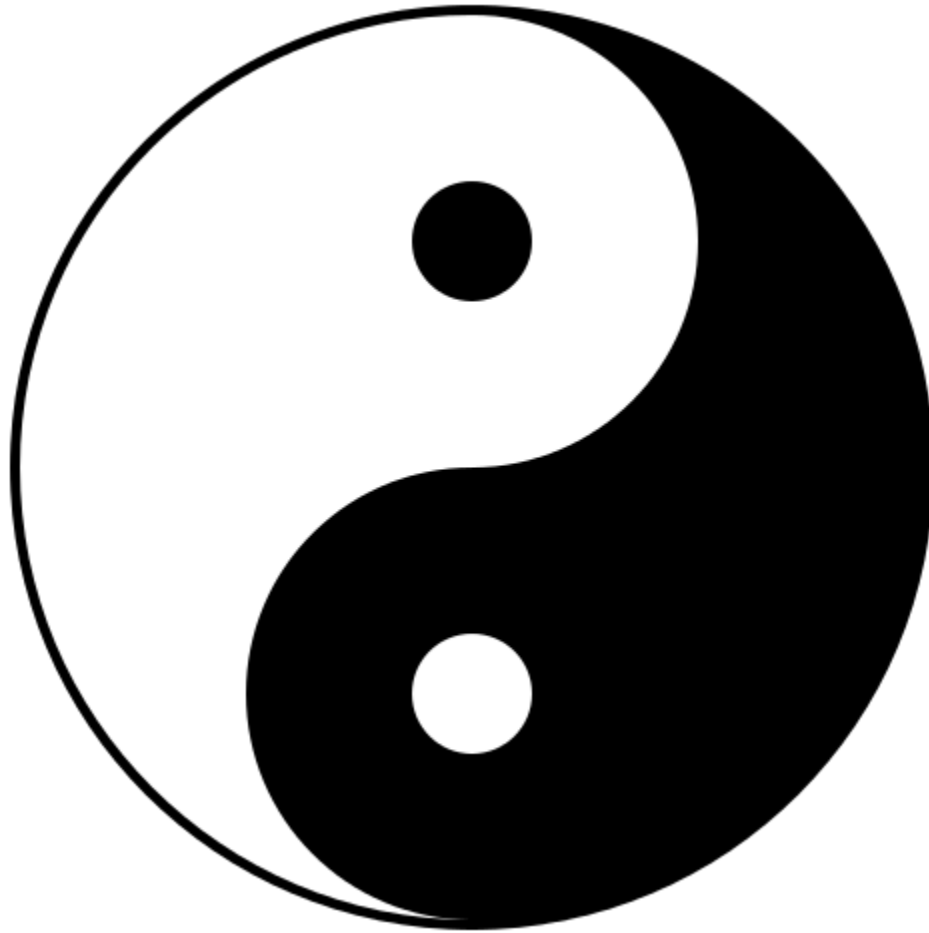
## Global

- Intrinsic
- Can be accidental at long distances – approximate
- Classify operators
- Can be spontaneously broken
- If unbroken can classify states
- Useful in classifying phases
- 't Hooft anomalies
- Not present in a theory of gravity

## Local (gauge)

- Ambiguous – duality
- Can emerge at long distances – exact
- All operators are invariant
- Not really a symmetry
- Hence it cannot be broken (Higgs description meaningful only at weak coupling)
- Cannot be anomalous
- Appears essential in formulating the Standard Model and in Gravity

# Duality



# Simple examples of duality

Simple dualities in free (solvable) theories (easy to establish)

Harmonic oscillator

$$H = \frac{1}{2m} p^2 + \frac{k}{2} q^2$$

Map  $H$  to itself:

- Classically, a canonical transformation ( $q \rightarrow \frac{p}{\sqrt{mk}}$ ,  
 $p \rightarrow -\sqrt{mk}q$ )
- Quantum mechanically, a Fourier transform

# Duality and emergent gauge symmetry

Again, easy to establish (essentially Fourier transform):

A free scalar field in  $2+1d$  is dual to Maxwell theory

$$F_{\mu\nu} \sim \epsilon_{\mu\nu\rho} \partial^\rho \phi$$

- The vector potential  $A_\mu$  is related in a non-local way to  $\phi$ .
- **No gauge symmetry in the description using  $\phi$**
- The equation of motion of  $F_{\mu\nu}$  is trivial in terms of  $\phi$ .
- The Bianchi identity of  $F_{\mu\nu}$  is the equation of motion of  $\phi$ .

# Duality and emergent gauge symmetry

Again, easy to establish (essentially Fourier transform):

Maxwell theory in  $3+1d$  is dual to a magnetic Maxwell theory ( $E \sim \tilde{B}$ ,  $B \sim -\tilde{E}$ )

$$F_{\mu\nu} \sim \frac{1}{2} \epsilon_{\mu\nu\rho\sigma} \tilde{F}^{\rho\sigma}$$

- The equation of motion of  $F_{\mu\nu}$  is the Bianchi identity of  $\tilde{F}_{\mu\nu}$  and vice versa.
- The new vector potential  $\tilde{A}_\mu$  is related in a non-local way to the original vector potential  $A_\mu$ .
- Its gauge symmetry is emergent (not present in the original formulation).

# Simple examples of duality

In all these examples duality is achieved by a Fourier transform.

Small fluctuations in  $q$   $\leftrightarrow$  large fluctuations in  $p$

Small fluctuations in  $p$   $\leftrightarrow$  large fluctuations in  $q$

Both descriptions are correct, but depending on the problem one description can be more useful than another.

# More examples of duality

These examples involve more subtle (interacting) theories

- Kramers–Wannier duality in the Ising model and its generalizations in higher dimensions including gauge systems
- Many examples in  $1+1d$ , especially in the context of bosonization, conformal field theory, etc.

As the previous examples, these can be established rigorously.



# More examples of duality

In the simple examples

large fluctuations  $\leftrightarrow$  small fluctuations

In the more subtle (interacting) examples

strong coupling  $\leftrightarrow$  weak coupling

large  $\hbar$   $\leftrightarrow$  small  $\hbar$

- Useful in solving complicated systems
- Duality is an intrinsically quantum mechanical phenomenon.

# Duality in $3+1d$ $\mathcal{N} = 4$ supersymmetry

- This is a scale invariant interacting theory of gluons characterized by a gauge group  $G$  and coupling constant  $\alpha$  (like the fine structure constant)
- The theory with  $G$  and  $\alpha$  is equivalent to the theory with another group  $\tilde{G}$  and  $\frac{1}{\alpha}$ 
  - Exchanges strong and weak coupling – large and small fluctuations
  - The gluons of  $\tilde{G}$  are magnetic monopoles of  $G$ . As above, the map between them is non-local.

# Duality in $3+1d$ $\mathcal{N} = 4$ supersymmetry

$$G, \quad \alpha \leftrightarrow \tilde{G}, \quad \frac{1}{\alpha}$$

- The duality is an exact equivalence of theories. Like the Fourier transform above, but cannot be performed explicitly (cannot prove it)
  - Same spectrum of states
  - Same spectrum of operators
  - Same correlation functions

# Duality in $3+1d$ $\mathcal{N} = 4$ supersymmetry

- The gauge symmetry of the dual description is emergent!
- Which of the two gauge symmetries is fundamental?
- Perhaps neither gauge symmetry is fundamental.
- Which set of gluons is elementary?
- Notion of “elementary particle” is ill-defined.

# Another notion of duality

The previous examples are **exact dualities** – two different descriptions of the same physics.

An alternate notion: two theories with the **same long distance (low energy) physics**

Hamiltonian at short distances (e.g. many electrons with some interactions)

Effective Hamiltonian describing the long distance physics (typically different degrees of freedom)



# Another notion of duality

Many known examples

- Sometimes we can identify the long distance degrees of freedom using the short distance variables; e.g.
  - Ginzburg–Landau theory of superconductivity
  - QCD at short distances; pions at long distances
- Sometimes the relation is not so easy – it might be non-local
  - The examples of emergent gauge symmetry
  - Particle-vortex duality in  $2+1d$
  - Emergent gauge fields in the FQHE

# Duality in interacting field theories

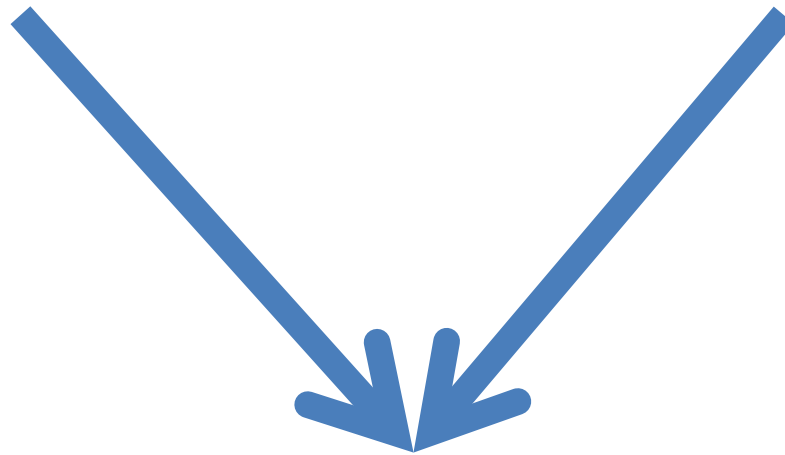
## $\mathcal{N} = 1$ supersymmetry

Electric theory

$G$

Magnetic theory

$\tilde{G}$



Non-trivial fixed point at long distances

# Duality in interacting field theories

## $\mathcal{N} = 1$ supersymmetry

Another option:

Electric theory

Based on  $G$



Approximately free theory (IR free)

Based on  $\tilde{G}$



# Duality in interacting field theories

## $\mathcal{N} = 1$ supersymmetry

At short distances an asymptotically free theory based on  $G$

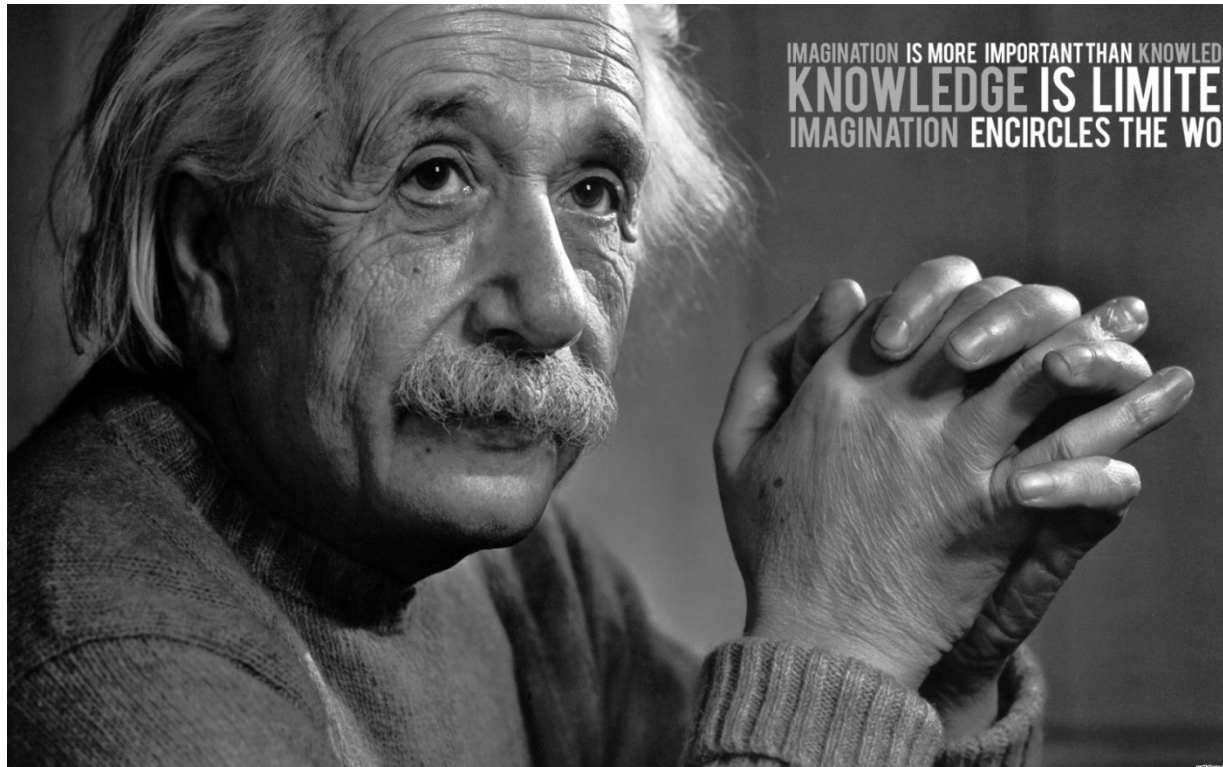
At long distances an almost free theory based on  $\tilde{G}$

Unlike ordinary  $QCD$ , at long distances this theory has a non-Abelian gauge theory –  $\widetilde{QCD}$ .

- The quarks and gluons of  $\widetilde{QCD}$  are composite. They can be thought of as magnetic monopoles of  $G$
- Their gauge symmetry is emergent.

Many more examples. Duality and emergent gauge symmetries are ubiquitous.

# Other Dualities and the Unity of Physics



# Emergent general covariance and emergent spacetime

- So far we discussed duality between two field theories
- String-string duality
  - T-duality
  - S-duality
  - U-duality
- String-fields duality
  - Matrix models for low dimensional string theories
  - BFSS M(atr ix) model
  - AdS/CFT
  - More generally gauge-gravity duality

# Simple Boson/Fermion dualities

- Bosnization in  $1+1d$ 
  - A theory of fermions is equivalent to a theory of bosons
  - No notion of spin of particles in  $1+1d$  (only spin of operators)
- Spin and statistical transmutation in  $2+1d$ . Flux attachment in the FQHE
  - Massive particles coupled to gauge fields with special interactions (Chern-Simons coupling)  
bosons  $\leftrightarrow$  fermions    or    bosons  $\leftrightarrow$  bosons'  
or    fermions  $\leftrightarrow$  fermions'

# Dualities in 2+1 dimensions

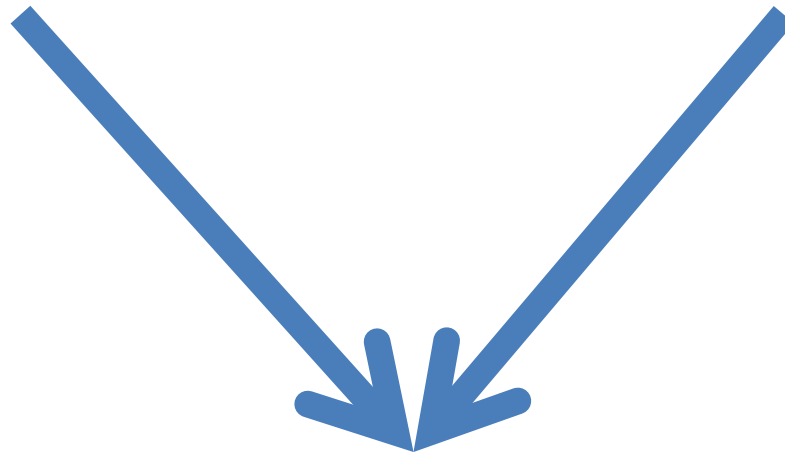
- Boson  $\leftrightarrow$  boson coupled to a gauge field –  
Particle/vortex duality [Peskin; Dasgupta, Halperin]
- New boson  $\leftrightarrow$  fermion dualities [... Chen, Fisher, Wu; Barkeshli, McGreevy; Giombi, Minwalla, Prakash, Trivedi, Wadia, Yin; Aharony; Karch, Tong; NS, Senthil, Wang, Witten; Hsin, NS; Aharony, Benini, Hsin, NS; ...]
- New fermion  $\leftrightarrow$  fermion dualities [... Son; NS, Senthil, Wang, Witten; Xu, You; Hsin, NS; Gomis, Komargodski, NS; ...]

# New dualities in 2+1 dimensions

Many  $2+1d$  examples of interacting theories of massless matter coupled to gauge fields

Bosons and/or fermions  
+ gauge fields

Bosons and/or fermions  
+ gauge fields



Non-trivial fixed point at long distances

# New dualities in 2+1 dimensions

Many  $2+1d$  examples of interacting theories of bosons and/or fermions at short distances flowing to free massless Fermions

Bosons and/or fermions + gauge fields



Free Fermions

# Motivation (unity of physics)

- Well established particle/vortex dualities
- Previously found dualities in supersymmetric theories
  - Many tests
- String duality
  - Many tests
- AdS/CFT and large  $N$ 
  - Same bulk theory is dual to different boundary theories
  - Checks at large  $N$
- Some suggestions in the condensed matter literature



# Checks and relations

- Can use these new dualities to derive well established dualities
- Relation to other conjectured dualities, which were subjected to many checks
- Relation to mathematics
  - Mirror symmetry
  - Level/rank duality
- Not independent: assuming some of these dualities are right we can derive others – **a web of dualities**

# Applications in Condensed Matter Physics

- Fractional Quantum Hall Effect
- Physics of first Landau level at half filling
- Gapped phases of topological insulators and topological superconductors
- ...

# Symmetries in Quantum Gravity and in Condensed Matter Physics

- No exact **global symmetries**
  - In gravity because of the physics of black holes
  - In condensed matter physics, if no exact global symmetry at short distances, all global symmetries at long distances are approximate
- Can have emergent **gauge symmetries** with emergent gauge fields.
  - Then there must be excitations with all allowed charges.

# Conclusions

- **Symmetries** are common in physics
- Both in a theory of gravity and in condensed matter physics (unity of physics)
  - No exact global symmetries
  - Exact gauge symmetries can be emergent (all charges must be present)
- **Dualities** are common in physics: two (or more) different descriptions of the same physics

# Conclusions

- **The unity of physics.** Insights from different branches of physics have recently converged leading to new dualities in  $2+1d$ .
  - Better understanding of known phenomena and mechanisms.
  - New phenomena and mechanisms.

# Conclusions

- Gauge symmetry is not fundamental.
  - It is often convenient to use it to make the description manifestly Lorentz invariant, unitary, and local.
  - But there can be different such (dual) descriptions.
- Look for a reformulation of quantum field theory, which makes the duality manifest.
  - We should not be surprised by duality!